

Bruno Maximovich Pontecorvo

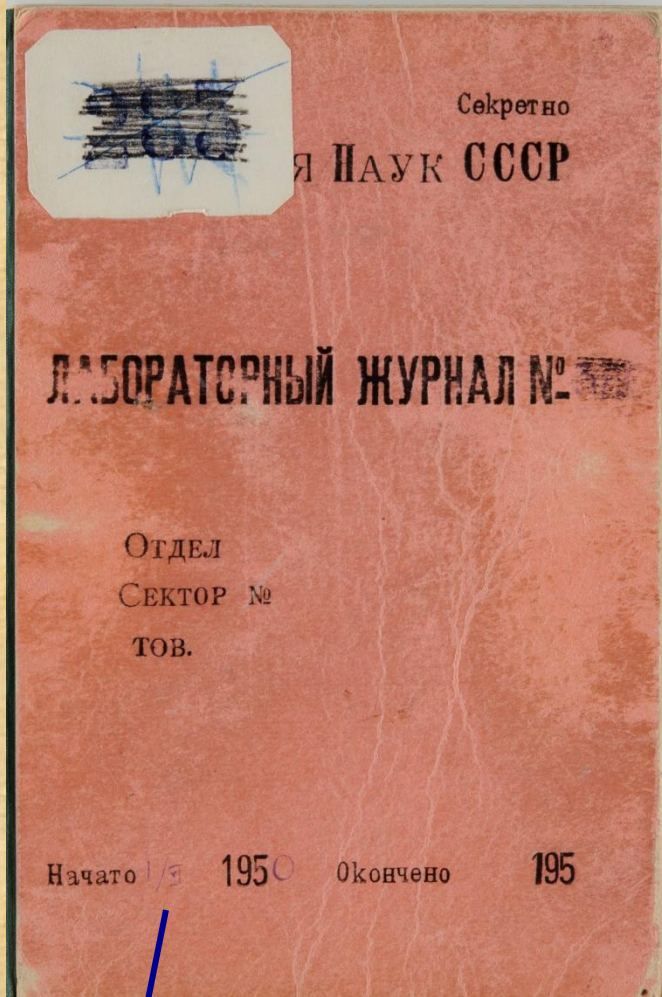
i due quaderni di appunti durante il primo anno e mezzo di lavoro a Dubna



SIF - Pisa
25 Settembre 2014

Gloria Spandre
Rino Castaldi
INFN-Pisa

I/XI 1950 Primo Quaderno: da 01/11/1950 a ≤ 30/11/1950
da 14/09/1950 a ≥ 24/03/1952

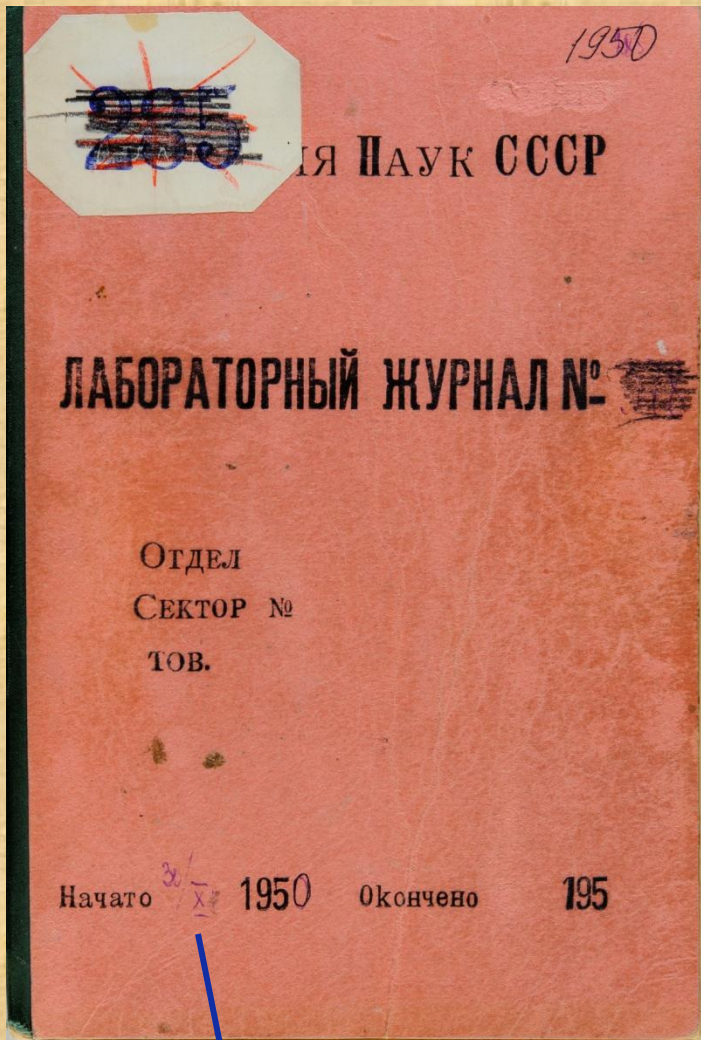


Lo scorso anno, nel centenario della nascita di Bruno Pontecorvo, abbiamo avuto, grazie al figlio Gil, un quaderno inedito di appunti che lo scienziato scrive di proprio pugno in lingua inglese nel suo primo anno e mezzo di attività di ricerca presso il Laboratorio di Problemi Nucleari di Dubna. La copertina e la prima pagina del quaderno riportano la data del primo novembre 1950, proprio il giorno in cui Pontecorvo inizia il suo lavoro al Laboratorio di Dubna.

Pontecorvo dopo solo nove pagine smette di scrivere i suoi appunti su questo quaderno per poi riprendere a scrivervi dall'ultima pagina, capovolgendolo, solo l'anno seguente, il 14 settembre 1951.

1/11 1950

30/XI 1950 Secondo quaderno: da 30/11/1950 a ≥ 18/07/1951



Cosa ha fatto Pontecorvo dal novembre 1950 al settembre 1951 e perché ha smesso di scrivere su questo primo quaderno solo pochi giorni dopo aver iniziato a lavorare al sincrotrone di Dubna?

Con l'aiuto del figlio Gil siamo riusciti a ritrovare un secondo quaderno in cui viene descritta la sua attività di ricerca durante i mesi mancanti nel primo quaderno.

La copertina riporta la data del 30 novembre 1950 mentre l'ultima data che si trova scritta sul quaderno è quella del 18 luglio 1951, nove pagine prima della fine del quaderno.

30/XI 1950

Attraverso la preziosa documentazione di questi due quaderni inediti (anche se, come vedremo, Pontecorvo molto spesso più che scriverci ci scarabocchia formule, numeri e grafici disordinati) abbiamo potuto ricostruire gli interessi scientifici, il lavoro sperimentale, la personalità e anche alcune grandi intuizioni che questo grande scienziato ha avuto nel suo primo anno e mezzo di permanenza in Russia.

... Ma procediamo con ordine ...

Chi è Bruno Pontecorvo come uomo e come scienziato quando, nel settembre del 1950 all'età di 37 anni decide di abbandonare tutto per andare a vivere in Russia?

Quali sono i motivi per cui prende questa drastica decisione? decisione che cambierà irrevocabilmente non solo tutta la sua vita ma anche quella della moglie e dei figli.

Persona che crede fermamente nel Comunismo

Bruno Pontecorvo è un comunista convinto che crede in una società socialista ispirata ad un profondo senso di giustizia e di eguaglianza:

Scrive in una nota autobiografica del 1988: *"Le mie opinioni politiche sono di sinistra. In origine esse erano dovute soprattutto al mio odio per il fascismo e, io penso ora, al senso di giustizia inculcatomi da mio padre....., opinioni dominate da una categoria non logica che io chiamo adesso "religione", una specie di "credo fanatico"....."*



Un fisico sia sperimentale che teorico con idee geniali

Nel 1934 Bethe e Peierls avevano valutato che i neutrini potevano penetrare ben 10^{+16} Km (~ 1000 anni luce) di materia solida prima di interagire, e pertanto concludevano che: *"it is therefore absolutely impossible to observe processes of this kind with neutrinos created in nuclear transformations"*.

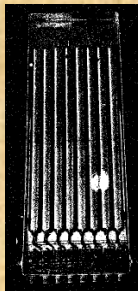
Pontecorvo nel 1945 propone di rivelare i neutrini con la reazione $\nu + {}^{37}\text{Cl}_{17} \rightarrow {}^{37}\text{Ar}_{18} + e^-$. È un'idea geniale che solo dopo molti anni verrà ripresa e messa in pratica da altri che per questo otterranno il premio Nobel.



Subito dopo l'esperimento di Conversi, Pancini e Piccioni e l'interpretazione datane da Fermi, Teller and Weisskopf, Pontecorvo scrive nel suo famoso articolo del '47: *"there exists fundamental analogy between β -processes and processes of emission and absorption of charged mesons"*.

Per primo concepisce l'universalità muone-elettrone nelle interazioni deboli !

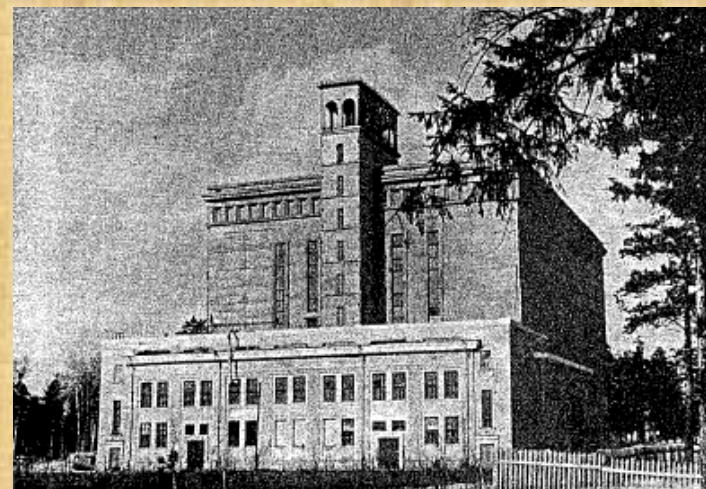
Quando nel '50 arriva in Russia ha appena pubblicato un articolo di rivista (*"Recent development in proportional counter technique"*) sui rivelatori a gas proporzionali con alta amplificazione; rivelatori non molto diversi dai moderni rivelatori a gas che si usano tutt'oggi. **Con essi ha studiato le proprietà fondamentali del μ e del suo decadimento:** la particella carica nel decadimento è un elettrone; è cinematicamente consistente con un decadimento in un elettrone e due neutrini; non si osservano decadimenti $\mu \rightarrow e + \gamma$



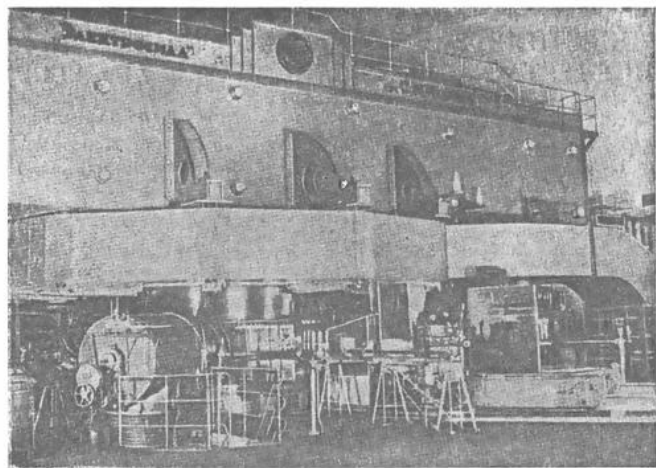
Una nuova vita e nuovi esperimenti a Dubna

Sicuramente Bruno Pontecorvo deve essere stato entusiasta di arrivare allo "Institute of Nuclear Problems" di Dubna all'inizio di Novembre 1950, dove poteva lavorare al sincrociclotrone dell'istituto che all'epoca era il più potente del mondo e per di più vivere in una società che proclamava di voler realizzare il vero comunismo.

La sua fama di geniale discepolo di Fermi lo precede e suscita grande entusiasmo tra i fisici del Laboratorio. È abitudine anche tra colleghi del laboratorio chiamarsi col nome seguito dal patronimico e risulta quindi a tutti molto imbarazzante chiamarlo semplicemente Bruno, il suo solo nome di battesimo. Il padre di Bruno si chiamava Massimo per cui decisero di chiamarlo Bruno Maximovich, nome che gli rimase per sempre.



Edificio del sincrociclotrone

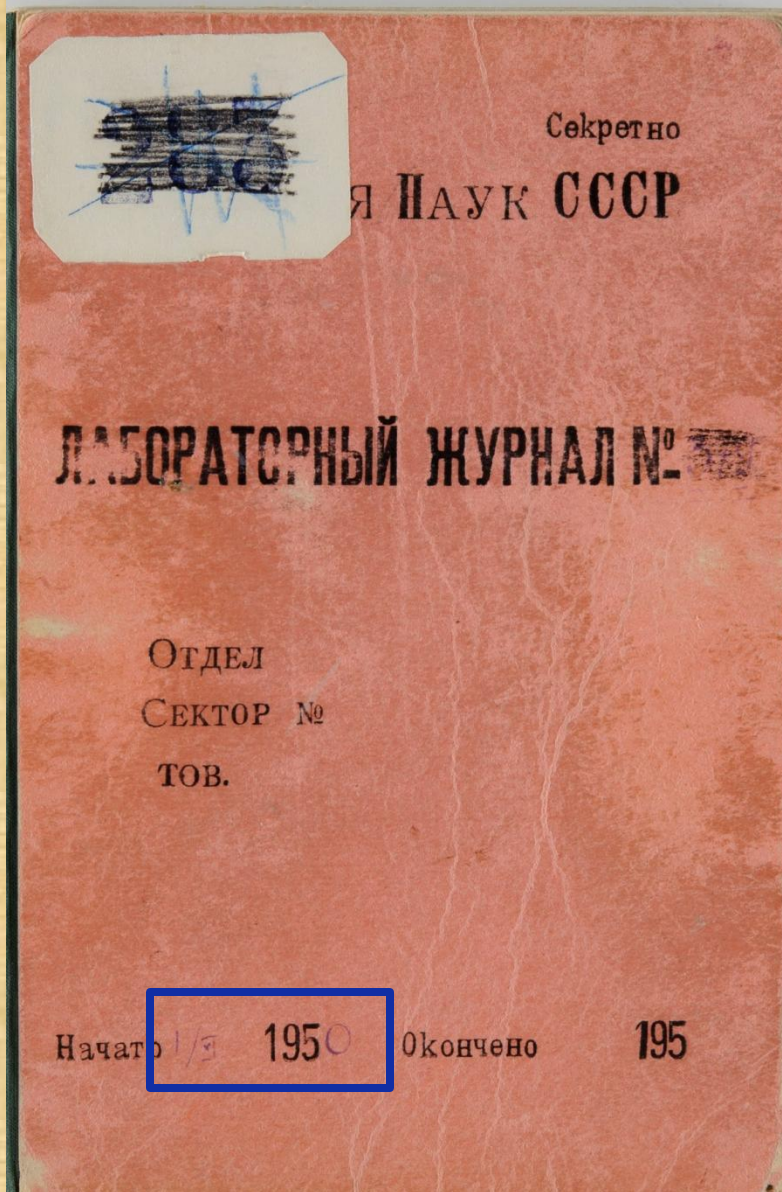


Il sincrociclotrone

	Kind of accelerated particles and their energy		
	280 MeV deuterons	560 MeV α 's	480 MeV protons
Internal target current (μA)	1	0,025	0,2–0,3
Extracted proton flux at a distance of 10 m from the magnetic channel ($cm^{-2} sec^{-1}$) . . .	—	—	$1 \cdot 10^6$ ($E_p = 460$ MeV)
Neutron flux at the maximum of the angular distribution 2 m from the internal target ($cm^{-2} sec^{-1}$)	$8 \cdot 10^7$	$2 \cdot 10^8$	$5 \cdot 10^6$
Neutron energy at the maximum of the energy distribution (MeV)	120	120	380
Halfwidth of the angular neutron distribution (radian)	0,17	0,35	0,55
Process responsible for neutron production . .	Stripping	α -particle disintegration	charge exchange

Parametri dei fasci disponibili nel 1950

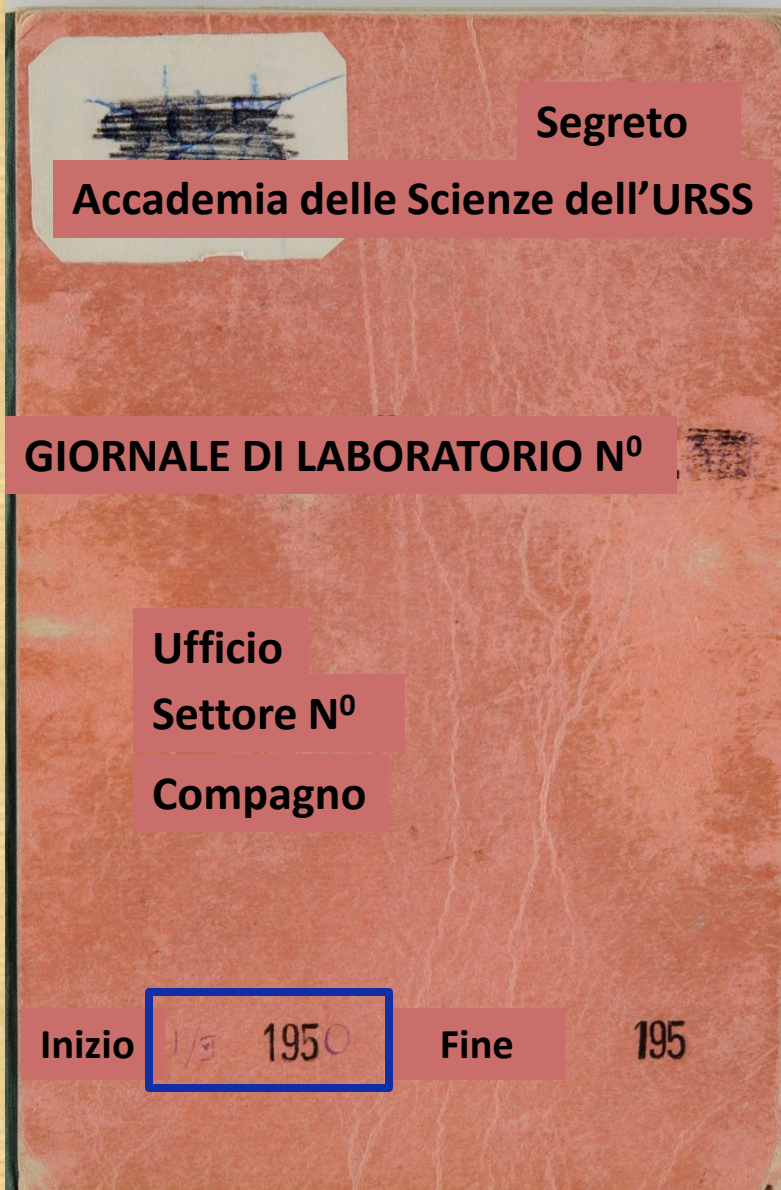
I/XI 1950 **Primo Quaderno:** da 01/11/1950 a ≤ 30/11/1950
da 14/09/1951 a ≥ 24/03/1952



1/11/1950: Pontecorvo inizia il suo lavoro di ricerca presso il sincrociclotrone dello "Institute of Nuclear Problems" di Dubna. La segretezza è rigorosa: non si può prendere appunti e scrivere formule se non su un quaderno ufficiale che viene riconsegnato a fine giornata e controllato periodicamente dalle autorità.

Questa è la foto del primo quaderno di appunti dove Pontecorvo annota giornalmente la sua attività scientifica presso il sincrociclotrone di Dubna a partire dal suo primo giorno di lavoro.

I/XI 1950 Primo Quaderno: da 01/11/1950 a ≤ 30/11/1950
da 14/09/1951 a ≥ 24/03/1952



1/11/1950: Pontecorvo inizia il suo lavoro di ricerca presso il sincrociclotrone dello "Institute of Nuclear Problems" di Dubna. La segretezza è rigorosa: non si può prendere appunti e scrivere formule se non su un quaderno ufficiale che viene riconsegnato a fine giornata e controllato periodicamente dalle autorità.

Questa è la foto del primo quaderno di appunti dove Pontecorvo annota giornalmente la sua attività scientifica presso il sincrociclotrone di Dubna a partire dal suo primo giorno di lavoro.

Pagina 1 del primo quaderno

Pontecorvo è subito interessato alla possibilità che offre il sincrotrone di produrre fasci di neutroni; immediatamente propone un metodo per stimarne la distribuzione in energia.

1 Novembre (1950)

- Neutron production by cyclotron particles -

120 Mordone

- Neutron production by cyclotron particles - 1

In the experiment with the water tank, one can get an idea of the neutron energy by measuring the space distribution of neutrons (for example measure $r^2|_{Av}$). A comparison at different energies is interesting. $r^2|_{Av}$ would be probably representative of the "evaporation" process, while the mean ~~value~~ relaxation length would be probably characteristic of the "see knock on" process.

"In the experiment with the water tank, one can get an idea of the neutron energy by measuring the space distribution of neutrons (for example measure $r^2|_{Av}$)."

(I neutroni sono prodotti a partire da un fascio di particelle α accelerate a 560 MeV e fatte collidere su una targhetta interna. Pertanto la loro distribuzione in energia non è conosciuta).

- Fission from highly excited states -
 The normal fission happens usually from low excited states (≈ 10 MeV), with high energy bombardment. Now, as the fission of medium A shows, there must be fissions ~~about~~ arising from very highly excited states, in very few cases. These fissions from highly excited states must release plenty of energy, ~~the~~ in U or Th. The difficulty in detecting them is "electrical" noise. This is stated to be $\sim 1/\text{min}$. It is possible to reduce it by gas amplification.

2

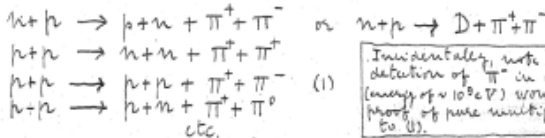
H⁴ problem - Is it possible to detect the H⁴ particles inside the chamber? One could use the magnetic field of the cyclotron to curve the electrons.

3 November

According to Anatoly Alexandrovich, the experiment with H⁴ is possible "inside the tank", with an arrangement of 3 counters in coincidence.

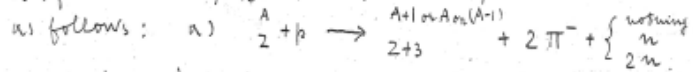
Multiple meson production

The threshold for multiple ^(double) production, for example:



Incidentally, note that the detection of π^- in alpha count (energy of $\sim 10^8$ eV) would be a proof of pure multiple, according to U.S.

is ~ 600 MeV in H. But in heavy materials the threshold is of the order of 300 MeV! An experiment can be done as follows:



I irradiate a target, and separate chemically the element Z+3 - Let us evaluate the σ for the emission of $2\pi^-$. It is:

(Heavy element) $\frac{\sigma_{(2\pi^-)}}{\sigma_{(2\pi^-)}} \approx \frac{\sigma_{(2\pi^-)}^{(500A)}}{\sigma_{(2\pi^-)}^{(100A)}} \times \frac{\sigma_{(2\pi^-)}^{(300A)}}{\sigma_{(2\pi^-)}^{(100A)}} = 10^{-26} \times \frac{10^{-28}}{10^{-24}} \approx 10^{-30} \text{ cm}^2$

In Pb, this gives a mean free path for double π^- equal to:

$l = \frac{200}{12 \times 0.6 \times 10^{24}} \text{ cm} = 3 \times 10^7 \text{ cm}$

Pontecorvo scrive nel quaderno le sue idee su quali siano gli esperimenti interessanti da fare con questo acceleratore e quali siano le tecniche e i rivelatori più idonei per realizzarli

- Fission from highly excited states -
The difficulty in detecting them is "electrical" noise. This is stated to be $\sim 1/\text{min}$. It is possible to reduce it by gas amplification

H⁴ problem - Is it possible to detect the H⁴ particles inside the chamber? One could use the magnetic field of the cyclotron to curve the electrons.
 3th November

According to Anatoly Alexandrovich, the experiment with H⁴ is possible "inside the tank", with an arrangement of 3 counters in coincidence.

Multiple meson production
 The threshold for multiple (double) production, for example:

$n+p \rightarrow p+n + \pi^+ + \pi^-$ or $n+p \rightarrow D + \pi^+ + \pi^-$
 $p+p \rightarrow n+n + \pi^+ + \pi^+$
 $p+p \rightarrow p+p + \pi^+ + \pi^-$
 $p+p \rightarrow p+n + \pi^+ + \pi^0$
 etc.

is ~ 600 MeV in H. But in heavy material the threshold is of the order of 300 MeV. An experiment can be done as follows:

A pagina 9 Pontecorvo scrive un ultimo breve commento "On the multiple production of mesons". Nella restante parte della pagina c'è scritto, in senso capovolto, il un draft di un articolo.

- On the multiple production of mesons -

In discussing the phenomenon of multiple production, from an experimental point of view, it is necessary to remember the possibility that an appearance of multiple production may be given by the production of heavy mesons (spin integer, strong interaction with matter), which of course, decay into π mesons immediately, giving the appearance of multiple production, while, in fact there maybe only one particle produced per hit.

- On the multiple production of mesons -

In discussing the phenomenon of multiple production, from an experimental point of view, it is necessary to remember the possibility that an appearance of multiple production may be given by the production of heavy mesons (spin integer, strong interaction with matter), which of course decay into π mesons immediately, giving the appearance of multiple production, while, in fact there maybe only one particle produced per hit.

... with a compensating filter of Al (2.5cm) in front of the collimator, equivalent (2.5cm) in... This method is preferable for small angle of detect(ion) to the(?) method.

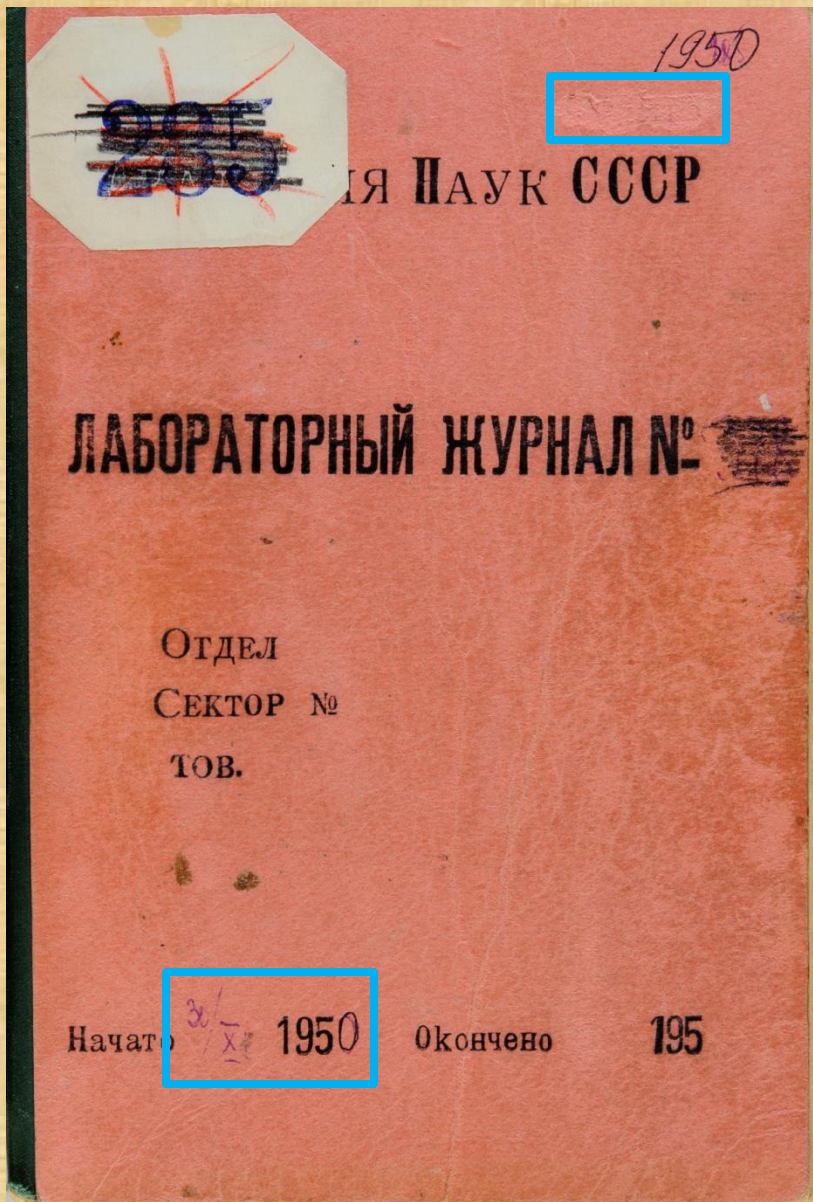
Dopo le prime 9 pagine Pontecorvo smette di scrivere su questo quaderno e riprende a scriverci il 14 settembre del 1951 girando il quaderno e iniziando dall'ultima pagina.

Nelle pagine successive, e fino alla fine del quaderno senza ulteriori interruzioni, descrive l'attività scientifica sua e del suo gruppo riportando giornalmente i progressi di questo e di altri successivi esperimenti, annotando i conteggi della presa dati e delle misure effettuate, commentando i risultati delle analisi dei dati raccolti e scrivendo infine le bozze dei relativi articoli che verranno poi pubblicati in russo come report interni del laboratorio.

L'ultima data che si trova sul quaderno è quella del 24 marzo 1952, poche pagine prima che termini il quaderno completando di scrivere questa pagina, la numero 9, che come abbiamo detto, aveva parzialmente scritto nel novembre del 1950.

Pontecorvo sta qui finendo di scrivere il draft dell'esperimento Production of neutral mesons by neutrons. Queste ultime righe riguardano le misure del flusso del fascio di neutroni.

30/XI 1950 Secondo quaderno: da 30/11/1950 a \geq 18/07/1951



In questo secondo quaderno, sulla cui copertina è riportata la data del 30 Novembre 1950, Pontecorvo continua ad annotare idee, calcoli, misure e appunti di varia natura su possibili esperimenti da fare con il sincrociclotrone dello "Institute of Nuclear Problems" di Dubna.

Per qualche strano motivo Pontecorvo dopo aver scritto solo nove pagine di appunti sul primo quaderno, alla fine di Novembre del 1950 inizia a scrivere su un nuovo quaderno identico al primo in cui approfondisce e comincia a mettere in pratica alcune delle idee a cui aveva accennato nelle prime pagine del primo quaderno. Si noti che la parola Секретно (Segreto) in alto a destra è stata cancellata; probabilmente è avvenuto successivamente quando queste ricerche sono diventate non più segrete o forse (ma questa è solo una mia fantasia) è stato lo stesso Pontecorvo a farlo ed è proprio questo il motivo per cui cambia quaderno quasi subito: non vuole un quaderno ufficiale ma solo un brogliaccio.

Pagina ① del secondo quaderno

$\sigma \propto \pi R^2$ $R \propto A^{1/3}$ $\sigma \propto A^{2/3}$ $A^{2/3}$ Heavy elements ①
 $\frac{1}{2}$ $z=1$ $\sigma \propto A$ A Light elements

Collision \Rightarrow Meson
 Meson collision $N+N$ $\left\{ \begin{array}{l} \text{Collision } \sigma_c \\ \text{Meson collision} \\ \text{scattering} \end{array} \right.$
 $M+N$ $\left\{ \begin{array}{l} \text{Disintegration} \end{array} \right.$

$\mu\mu$ 10^{-28} cm^2

$\pi^+ + \left\{ \begin{array}{l} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \right.$ σ_{π} nuclear High energies
 $\mu\mu$ 100 MEV
 $N+N \rightarrow \left\{ \begin{array}{l} P+\pi^- \\ \text{---} \end{array} \right.$ $\left\{ \begin{array}{l} \pi^+ + P \rightarrow N + \pi_0 \\ \pi^+ + P \rightarrow N + 2\pi^+ \text{ Cascade} \\ \pi^+ + N \rightarrow P + \pi_0 \\ \text{scattering processes } \sigma_{\pi} \\ P + \pi^- \rightarrow N + N \end{array} \right.$

$\sigma = \pi R^2 \propto A^{2/3}$

σ_{tot} (mb)	E (MeV)
4	1
1	10
0.07	100
0.03	300

$\lambda \ll R$ $(\mu\mu)^{10^{-15}}$ πR^2

3×10^{-26} 10^{-28}

$\xi = \frac{E-200}{200}$ $\frac{1}{2} \frac{M^2}{2M} v^2$

$\sigma \propto A^{2/3}$ $A^{2/3}$ Heavy elements
 $\sigma \propto A$ A Light elements

Pontecorvo inizia il quaderno valutando con considerazioni di tipo geometrico l'ordine di grandezza delle sezioni d'urto totale e di produzione di mesoni nell'urto nucleone-nucleone e nucleone-nuclei.

$\sigma = \pi R^2 \propto A^{2/3}$

$\lambda \ll R$

Pagina (2) del secondo quaderno

1) $\sigma \ll \pi \left(\frac{\hbar}{mc}\right)^2$ $\sigma \ll 6 \times 10^{-26} \text{ cm}^2$ nucleon, nucleon, 2

when $\lambda \ll \frac{\hbar}{mc}$

2) In a collision of parameter b the predominant v will be that such that $2\pi v \times b = u$ (relative velocity). Then the collision will take place with great probability only if b is so small, that v corresponds to the relative energy E . $\frac{\hbar v}{2\pi} = E$

~~$b \leq \frac{u \hbar}{2\pi E}$~~
 ~~$\sigma \leq \pi \left(\frac{u \hbar}{2\pi E}\right)^2 = \frac{u^2 \hbar^2}{4\pi^2 E^2} = \frac{u^2 \hbar^2}{4\pi^2 \frac{1}{2} m u^2} = \frac{2 \hbar^2}{\pi m}$~~

$b \leq \frac{u \hbar}{2\pi E} = \frac{2 \hbar}{\pi m}$

$E = \frac{1}{2} m u^2$
 $\frac{u}{E} = \frac{2}{m u}$

$\sigma \leq 4\pi \lambda^2 \ll \frac{4\pi \hbar^2}{2mE}$ $\sigma_{p > 300} \rightarrow 10^{-28}$ (D.K.)

3.50 MeV
 10^{-27}

Pontecorvo continua per alcune pagine a valutare con considerazioni di tipo geometrico l'ordine di grandezza delle sezioni d'urto totale e di produzione di mesoni nell'urto nucleone-nucleone e nucleone-nuclei.

1) $\sigma \ll \pi \hbar^2 / (mc)^2$ $\sigma \ll 6 \times 10^{-26} \text{ cm}^2$
 nucleon, nucleon,
 when $\lambda \ll \hbar / mc$

2) In a collision of parameter b the predominant v will be that such that $2\pi v \times b = u$ (relative velocity) Then the collision will take place with great probability only(?) if b is so small that v corresponds to the relative energy $\hbar v / 2\pi = E$ (?! $\hbar v / 2\pi = E \rightarrow v = E / 2\pi \hbar$)

$b \leq u / 2\pi v = u / 2\pi (2\pi \hbar / E) = 2\hbar / mu = 2\hbar / p = 2 \lambda$
 $\{E = 1/2 mu^2 \rightarrow u/E = 2/mu = 2/p\}$

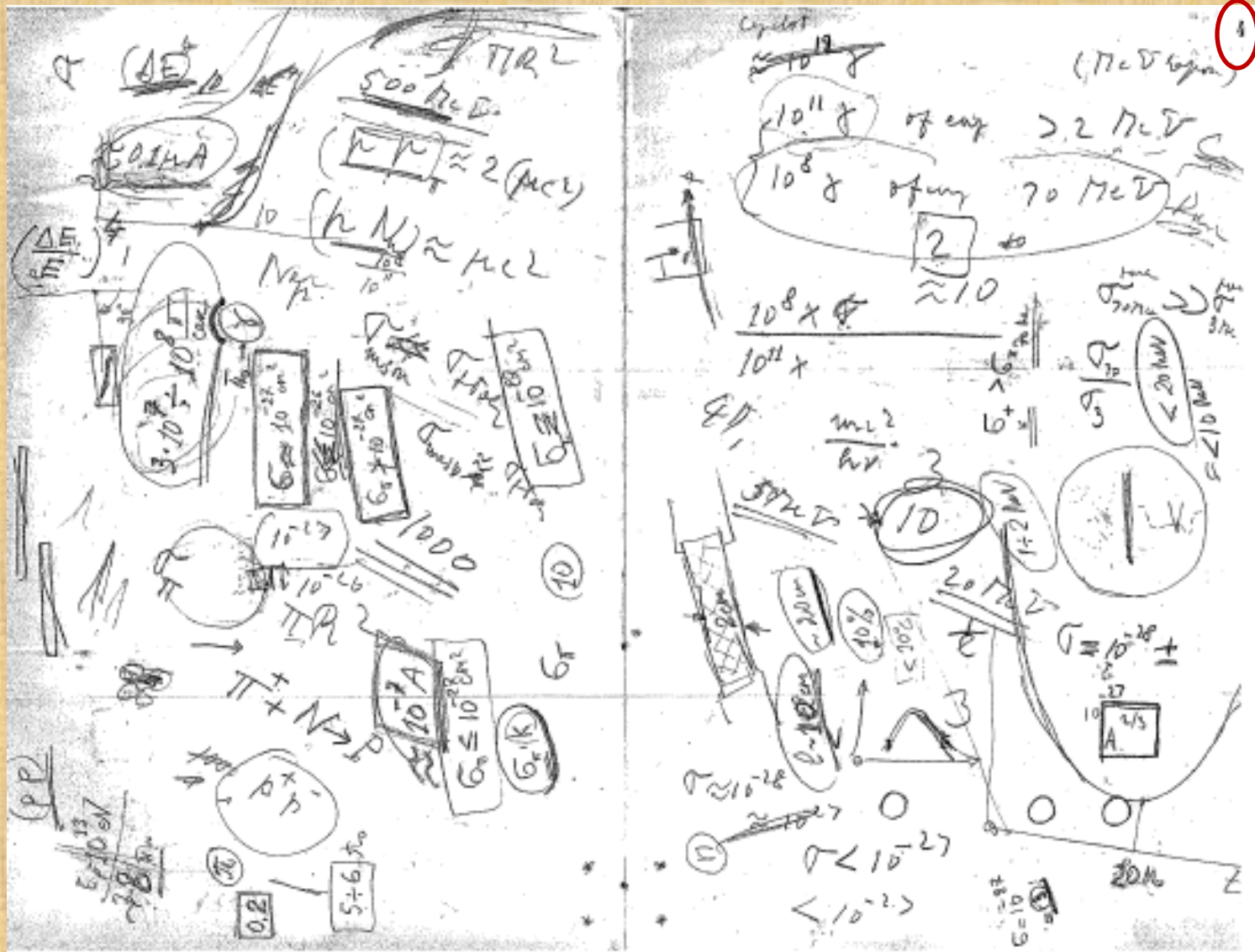
$\sigma \leq \pi b^2 \leq 4\pi \lambda^2 = 4\pi \hbar^2 / (mu)^2 \rightarrow$
 $\sigma \ll 4\pi \hbar^2 / (2mE)$

3.50 MeV
 10^{-27}

$\sigma_{p > 300} \rightarrow 10^{-28}$ (D.K.)

Pagina 4 del secondo quaderno

Pontecorvo molto spesso usa le pagine di questi quaderni solo come un "brogliaccio" scarabocchiandovi in modo disordinato e senza alcun ordine logico numeri, formule e grafici, forse solo per fissare i pensieri e le idee che gli turbinano nella mente. Risulta pertanto spesso difficile capire con esattezza, gli esperimenti e le misure che Bruno e il suo gruppo vogliono fare.



Page 33-36: Draft "Measurement with a "star detector" of total cross sections for neutrons produced in the bombardment of Be with 475 400 Megavolt protons -"

~~Measurement with a "star detector" of total cross sections for neutrons produced in the bombardment of Be with 475 400 MeV Megavolt protons -~~ 33

~~for several nuclei~~

Introduction - Total cross sections ~~of~~ ^{of} ~~several~~ ^{of} ~~nuclei~~ ^{nuclei} have been ~~measured~~ ^{measured} for energies of ~~about 90 MeV~~ ^{~ 40, 90 MeV () and ~ 280 MeV ()} for ~~energies of ~ 40 MeV~~ ^{for H and C are also available measurements at 40 MeV () and 156 MeV ()} and ~~280 MeV ()~~ ^{and 156 MeV ()}.

Total cross sections are ~~usually~~ ^{usually} best measured by an attenuation method, in "good geometry conditions", ~~that is~~ ^{with a great distance between the attenuator and the detector, in essence, one of the conditions} with a great distance between the ~~attenuator~~ ^{attenuator} and the detector, ~~in essence, one of the conditions~~.

Until now the following detectors have been used:

- 1) Radioactive indicators, such as C^{11} , (threshold ≈ 20 MeV) () produced in a n-2n reaction (threshold \approx
- 2) Bi fission chamber, which ~~has the threshold ≈ 50 MeV~~ ()
- 3) A telescope counting high energy recoils protons from a radiator placed in the beam. The telescope ~~consists of proportional counter~~ ^{consists of proportional counter} and a scintillating counter in coincidence with the threshold of the detector.

Measurement with a "star detector" of total cross sections for neutrons produced in the bombardment of Be with 400 Megavolt protons

Introduction - Total cross section of several nuclei have been measured for energies of ~ 40 , 90 MeV () and ~ 280 MeV (). For H and C are also available measurements at 40 MeV () and 156 MeV ().

Total cross sections are ~~usually~~ ^{usually} best measured by(?) an attenuation method, in "good geometry conditions", with a great distance between the attenuator and the detector.

Until now the following detectors have been used:

- 1) Radioactive indicators, such as C^{11} , (threshold ~ 20 MeV) () produced in a n-2n reaction (threshold \approx
- 2) Bi fission chamber (threshold ≈ 50 MeV) ()
- 3) A telescope counting high energy recoils protons from a radiator placed in the beam () The telescope consists of proportional and a scintillating counters in coincidence (?) the threshold of the detector is

Pagine 33-36: Draft "Measurement with a "star detector" of total cross sections for neutrons produced in the bombardment of Be with 475 400 Megavolt protons -"

insured by placing absorbers between the counters of telescope. ()

In the course of a separate experiment on the production of secondary neutrons produced by the high energy neutron beam of ~~the proton~~ obtained under 475 MeV bombardment of the proton, we realized ~~that~~ the interest of using ~~atypical~~ detectors of high energy neutrons. The secondary neutrons of measuring one production of ~~heavy~~ neutrons in a heavy element.

~~This is not a typical detector based on this principle - the "star" detector - will be described in the second part of this article. We will describe measurements of the total cross sections~~

In this article we will ^{first} describe a detector based on this principle, ^{which we shall call} the "star" detector - and ^{then} report of measurements of total cross section of various elements for neutrons of energy ϵ , which were obtained ~~by~~ means of one star detector.

The star detector as indicator of ~~high energy~~ ^{high energy} particles, ^{as observed in photograph emulsions of the counter} give in principle a good method of studying high energy particles. However the star detector

insured by placing absorbers between the counters of telescope. ()

In the course of a separate experiment on the production of secondary neutrons produced by the high energy neutron beam obtained under 475 MeV bombardment(??) of the proton, we realized the interest of detecting high energy neutrons by measuring one secondary production of neutrons in a heavy element.

In this article we will first describe a detector based on this principle - which we shall call the "star detector"- and then report the measurements of total cross section of various elements for neutrons of energy , which were obtained by means of the star detector.

The stars, as observed in photograph emulsions are typical big energy events, and give in principle a good method of detecting high energy particles. However the star detector

Page 33-36: Draft "Measurement with a "star detector" of total cross sections for neutrons produced in the bombardment of Be with 475 400 Megavolt protons -"

based on photographs would be because of the long time necessary for the analysis of the measurements. An alternative star detector could consist of one ionizing chamber detecting the ionization of charged particles, mainly protons and α 's released in the evaporation process. Such a chamber could be used to register only pulses bigger than a certain value. However it is easy to see that a star ionization chamber has a very small sensitivity, unless very high pressures of a heavy gas like Xe are used. A different solution of the process is to detect neutrons (order of magnitude of few MeV) emitted in the evaporation process. This method has the advantage that a very high thickness of star producing material (order of a mean free path of the high energy neutrons) can be used. It is true that the secondary neutrons will be detected with a small efficiency ($\approx 10^{-4}$); however the high multiplicity of neutron production per inelastic collision (≈ 10 in a heavy element) and the very large thickness which can be used make the overall sensitivity high. For example, a neutron beam of energy ≈ 400 MeV, and intensity 3×10^3 / sec / cm² gave

based on photographs would be because of the long time necessary for the analysis of the measurements. An alternative star detector could consist of one ionizing chamber detecting the ionization of charged particles, mainly protons and α 's released in the evaporation process. Such a chamber could be used to register only pulses bigger than a certain value. However it is easy to see that a star ionization chamber in practice has a very small sensitivity, unless very high pressures of a heavy gas like Xe are used. A different solution of the process is to detect neutrons (order of energy (?) of few MeV) emitted in the evaporation process. This method has the advantage that a very high thickness of star producing material (order of a mean free path of the high energy neutrons) can be used. It is true that the secondary neutrons will be detected with a small efficiency ($\approx 10^{-4}$); however the high multiplicity of neutron production per inelastic collision (≈ 10 in a heavy element) and the very large thickness which can be used make the overall sensitivity high. For example, a neutron beam of energy ≈ 400 MeV, and intensity 3×10^3 / sec / cm² gave

Pagine 33-36: Draft "Measurement with a "star detector" of total cross sections for neutrons produced in the bombardment of Be with 475 400 Megavolt protons -"

in our detector (Pb + BR in paraffine block)
~~3600 pulses/min~~ ²⁵⁰⁰ pulses/min.

The ~~secondary~~ ^{secondary} neutrons were detected in an arrangement consisting of a BF₃ counter ^{long on diameter} imbedded in a paraffine block ^{as in fig 1}. The characteristic of this ^{long on diameter} counter is ^{not} ~~the~~ ^{the} ~~plot~~ ^{plot} ~~response~~ ^{response} of such detector in ^{the} ~~the~~ ^{the} ~~energy~~ ^{energy} region ^{150 KV} ~~10 KV~~ ^{10 MeV} ~~10 MeV~~. This important feature, ^{however,} ~~is~~ ^{is} ~~not~~ ^{not} ~~important~~ ^{important} in the experiment on total cross section ^{although} ~~it~~ ^{it} ~~may~~ ^{may} ~~be~~ ^{be} ~~very~~ ^{very} ~~useful~~ ^{useful} in other applications of the "star detector".

Another ^{advantage} ~~advantage~~ of the detector, for use ^{with} ~~in~~ ^{with} ~~pulsed~~ ^{pulsed} sources, in particular as a monitor, is the fact that the ~~process~~ ^{process} of diffusion of the thermal neutrons, which are detected in the counter, smooth out the distribution of intensity of the ^{fast} ~~fast~~ ^{fast} neutrons during the cyclotron proton pulse (the proton pulse distribution is known to have peak and valleys). As a consequence of this fact it is possible to count

in our detector (Pb + BF₃ in paraffine block) gave 2500 pulses/min. The secondary neutrons were detected in the arrangement, usually called "long counter" consisting of a BF₃ counter long, and Diameter imbedded in a paraffine blok. The characteristic of this long counter in paraffine is that its neutron sensitivity is not very sensitive of the energy in Fn (?) region 50 KeV- 10 MeV. While this feature is not very important on the experiment on total cross section, it may be very useful in other application of the "star detector". Another advantage of the detector, for use with pulsed sources, in particular as a monitor, is the fact that the process of diffusion of the thermal neutrons, which are detected in the counter, smooth out the distribution of intensity of the (?) neutrons during the cyclotron proton pulse (the proton current pulse distribution is known to have peak and valleys) As a consequences of this fact it is possible to count

I contatori per neutroni BF³ (pag. 17)

March 10, 1951

BF³ counters

No. 1, brass, 30 cm pressure BF³, contains effectively:
 20.8×10^{19} atoms of B¹⁰ in total.

No. 2, glass, 10 cm pressure BF³, contains effectively:
 3.4×10^{19} atoms of B¹⁰ in total.

kompare counter

glass
 $2 \times 64 + 44$

brass
 $32 \times 64 + 60$

$3 \times 64 + 28$

$29 \times 64 + 4$

90

March 10, 1951

BF³ counters

N. 1, brass, 30 cm pressure BF³, contains effectively:
 20.8×10^{19} atoms of B¹⁰ in total

N. 2, glass, 10 cm pressure BF³, contains effectively:
 3.4×10^{19} atoms of B¹⁰ in total

kompare counter

glass
 $2 \times 64 + 44$
 $3 \times 64 + 28$

brass
 $32 \times 64 + 60$
 $29 \times 64 + 4$

Page 33-36: Draft "Measurement with a "star detector" of total cross sections for neutrons produced in the bombardment of Be with 475 400 Megavolt protons -"

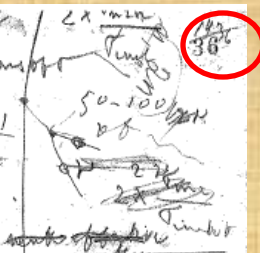
The cube does not appear one way or the other when the ~~cube~~ beam is absent. We want to discuss now the energy dependence of the star detector.

- Energy response of the star detector

The ~~dependence~~ sensitivity of the "star detector" to the neutron event is roughly the number of secondary neutrons per collision, as a function of the energy.

between 100-400 MeV is not known. The value(?) of ν in the region 100-400 MeV is not known. The detector consequently has not a real threshold, if the secondary neutrons are detected, as in our work, incoherently. (if the secondary neutrons, however, were detected in coincidence, it would be possible to impose a high threshold of sensitivity.) For this reason the absence of

(see page 50)



the cube does not affect the reading when the collimator of the beam is absent. These points are important for the determination of σ . We want to discuss now the energy dependence of the star detector.

- Energy response of the star detector

The sensitivity of the "star detector" to the neutron energy is given by roughly the number of secondary neutrons per collision, as function of the energy. The value(?) of ν in the region 100-400 MeV is not known. The detector consequently has not a real threshold, if the secondary neutrons are detected, as in our work, incoherently. (if the secondary neutrons, however, were detected in coincidence, it would be possible to impose a high threshold, at loss of sensitivity) Because of the absence of (see page 50)

Page 50:
a threshold, it was necessary to ensure that the presence of low energy neutrons ($\ll \approx 100$ MeV) did not seriously affect the measurement. This was done...etc

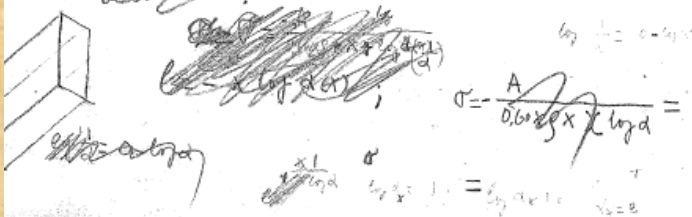
Pagine 50-52: Draft "Measurement with a "star detector" of total cross sections for neutrons produced in the bombardment of Be with 475 400 Megavolt protons -"

detector. The same conclusion ^{was} reached
 of measuring the attenuation ^{in Cu} (for $X_{Cu} \cong$ one mean
 free path) with and without a filter
 of paraffine ^{with thickness \cong one mean free path}.
 The fact that the attenuation in Cu ^{with} and
 without filter, was equal with the quite
 small statistical error, ^{is another} information of the
 stable import of neutrons $<$ in our
 measurements.

Results

Once ^{we have} established that the absorption
 curve were exponential, ^{we} performed
 all our measurements with thickness
 of absorber approximately equal to the
 m f p. This is preferable for the
 following reasons:

- 1) ~~The cross section, ^{is given by} $\sigma = \frac{A}{0.693 \times X \times \log d}$~~
~~knowing the cross section, ^{is given by} $\sigma = \frac{A}{0.693 \times X \times \log d}$~~
~~relation:~~



detector. The same conclusion can be reached
 by measuring the attenuation in Cu (for $X_{Cu} \cong$
 one mean free path) with and without a filter
 of paraffine ... long: such paraffine thickness
 is the m f path in paraffine for neutrons of
 The fact that the attenuation in Cu, with and
 without filter, was equal with the quite small
 statistical error, is another information of the
 stable (?) import (?) of neutrons $<$ in our
 measurements.

Results

Once we **established that the absorption curve
 were exponential**, we performed all our
 measurements with thickness of absorber
 approximately equal to the m. f. p. This is
 preferable for the following reasons:

- 1) ~~The cross section is given by the relation:~~

~~$\sigma = \frac{A}{0.693 \times X \times \log d}$~~

Measurements in Cu of "Total σ with star detector"

(pag.22)

Results

(Normalised)

Total σ with star detector Results (Normalised)		
Cu 60cm + Pb cube	0.0405 ± 0.0012	
Cu 60cm, no Pb cube	0.0389 ± 0.0012	
No Cu, no Pb cube	0.0438 ± 0.006	
No Cu, Pb cube	0.542 ± 0.010	
Cu 10cm, Pb cube	0.214 ± 0.004	
Cu 20cm, Pb cube	0.106 ± 0.002	
Cu 30cm Pb cube	0.0627 ± 0.0012	
No Cu Pb cube	0.543 ± 0.010	
Cu 15cm Pb cube	0.158 ± 0.002	
Cu 5cm Pb cube	0.405 ± 0.006 0.360 ± 0.006	✓
Cu 25cm Pb cube	0.0845 ± 0.002	✓
Cu 10cm Pb cube	0.218 ± 0.004	✓
No Cu, Pb cube	0.544 ± 0.010	22
No Cu, No Pb cube	0.0469 ± 0.0018	0
Cu 10cm, No Pb cube	0.0413 ± 0.0012	0
Cu 20cm, No Pb cube	0.0413 ± 0.0012	0
<p>~~~~~ \downarrow Pb position is not exactly in the fore, because keeping diameter. However as close as possible it was put</p>		
No Cu, Pb cube (no paraffin)	0.555 ± 0.008	
No Cu, Pb cube paraffin, 36cm before collimator	0.352 ± 0.005	✓
Cu 10cm, Pb cube, paraffin paraffin 36cm before Pb cube	0.156 ± 0.003	✓
Cu 10cm, Pb cube (no paraffin)	0.235 ± 0.0035	
No Cu, Pb cube	0.551 ± 0.012	
~~~~~		
AC 30cm, Pb cube absorber	0.216 ± 0.003	

Risultati di misure di sezione d'urto totale su alcuni nuclei complessi (Cu,Pb,Al)
 Per ottenere le sezioni d'urto totali pensa di dover graficare in carta logaritmica
 i conteggi in funzione dello spessore della targhetta (metodo di trasmissione).
 Nel frattempo prova a scrivere in russo la richiesta di cittadinanza sovietica.

Summary of results

Total σ :

$\sigma = 1.19 \pm 0.02$

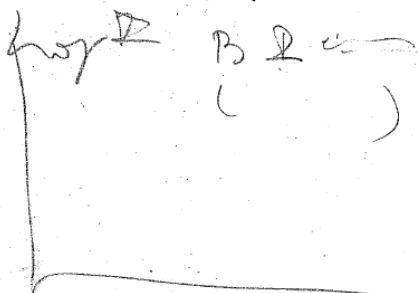
$\sigma_{Pb} = 2.89 \pm 0.07$

$\sigma_{Al} = 0.60 \pm 0.02$

Председателю Президиума Верховного
 Совета СССР
 тов. Швернику

~~Прошу Вас, т. Шверник, разрешить~~
~~подданство~~ ~~меня~~ ~~и~~ ~~мою~~ ~~семью~~
 перейти в Советское подданство.

Прошу Вас, т. Шверник, разрешить мне и моей семье
 перейти в Советское подданство.



then

Summary of results

Total σ :

$\sigma_{Cu} = 1.19 \pm 0.02$

$\sigma_{Pb} = 2.89 \pm 0.07$

$\sigma_{Al} = 0.60 \pm 0.02$

-- Presidente del Praesidium del
 Soviet supremo dell'URSS
 Compagno Shvernik (товарищу швернику)

Vi chiedo, compagno Shvernik, di permettere a
 me e alla mia famiglia di diventare cittadino
 sovietico.

(nella parte cancellata c'è scritto male in russo la stessa cosa)

↑ Log I

→ Th Cu

Questo plot su carta logaritmica si trova inserito alla pagina delle misure precedenti (pag.22)

Attenuation curve in Cu

Attenuation curve in Cu

~~** without paraffine filter~~

* Point without paraffine filter

o Point with paraffine filter

~~X normalized to (The intensity is normalized ?? so that the intensity without absorber ?? X 3~~

Attenuation curve in Cu

~~** without paraffine filter~~
~~* Point without paraffine filter~~

o Point with paraffine filter X normalized to (The intensity is normalized ?? so that the intensity without absorber ?? X 3

Log Intensity

Log Intensity

1000

100

Cu σ^*

Cu σ with ~ 36 cm filter of paraffine o

Without paraffine filter

With paraffine filter ~ 40 cm

Cu σ^*

Cu σ with ~ 36 cm filter of paraffine o

Without paraffine filter

With paraffine filter ~ 40 cm

Thickness of Cu

Thickness of Cu

10

0

5

10

15

20

25

30

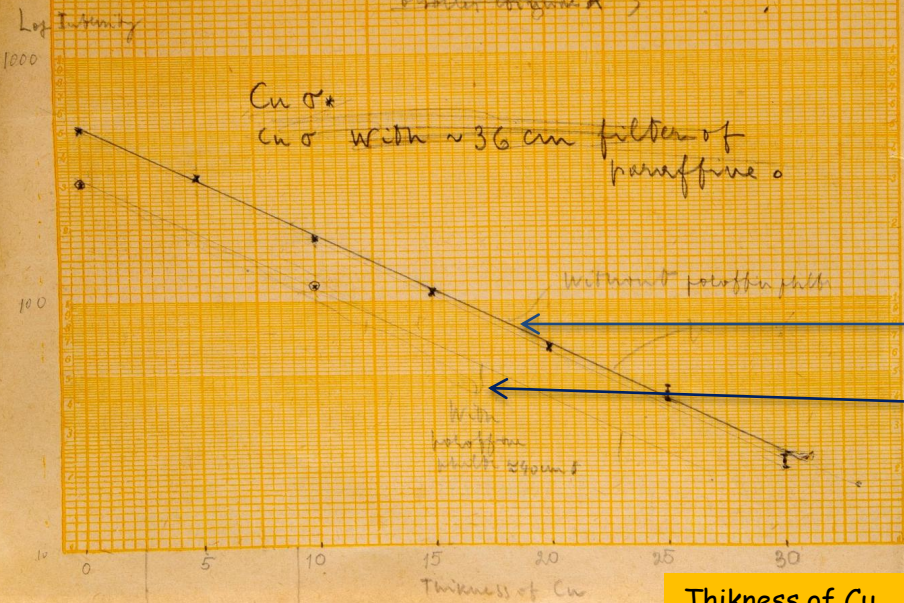
X - *момент без парафин* *аннотация и анкетам*
 O - *момент с парафин* *аннотация и анкетам с парафин*

X - *момент без парафин* *аннотация и анкетам*
 O - *момент с парафин* *аннотация и анкетам с парафин*

Attenuation curve in Cu

$*$ without paraffine filter
 \circ with paraffine filter

X normalized to (The intensity is normalized ?? so that the intensity without absorber ?? $X \cdot 3$)



Attenuation curve in Cu
 $**$ without paraffine filter
 $*$ Point without paraffine filter
 o Point with paraffine filter
 X normalized to (The intensity is normalized ?? so that the intensity without absorber ?? $X \cdot 3$)

$Cu \sigma *$
 $Cu \sigma$ with ~ 36 cm filter of paraffine o

Without paraffine filter

With paraffine filter ≈ 40 cm

Questo plot su carta logaritmica si trova inserito alla pagina delle misure precedenti (pag.22)

Le misure di sezione d'urto totale vengono estese a altri nuclei complessi

Programma Cu²⁺ - Sn, Fe, U, C, Rosoff

- I) Pb, no absorber
- II) No Pb, no absorber
- III) No Pb, 3 or 4 absorbers ~~Sn, Fe, C~~
- IV) Pb, no absorber
- V) Pb + U
- VI) Pb + Sn
- VII) Pb + Fe
- VIII) Pb + U
- IX) Pb + C
- X) Pb + Rosoff
- XI) Pb + Cu (5cm)
- XII) Pb no absorber
- XIII) Pb + Rosoff
- XIV) Pb + purple
- XV) 5 Cu
- XVI) 10cm Cu
- XVII) No Pb no absorber

Pag. 27

A - Normalized ^{results} Total σ

Pb, no absorber 0.612 ± 0.009

No Pb, no absorber 0.0413 ± 0.0016

No Pb, Sn 20cm 0.0384 ± 0.0016

No Pb, Fe 12cm 0.0397 ± 0.0020

No Pb, C 40cm 0.0405 ± 0.0020

Pb, no absorber 0.605 ± 0.009

Pb, Sn 20cm 0.166 ± 0.002

Pb, Fe 12cm 0.220 ± 0.003

Pb, U 10cm 0.154 ± 0.002

Pb, no absorber 0.597 ± 0.006

Pb, C 40cm 0.224 ± 0.003

Pb, Rosoffine 8cm 0.210 ± 0.003

Pb, no absorber 0.603 ± 0.008

Pb, Rosoffin 8cm 0.207 ± 0.003

Pb, C 40cm 0.221 ± 0.002

Pb Cu 5cm 0.360 ± 0.005

	185/195	280	156/153	270/260	150/140
H	0.034 ± 0.005	0.033 ± 0.003	0.046 ± 0.001	0.038 ± 0.002	0.039
C	0.325 ± 0.006	0.279 ± 0.004	0.230 ± 0.003	0.288 ± 0.003	0.31
Al	0.60 ± 0.02	0.566 ± 0.018		0.555 ± 0.008	
Fe	1.11 ± 0.03				
Cu	1.19 ± 0.02	1.19 ± 0.02		1.145 ± 0.0015	
Sn	2.02 ± 0.04	1.83 ± 0.03		1.87 ± 0.03	
Pb	2.84 ± 0.07	2.83 ± 0.03		2.84 ± 0.03	
U	3.40 ± 0.07	3.14 ± 0.05		3.29 ± 0.03	

Pag. 31

Pag. 29

Pontecorvo prova ad interpretare le misure di sezione d'urto totale su nuclei complessi in termini del "opaque nucleus model"

R (10 ⁻¹³ cm)	Z	$2\pi R^2$	$\frac{\sigma_{total400}}{2\pi R^2}$	$1 - \frac{\sigma_{total400}}{2\pi R^2}$	$\frac{\sigma_{total95}}{2\pi R^2}$	$1 - \frac{\sigma_{total95}}{2\pi R^2}$	$\frac{\lambda^{400}}{\lambda^{90}}$	λ
3.5	C	0.62	0.51	0.49	0.80	0.20	6.2	
4.1	Al	1.06	0.57	0.43	0.94	0.06	6.4	
5.5	Cu	1.90	0.65	0.35			6.4	
6.7	Sn	2.80	0.72	0.28			6.4	
8.1	Pb	4.10	0.70	0.30			8.1	
8.5	U	4.50	0.75	0.25			7.0	

Start all over again:

R , opaque:

$$R = (1.3 + 1.37 A^{1/3}) \times 10^{-13} \text{ cm}$$

Probability of escape when the neutron hits the nucleus = $1 - \frac{\sigma_{inel}}{\pi R^2} \approx 1 - \frac{\sigma_{total}}{2\pi R^2} =$

$$= \int_0^R \frac{2\pi x}{\pi R^2} e^{-\frac{2\sqrt{R^2-x^2}}{\lambda}} dx = \text{(where } \lambda \text{ is the mean free path in nuclear matter)}$$

$$= \left(\frac{\lambda^2}{2R^2}\right) \left[1 - \left(1 + \frac{2R}{\lambda}\right) e^{-\frac{2R}{\lambda}}\right] = \text{(let us call } R/\lambda = a)$$

$$= \frac{1}{2a^2} \left[1 - (1 + 2a) e^{-2a}\right]$$

$$R_1 = R + 1.3 \times 10^{-13} = (1.3 + 1.37 A^{1/3}) \times 10^{-13} \text{ cm}$$

$$R (10^{-13} \text{ cm}), Z, 2\pi R^2, \sigma_{total400}/2\pi R^2, 1 - \sigma_{total400}/2\pi R^2, \sigma_{total95}/2\pi R^2, 1 - \sigma_{total95}/2\pi R^2, \lambda^{400}/\lambda^{90}$$

Start all over again:

R , opaque:

$$R_{(1)} = (1.3 + 1.37 A^{1/3}) \times 10^{-13} \text{ cm}$$

Probability of no collision when the nucleon hits the nucleus = $1 - \sigma_{inel}/\pi R^2 \approx 1 - \sigma_{total}/2\pi R^2 =$

$$= \int_0^R \frac{2\pi x}{\pi R^2} e^{-\frac{2\sqrt{R^2-x^2}}{\lambda}} dx = \text{(where } \lambda \text{ is the}$$

mean free path in nuclear matter) ($R/\lambda = a$)
 $= 1/(2a^2)[1 - (1 + 2a)e^{-2a}]$

$$R_1 = R + 1.3 \times 10^{-13} = (1.3 + 1.37 A^{1/3}) \times 10^{-13} \text{ cm}$$

The "nuclear radius" R is defined conventionally by the empirical relation $R = (1.3 + 1.37 A^{1/3}) \times 10^{-13}$, which is the best fit to data of total cross section for neutrons of $\sim 15 \text{ MeV}$ (opaque nucleus model)

The "nuclear radius" R is defined conventionally by the empirical relation $R = (1.3 + 1.37 A^{1/3}) \times 10^{-13}$, which is the best fit to data of total cross section for neutrons of $\sim 15 \text{ MeV}$ (opaque nucleus model)

Pontecorvo prova ad interpretare le misure di sezione d'urto totale su nuclei complessi

Element	R_1 (10^{-13} cm)	$2\pi R_1^2$ (10^{-24})	$\frac{\sigma_{total400}}{2\pi R_1^2}$	$\frac{\sigma_{total95}}{2\pi R_1^2}$	$1 - \frac{\sigma_{400}}{2\pi R_1^2}$	$1 - \frac{\sigma_{95}}{2\pi R_1^2}$	λ^{400} (10^{-13})	λ^{90}	R_1
C	3.15	1.2	0.27	0.42	0.73	0.58	12	8	4.4
Al	4.1	1.8	0.33	0.55	0.67	0.45	13	7	5.4
Fe	5.5	2.9	0.41	0.69	0.59	0.31	14		6.8
Cu	6.7	4.0	0.50	0.79	0.50	0.21			8
Pb	8.1	5.6	0.52	0.78	0.48	0.22			9.4
U	8.5	6.6	0.56	0.80	0.44	0.20			9.8

Conclusion



Z does not work -

$$3 \times 10^{-26} / 2 \times \sqrt{2} =$$

$$\frac{\pi \times 1.3^2 \times 10^{-26}}{2} = 5 \times 10^{-26}$$

What is the evidence that the n-p σ increases at high energies again?

$$0.27 \times (1.3 + 1.37 \sqrt{56}) = 1.11$$

$$1.11 \times 2.65 = 2.94$$

$$2.94 - 1.11 = 1.83$$

$$1.83 \times 6.28 \times 6.5^2 = 1.11$$

$$1.3 + 1.37 \times 3.8 = 6.5$$

Table

Element, R (10^{-13} cm), $2\pi R_1^2$ (10^{-24}), $\sigma_{total400}/2\pi R_1^2$, $\sigma_{total95}/2\pi R_1^2$, $1 - \sigma_{400}/2\pi R_1^2$, $1 - \sigma_{95}/2\pi R_1^2$, λ^{400} (10^{-13}), λ^{90} , R_1

Conclusion

$$A^{2/3}$$

Z does not work -

What is the evidence that the n-p σ increases at high energy again?

Page 39-47: Draft "Determination of the relative yield of evaporation neutrons from various elements bombarded by ~ 350 MeV neutrons."

Determination of the relative yield of evaporation neutrons from various elements bombarded by ~ 350 MeV neutrons

Introduction

The study of nuclear disintegrations produced by high energy particles have been usually investigated with photographic plates. While this technique is extremely powerful, there are some features of the disintegration process which cannot be definitely studied in photographic plates. These are:

- 1) The dependence of the star properties upon the atomic number of the bombarded nuclei exploding nucleus
- 2) The neutron emission in the explosion.

The aim of this work was to investigate the production of neutrons in various materials bombarded by the relative number of neutrons emitted per inelastic collision (per star) in neutrons of ≈ 360 MeV. Measurements of the number of neutrons were made with the calorimeter (star) was made.

Determination of the relative yield of evaporation neutrons from various elements bombarded by ~ 350 MeV neutrons.

Introduction

Nuclear disintegrations produced by high energy particles have been usually investigated with photographic plates. While this technique is extremely powerful, there are 2 features of the disintegration process which are **not conveniently studied in photographic plates**. These are:

- 1) The dependence of the star properties upon the atomic number of the exploding nucleus
- 2) The neutron emission in the explosion.

The aim of this work was to investigate the production of neutrons in various material bombarded by neutrons of ≈ 360 MeV. **Relative measurements of the number of neutrons emitted per inelastic collision (star) were made.**

Pagina 40: Draft "Determination of the relative yield of evaporation neutrons from various elements bombarded by ~ 350 MeV neutrons."

The results ~~will be compared~~ ^{will be compared} in relation ^{compare} with photograph plate evidence, ~~and with~~ ^{and with} cosmic ray experiments on neutron production.

(For neutrons of 120 MeV the average number of ~~neutrons~~ ^{evaporation} ~~evaporation neutrons~~ ^{neutrons} ~~was recently measured~~ ^{was recently measured} in U and Pb ~~with an absolute method~~ ^{with an absolute method} and found to be about 14 and 8).

The method ~~used for~~ ^{used for} detecting neutrons was selected for his simplicity: it consists of a BF_3 counter imbedded in paraffine. It ~~is known that~~ ^{is known that} ~~there~~ ^{there} differences in the spectrum of evaporated neutrons may exist, however, it is ~~necessary~~ ^{necessary} to ~~be sure~~ ^{be sure} that ~~the~~ ^{the} detector is at least approximately equally sensitive to constant ~~independ~~ ^{independ} of the energy. For this reason the ~~form~~ ^{form} of the counter was of the "long counter" type ^{of} which ~~the~~ ^{the} ~~approximate~~ ^{approximate} independence of the energy between 5000 eV on 10^6 MeV. ~~The~~ ^{The} detector is consequently

The results will be compared with photograph plate evidence, and with cosmic ray experiments on neutron production. (for neutrons of 120 MeV the average number of evaporation neutrons was recently measured in U and Pb with an absolute method and found to be about 14 and 8).

The method for detecting neutrons was selected for his simplicity: it consists of a BF_3 counter imbedded in paraffine, since differences in the spectrum of evaporated neutrons may exist, however, it is necessary that the sensitivity of the detector is at least appromimately equally sensitive to constant independent of the energy. For this reason the geometry of the counter was of the "long counter" type ^{of} ,, of what(?) ?? approximate independence of the energy between 50000 eV on 10^6 MeV.

The detector is consequently satisfactory to detect "evaporation neutrons from stars, of which by far the great majority have energy < 20 MeV ^{of}

Pagina 47 : Draft "Determination of the relative yield of evaporation neutrons from various elements bombarded by ~ 350 MeV neutrons."

Acknowledgements

~~It is a pleasure to thank M G for his interest on for a discussion on stars, for I S a jerk for his competent help.~~

This value is ~~added to~~ v_a , then the ratio $(v_Z/v_{pb})_a^{calc}$ would be about 1.4, to compare with the observed value 1.4. Such good agreement is obviously a matter of chance.

Conclusion page 43

Acknowledgements

Mich. Gerig, George Ivon,
Crima, Dimchev

Acknowledgements

~~It is a pleasure to thanks M G for his interest on for a ..discussion on stars ?? I S A ?? For his competent help.~~

This value is added to v_a , then the ratio $(v_Z/v_{pb})_a^{calc}$ would be about 1.4, to compare with the observed value 1.4. Such good agreement is obviously a matter of chance.

Conclusions page 43

Acknowledgements

Mich. G.,

Pagina 43: Draft "Determination of the relative yield of evaporation neutrons from various elements bombarded by ~ 350 MeV neutrons."

is approximately known to be $\frac{1}{2}$ of the known total cross section (σ) i.e. $1.7 \times 10^{-24} \text{ cm}^2$. The fission cross section is about 0.5×10^{-24} for neutrons of 340 MeV, and assuming the same value of fission section in U for our neutron beam, the number of fissions per elastic collision is ≈ 0.3 .

Since the number of neutrons produced in the fission process is of the order of 2-3, a fission in U will result in:

$$N = N_0 + (2-3) \times \frac{\sigma_f}{\sigma} = N_0 + (2-3) \times 0.3 \approx N_0$$

Conclusions

- 1) The ratio $\frac{N_2}{N_1}$ obtained was comparable with the ratio obtained in studying the neutron production in our measurements with multiplicities in our detectors. It is of moderate energy neutrons (3-15 MeV) in cosmic radiation that the great majority of cosmic ray neutrons are produced in "stars" produced by nucleons of few hundred MeV.
- 2) The high multiplicity number of neutrons produced per collision in heavy elements by neutrons of few 100 MeV, and the fact that the secondary neutrons can easily escape from a big dimensions bloc, support the use of a such

Conclusions

- 1) The comparison of our measurements with measurements of multiplicities of moderate energy neutrons (3-15 MeV) in cosmic radiation confirms directly that the great majority of cosmic ray neutrons ??? in "stars" produced by nucleons of few hundred MeV.
- 2) The high number of neutrons produced per collision in heavy elements by neutrons of few 100 MeV, and the fact that the secondary neutrons can easily escape from a big dimensions bloc, support the use of a such

Pagina 44: Draft "Determination of the relative yield of evaporation neutrons from various elements bombarded by ~ 350 MeV neutrons."

block in connection with a ~~long~~ BF₃ counter, or
 as a convenient indicator of high energy neutrons (star detector). (1)

3) ~~The ratio of~~ ^{transmitted} ~~photons emitted~~
 A ~~block~~ ^{photographic plate} evidence with ~~our results~~ ^{photographic plate} not
~~explained~~ ^{by taking into account} the ~~neutron~~ ^{neutron} ~~loss~~ ^{and} ~~absorption~~ ^{of the} ~~in~~ ^{evaporation}
 the ~~most~~ ^{neutrons} ~~of~~ ^{of an explosion in} ~~various~~ ^{elements} are ~~photo~~ ^{photo} ~~neutron~~ ^{neutron} ~~deficient~~ ^{deficient}.

4) ~~Our~~
 3) ~~Comparison~~ It is possible to obtain ~~on the basis of photographic plate work, and the present work,~~ ^{It is possible} to ~~perform~~ ^{analyze} satisfactorily ~~the~~ ^{our} relative values of neutron production on the basis of ~~our photographic plate evidence~~ (see the ~~present work~~ ^{present work}).

4) ~~Another~~ ^{the} ~~best~~ ^{best} ~~method~~ ^{method} ~~of~~ ^{of} ~~estimating~~ ^{estimating} ~~the~~ ^{the} ~~number~~ ^{number} ~~of~~ ^{of} ~~neutrons~~ ^{neutrons} emitted in a ~~star~~ ^{star} ~~is~~ ^{is} about 1.7 times the number of ~~charged particles~~ ^{charged particles}. The absolute ~~number~~ ^{number} of evaporated neutrons per disintegration in various elements is presented in column v^a of table II.

Bloc in connection with a "long" BF₃ counter detector of neutrons of the order of few MeV is a convenient indicator of high energy neutrons (star detector). (1)

3) It is possible to analyze satisfactory our relative values of neutron production on the basis of photographic plate evidence (see $(v_Z/v_{Pb})_a^{calc}$, $(v_Z/v_{Pb})^{obs}$ in the table II) According to the analysis the average number of neutrons emitted in a star in Ag or Pb is about 1.7 times the number of ??? charged particles. The absolute ??? average number of evaporated neutrons per disintegration in various element is presented in column v^a of table II.

Schematic estimate of the average value of neutrons evaporated in a star with $1/\xi$ charged prongs

in addition to protons (). The ratio of the double charged to single charged particle is taken from ref (). The results calculated are in table II

The ratio of the average number of neutrons produced per star, calculated on the ???, can be compared with the observed one in the 2 last columns.

Schematic estimate of the average value of neutrons evaporated in a star with $1/\xi$ charged prongs

Element	Z	Atomic weight	ξ	Double ch/Single ch	Number of single charged	Z _{fin}	Final Nucleus	Number of evaporated neutrons	$(v_Z/v_{Pb})_{calc.}$	$(v_Z/v_{Pb})_{obs.}$	Note
U	92	238	0.18	0.3	4.2	85	At ²¹⁵	16.0	1.3	1.6201	
Pb	82	207	0.10	0.3	3.9	76	Pb ²¹⁵	12.6	1.0	1.0	
Sn	50	119	0.15	0.3	3.7	45	Al ¹¹⁵	9.0	0.7	0.62007	
Ag	47	108	0.27	0.3	3.0	42	Mn ¹⁰⁵	8.0	0.6		
Ba	35	80	0.33	0.4	2.2	31	Ca ⁶⁸	6.0	0.5		
Cu	29	64	0.40	0.5	1.7	26	Fe ⁵⁵	4.6	0.3	0.341004	
(C)	6	12	0.7	1.3	0.6	4	Be ⁷	0.25	0.05	0.2101	This ?? is probably meaningless in light element stars

In addition to protons (). The average ratio double/single double charged to single charged particle is taken from ref (). The results calculated are in table II
 The ratio of the average number of neutrons produced per star, calculated on the ???, can be compared with the observed one in the 2 last columns.

Table II

- Schematic estimate of the average value of neutrons evaporated in a star with $1/\xi$ charged prongs-
 Element, Z, Atomic Weight, ξ , Double ch/Single ch, number of single charged, Z_{fin}, Final Nucleus, Number of evaporated neutrons, $(v_Z/v_{Pb})_{calc.}$, $(v_Z/v_{Pb})_{obs.}$, Note

This ?? is probably meaningless in light element stars

 Number in a $(\xi+1)/\xi$ prong stars

- Light element meaningless -



Schematic estimate of the average value of neutrons evaporated in a star with $1/\xi$ charged prongs

Table II ^{approx}
 Schematic estimate of the ~~possible~~ value of neutrons evaporated in a star with ~~1/ξ~~ charged prongs - 1/ξ

Element	Z	Atomic weight	ξ	Double ch. Single ch.	Number of single charged	Z bin	Fission Nucleus	Number of evaporated neutrons	$(\frac{N_n}{N_{fiss}})_{cal}$	$(\frac{N_n}{N_{fiss}})_{obs}$	Note
U	92	238	0.18	0.3	4.2	85	At ₈₅ ²¹⁰	16.0	1.3	1.6 ± 0.1	
Pb	82	207	0.20	0.3	3.9	76	Pb ₈₂ ¹⁸⁵	12.6	1.0	1.0	
Sn	50	119	0.25	0.3	3.1	45	Pb ₄₅ ¹⁰²	9.0	0.7	0.62 ± 0.07	
Ag	47	108	0.27	0.3	3.0	42	Mo ₄₂ ⁹³	8.0	0.6		
Bz	35	80	0.33	0.4	2.2	31	Cr ₃₁ ⁶⁸	6.0	0.5		
Cu	29	64	0.40	0.5	1.7	{25 {26	{Fe ₂₅ ⁵⁵ {Fe ₂₆ ⁵⁴	{3.6 {4.6	0.3	0.34 ± 0.04	
(C)	6	12	0.7	1.3	0.6	4	Be ₄ ⁷	1 0.5	0.55	0.2 ± 0.1	First fission is probably measurement on light elements & 2000

Table II

Schematic Aestimate of the average value of neutrons evaporated in a star with $(\xi+1)/\xi$ charged prongs

Element, Z, Atomic Weight, ξ , d/s ratio of double to single ch, number of single charged, Z_{fin} , Mass number $A_{final}(a)$, , Mass number $A_{final}(b)$, v_a^{calc} , v_b^{calc} , $(v_Z/v_{Pb})_a^{calc}$, $(v_Z/v_{Pb})_b^{calc}$, $(v_Z/v_{Pb})^{obs}$.

Table II

Schematic Aestimate of the average value of neutrons evaporated in a star with $(\xi+1)/\xi$ charged prongs

Element	Z	Atomic weight	ξ	d/s ratio of double to single charged	number of single charged	Z final	Mass number A final (a)	Mass number A final (b)	v_a^{calc}	v_b^{calc}	$(v_Z/v_{Pb})_a^{calc}$	$(v_Z/v_{Pb})_b^{calc}$	$(v_Z/v_{Pb})^{obs}$
U	92	238	0.18	0.3	5.1	84	212	207	13.4	18.4	1.3	1.4	1.4
Pb	82	207	0.20	0.3	4.6	74	184	181	10.4	13.4	1	1	1
Sn	50	119	0.25	0.3	3.8	44	102	97	7.5	12.5	0.7	0.9	0.62
Ag	47	108	0.27	0.3	3.6	41	93	90	6	9	0.6	0.7	
As	35	80	0.33	0.4	2.9	30	67	65	5	7	0.5	0.5	
Cu	29	64	0.40	0.5	2.3	24	53	51	3	5	0.3	0.4	0.34

Tabella II della slide precedente, rimessa a pulito e trascritta su un foglio separato.

Pontecorvo si tiene ben informato sugli esperimenti fatti in occidente su π^+/π^-

Data on π^+/π^-

76,
(1467A-

76,588,A
77,526,A
79,198

Hoyer, Maden, Hilsenrath, Knoble, Hales, Bull. Am. Phys. Soc., Aprile 1951, pp 15: $p+p \rightarrow \pi^0$

78, 497, 1950 - Chew and Steinberger

Pauli principle, in normal interpretation, 2 effects:

- 1) $\frac{\pi^+}{\pi^-}$ increase at high meson energy, for a given proton energy
- 2) $\frac{\pi^+}{\pi^-}$ increase with decreasing proton energy.
- And viceversa with neutron bombardment.

- 78, 85, 1950, Richmon and Wilson -

340 MeV protons: $\frac{\pi^+}{\pi^-}$ in C = 5 ± 1.5

Weissbluth, 78, 86, 1950

$\frac{\pi^+}{\pi^-} = 1.5 \pm 1$ in Pb, p 340 MeV

Brodner and Jones, 78, 90, 1950

$\frac{\pi^+}{\pi^-}$	Be	2.7 ± 2
	C	4.8 ± 5
	Al	5.4 ± 1
	Cu	4.3 ± 2

Data on π^+/π^-

75,1467 A; 76,588, A; 77,526 ; 79,198

Mayer (?) et al, , Aprile 1951, pp 15. $p+p \rightarrow \pi^0$

78, 497, 1950 - Chew and Steinberger

Pauli principle, in normal interpretation, 2 effects:

- 1) π^+/π^- increase at high meson energy, for a given proton energy
 - 2) π^+/π^- increase with decreasing proton energy
- And viceversa with neutron bombardment

- 78, 85, 1950, Richmon and Wilson -

380 MeV protons: π^+/π^- in C = 5 ± 1.5

Weissbluth, 78, 86, 1950

$\pi^+/\pi^- = 1.5 \pm 1$ in Pb, p 340 MeV

Brodner and Jones, 78, 90, 1950

π^+/π^- Be 2.7 ± 2 ; C 4.8 ± 5 ; Al 5.4 ± 1 ; Cu 4.3 ± 2

Brauner et al. Phys. Rev. 79, 720

n 270 MeV $\pi^+/\pi^- = 14$

Brauner, not published:

π^+/π^- , protons 260-270 MeV = ≈ 15

Brauner, Cornell, Princeton

Phys. Rev. 79, 720.

n 270 MeV.

$\frac{\pi^-}{\pi^+} = 14$

61

50 MeV mesons

Brauner, not published:

$\frac{\pi^+}{\pi^-}$

protons 260-270 MeV

= ≈ 15

Pag. 64: Seminario di Pontecorvo sul suo programma di ricerca sulle collisioni nucleone-nucleone e sulla misura del rapporto π^+/π^- .

~~Seminario~~ - Seminario

I want to talk on nucleon-nucleon collisions, and $\frac{\pi^+}{\pi^-}$ ratios -
 Discussion of experimental method - Conclusion on what are the collision which produce mesons -

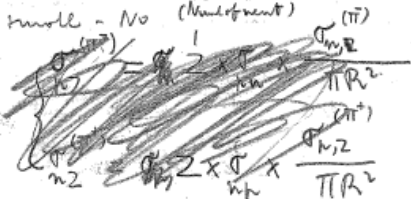
Usually $p+p \rightarrow \pi^+ + n + p$ ✓
 $n+n \rightarrow \pi^- + n + p$ ✓
 $n+p \rightarrow \begin{cases} \pi^+ + n + n & \text{v ✓ (neutron beam (n+Z))} \\ \pi^- + n + p & \text{v ✓ (proton beam (p+Z))} \end{cases}$

What is the evidence?

~~p+Z~~ $p+Z \rightarrow \pi^-$
 $n+Z \rightarrow \pi^+$

since $\begin{cases} p+p \rightarrow \pi^- \\ n+n \rightarrow \pi^+ \end{cases}$, must be $p+n$
 must be $n+p$.
Not true proof, reasonable

Z p n,n
 Too small - No (Misstatement)



Assume $n+n \rightarrow \pi^-$, then

$$\sigma_{n,Z}^{\pi^-} = \sigma_{\text{neut, recoil}} \times \frac{\sigma_{n,Z}^{\pi^-}}{\pi R^2}$$

since $\sigma_{\text{neut, recoil}} / \pi R^2$

is a big number, let us say $\approx 1/5$, it is clear

that this hypothesis not against facts.

basically from the fact that $n+Z \rightarrow \pi^+$

not possible conclude that etc. Δh

Conclusion:

only things we know:

$p+p \rightarrow \pi^+$
 $p+p \rightarrow \pi^0$
 $n+n \rightarrow \pi^0$

- Seminaire -

I want to talk on nucleon-nucleon collisions, and π^+/π^- ratios -

Discussion of experimental method -

Conclusion on what (??) are the collision which produce mesons

Usually $p+p \rightarrow \pi^+ + n + p$
 $n+n \rightarrow \pi^- + n + p$
 $n+p \rightarrow \begin{cases} \pi^+ + n + n & \text{(neutron beam (n+Z))} \\ \pi^- + p + p & \text{(proton beam (p+Z))} \end{cases}$

What is the evidence?

$p+Z \rightarrow \pi^-$
 $n+Z \rightarrow \pi^+$

Since $\begin{cases} p+p \rightarrow \text{no } \pi^- \\ n+n \rightarrow \text{no } \pi^+ \end{cases}$, must be $p+n$
 must be $n+p$

Not true proof, reasonable

Z p n,n

Too small - No

Assuming only (?) $n+n \rightarrow \pi^-$, then

$$\sigma_{p,Z}^{\pi^-} = \sigma_{\text{neut, recoil}} \times \sigma_{n,Z}^{\pi^-} / \pi R^2$$

since $\sigma_{\text{neut, recoil}} / \pi R^2$

is a big number, let us say $\approx 1/5$, it is clear that this hypothesis not against facts.

Similarly from the fact that $n+Z \rightarrow \pi^+$, not possible conclude that $n+p$ etc. $p+p$

Conclusion:

only ??? we know:

$p+p \rightarrow \pi^+$
 $p+p \rightarrow \pi^0$
 $p+n \rightarrow \pi^0$

Valutazioni sulla fattibilità di un "H⁴ experiment" (pag. 13)

- H⁴ experiment -

~~Energy~~ Possibility of measuring α recoil; which will be of the order of (neglect neutrinos)

$$Mv = \frac{E_\alpha}{c} \quad E_\alpha = \frac{1}{2} M_\alpha \frac{E^2}{M_\alpha^2 c^2} = \frac{1}{2} \frac{E^2}{M_\alpha c^2}$$

$$= \frac{1}{2} \frac{E^2}{M_\alpha c^2} = \frac{1}{2} \frac{(20 \times 10^6)^2}{4 \times 10^9} = \frac{1}{2} \times 10^5 \text{ eV} \approx$$

$\approx 50000 \text{ eV}$. This could be

measured in a small prop. counter

Let us assume that the H⁴ track have 20 cm air range. The number stopping in the counter will be of the order of $\frac{\text{counter diameter (cm)}}{\text{range (cm)}}$

If we have a 0.1 cm counter with A, then the number stopping is of the order of $\frac{1}{200} \approx \frac{1}{1000}$ of those produced in a thickness = range, in solid matter, i.e.: in a volume = $2\pi \times 0.1 \times 3$

- H⁴ experiment -

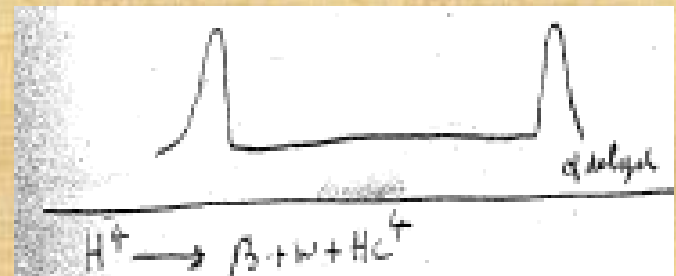
Possibility of **measuring α recoil** which will be of the order of (neglect neutrinos)

$$Mv = E_\alpha / c \quad E_\alpha = \frac{1}{2} M_\alpha E^2 / (M_\alpha^2 C^2) = 1/2 E^2 (M_\alpha / C^2)$$

$$= \frac{1}{2} (20 \times 10^6)^2 / 4 \times 10^9 = \frac{1}{2} 10^5 \text{ eV}$$

$\approx 50000 \text{ eV}$. This could be measured in a **small prop. counter**.

Let us **assume** that the H⁴ track have **20 cm avr (?) Range**. The number stopping in the counter will be of the order of $\frac{\text{counter diameter (cm avr)}}{\text{Range (cm avr)}}$. If we have a 0.1 cm counter with A, then the number of stopping is of the order of $1/200 \approx 1/1000$ of those produced in a thickness range, in solid matter, i.e.: in a volume = $2\pi \times 0.1 \times 3$



H⁴ ipotetico stato legato di tre neutroni ed un protone: non è mai stato osservato.

Valutazioni sulla fattibilità di un "H⁴ experiment" (pag. 13)

The number produced is, per cm³,
 $10^{-28} \times N \times 2 \times 10^{19} / \text{sec}$. To have
 an intensity of 10/millisec, or 10⁴/sec aver,
 is necessary a ~~fantastic~~ **fantastic** intensity.
 No good. (Unless: we irradiate with 0.1 per A produced)

13

Let us calculate with a big counter.
 Let us say. Length 40 cm,
 diameter 6 cm. Volume
 $\pi \times 9 \times 40$ cm 1 liter, atmospheric pressure.

This gives a rate of H⁴ equal to
 $10^{-28} \times 2 \times 10^{22} \times N \cdot \frac{\text{intens av}}{\text{sec}}$ Per Pulse

~~$10^{-28} \times 2 \times 10^{22}$~~ Intensity ave per pulse =

$\frac{1}{40} \times 10^{-28} \times 2 \times 10^{22} \times N = 10$, If this = 10,

then N must be = 10⁸/sec/cm². We could
 put that in the channel, but the
 intensity is too big at "start down".

The number produced is, per cm³,
 $10^{-28}(\sigma) \times N \times 2 \times 10^{19} / \text{sec}$. To have
 an intensity of 10/millisec, or
 10⁴/sec aver. **is necessary a ~~N~~of
 fantastic intensity. No good.**
 (unless: we irradiate ?? 0.1 per A
 protons (?))

Let us calculate **with a big
 counter**, let us say Length 40 cm,
 diameter 6 cm Volume $\pi \times 9 \times 40$
1 liter, atmospheric pressure. This
 give a rate of H⁴ equal to
 $10^{-28} \times 2 \times 10^{22} \times N$ intens av/sec.
 Intensity ave per pulse = $1/40 \times$
 $10^{-28} \times 2 \times 10^{22} \times N = 10$
 If this = 10, then N must be
 = **10⁸/sec/cm²**. We could ?? that in
 the channel, but the **intensity is
 too big** at "?? ??"

Valutazioni sulla fattibilità di un "H⁴ experiment" (pag. 14)

Ex. If we have a proton beam of 10^8 particles, then we might use the big counter technique.

Other version of H⁴ experiment

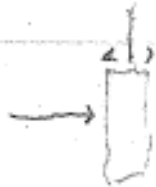
A "solution" scintillating (≈ 1 liter) is irradiated by a beam, One measures delayed β , biased.

Assuming first a 10^6 n/cm²/sec beam, like it would be inside the window, in the hole. Then the intensity of H⁴ is $10^{-29} \times 10^{25} \times 10^6 = 100$ /sec ~~at 10^6 sec~~

OK!!!!

~~10^6 sec~~
 10^7 sec

- Check the range of β particles of 10 MeV and 20 MeV -



10^{-32} cm

If we have a proton beam of 10^8 particles, then we might use the big counter technique.

Other version of H⁴ experiment

A "solution" scintillating (≈ 1 liter) is irradiated by a beam, One measures delayed β , biased.

Assuming ?? a 10^6 n/cm²/sec beam, like it would be inside the ??, in the hole. Then the intensity of H⁴ is $10^{-29} \times 10^{25} \times 10^6 = 100$ /sec ~~≈ 1 min too small!~~

OK!!!!

- Check the range of β particles of 10 MeV and 20 MeV -

Fisica ma anche Poesia

50kV

$4 \times \frac{1800}{276}$

77

$\frac{245}{10}$

4

$H^+ \rightarrow \beta + \nu + He^4$		$Li^+ - \text{unstable probably}$
$Be^8 \rightarrow \beta + \nu + Be^8$	x o x	$Be^8 \rightarrow \beta + \nu + Li^8$
$B^{12} \rightarrow \beta + \nu + C^{12}$	x o?	$N^{12} \rightarrow \beta + \nu + C^{12}$
$N^{16} \rightarrow \beta + \nu + O^{16}$	x	$F^+ \rightarrow \text{unstable}$
$Na^{20} \rightarrow \beta + \nu + Ne^{20}$	x	$Na^{20} \rightarrow \beta + \nu + Ne^{20}$

- For meson work -
Delayed α emission debate:

1) $He^4 + \pi^- \rightarrow He^4$ (fast) \downarrow β -emission

2) a) $C^{12} + \pi^- \rightarrow B^{12}$ (fast)?

b) $C^{13} + \pi^- \rightarrow B^{12} + n$

$C^{13} + \pi^+ \rightarrow N^{12} + n$

$C^{12} + \pi^+ \rightarrow N^{12}$ (fast)

$Ne^{20} + \pi^+ \rightarrow Na^{20}$ (slow)

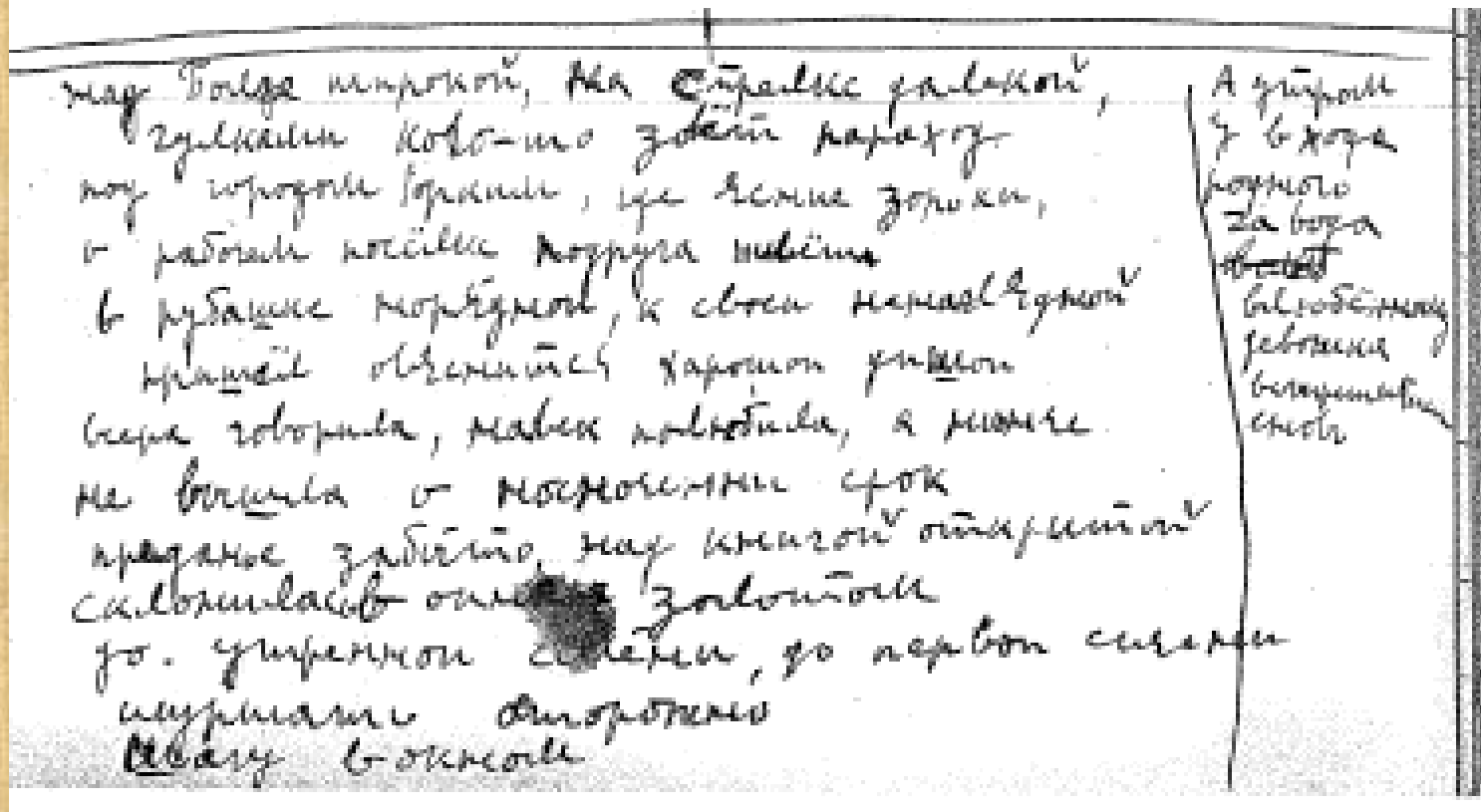
на Волге широкой, на Сормовской долине,
 великим колесом зовут пароходы
 по широкой реке, на яхте золотой,
 в рубашке белой пошла невеста
 в рубашке белой, к слову неслучайно
 пришла обвенчаться паромом ужином
 была золотая, была белая, а невеста
 на Волге в просторном спок
 преданье забыто, на книжке отпечатан
 саблонная опись золотой
 до уфимской слесари, до первого сигнала
 шумными пароходами
 везут в Сормовск

$H^4 \rightarrow \beta + \nu + He^4$ | $Li^4 \rightarrow \text{unstable probably}$
 $Li^8 \rightarrow \beta + \nu + Be^8$ x o x | $Be^8 \rightarrow \beta + \nu + Li^8$ o
 $B^{12} \rightarrow \beta + \nu + C^{12}$ x o? x | $N^{12} \rightarrow \beta + \nu + C^{12}$ o
 $N^{16} \rightarrow \beta + \nu + O^{16}$ x | $F^{16} \rightarrow \text{unstable}$
 $F^{20} \rightarrow \beta + \nu + Ne^{20}$ x x | $Na^{20} \rightarrow \beta + \nu + Ne^{20}$ o

- For meson work -
Delayed α emission ??? :
- 1) $He^4 + \pi^- \rightarrow H^4$ (fast)
 $\rightarrow \beta^-$ and He^4 recoil
 - 2) a) $C^{12} + \pi^- \rightarrow B^{12}$ (fast)? | $C^{13} + \pi^+ \rightarrow N^{12} + n$
 | $C^{12} + \pi^+ \rightarrow N^{12}$ (fast)
 b) $C^{13} + \pi^- \rightarrow B^{12} + n$ | $Ne^{20} + \pi^+ \rightarrow Na^{20}$ (slow)

Lirica di Sormovo
(Сормовская лирическая)
 Canzone d'amore (diventata poi canzone) di
Yevgeniy Aronovich Dolmatovsky
*Parla di un fidanzato che si veste a festa e
 va nella cittadina di Sormovo a trovare la
 sua bella, ma viene da lei deluso.*

Il testo russo della poesia



На Волге широкой,
На стрелке далёкой
Гудками кого-то зовёт парашод.
Под городом Горьким,
Где ясные зорьки,
В рабочем посёлке подруга живёт.

В рубашке нарядной
К своей ненаглядной
Пришёл объясниться хороший дружок:
Вчера говорила —
Навек полюбила,
А нынче не вышла в назначенный срок.

Свиданье забыто,
Над книгой раскрытой
Склонилась подруга в окне золотом.
До утренней смены,
До первой сирены
Шуршат осторожно шаги под окном.

Ой, летние ночки,
Буксиров гудочки...
Волнуется парень и хочет уйти.
Но девушки краше,
Чем в Сормове нашем,
Ему никогда и нигде не найти.

А утром у входа
Родного завода
Влюблённую девушка встретится вновь
И скажет: «Немало
Я книжек читала,
Но нет ещё книжки про нашу любовь».

<http://poetryrain.com/authors/dolmatovskiy-evgeniy/10919>

<https://www.youtube.com/watch?v=lg0g5-2fk5k>

Pagina 85: Siamo ad estate 1951 (~ July 11) e Pontecorvo cerca di definire il programma per l'anno successivo.

Plan for the year

85

I) Anatol Alex + Vladimir.

Investigation of the possibility of ^(a stable state) existence of H^4 , and, in case of existence of H^4 , inquiry of its properties.

The problem is interesting from the point of view of the ?? of light number ????? on the ?? β decay. Preliminary experiments should be made within the next few months. The first experiment consists in curving the (hypothetical) β particles from H^4 in 3 counters in coincidence, placed at a distance of ≥ 10 cm from the cyclotron target. This method should give a rough measurement of H^4 , while a measurement of the lifetime (expected value $\approx 10^{-3}$ sec) is made by electronic methods. The H^4 could be produced in the target by nuclear interactions (?): for example Li^8 excite $\rightarrow H^4 + He^4$. If the first experiment is successful, it will take about a year to investigate the properties of H^4 (spectrum, lifetime), and especially to study in what condition it is produced. If the first experiment is not successful, other methods should be considered: these are a) detection of H^4 in a long liquid organic scintillation counter, biased so that it does not register β particles of $E_{\beta} \leq 14$ MeV, and ?? by high energy neutron. b) Detection (on a prop. count. filled with He) ? of delayed α particles (\approx few tens KV) corresponding to the recoil of β from H^4 experiment. The 50 MeV π^- meson should be used to produce H^4 according to $\pi^- + He^4 \rightarrow H^4 + He^4$. Unfortunately the interest of the problem

Plan of the year

85

I) Anatol Alex + Vladimir

Investigation of the possibility of a stable state of H^4 , and, in case of existence of H^4 investigating of its properties.

The problem is interesting from the point of view of the ?? light number ????? on the ?? β decay. Preliminary experiments should be done within the next few months. The first experiment consists in curving the (hypothetical) β particles in the cyclotron magnetic field and registering them in 3 counters in coincidence, placed at a distance of ≥ 10 cm from the cyclotron target. This method should give a rough measurements of H^4 , while a measurement of the lifetime (expected value $\approx 10^{-3}$ sec) is made by electronic methods. The H^4 could be produced in the target by nuclear interactions (?): for example Li^8 excite $\rightarrow H^4 + He^4$. If the first rough experiment is successful, it will take about a year to investigate the properties of H^4 (spectrum, lifetime) and ?? also to study in what condition it is produced. If the first experiment is not successful, other methods are considered: these are a) detection of H^4 in a long liquid organic scintillation counter, biased so that it does not register β particles by Energy ≤ 14 MeV, and ?? by high energy neutron. b) Detection (on a prop. count. filled with He) ? of delayed α particles (\approx few tens KV) corresponding to the recoil of β from H^4 $H^4 \rightarrow He^4 + \beta + \text{neutrino}$. In this experiment, the 50 MeV π^- meson should be used to produce H^4 according to $\pi^- + He^4 \rightarrow H^4 + He^4$ but unfortunately the interest of the problem

Pagina 86: Siamo ad estate 1951 e Pontecorvo cerca di definire il programma per l'anno successivo.

of H^+ is only if a stable (versus ???) H^+ exists. In other words a negative result on H^+ is not significant.

However, even if the H^+ experiments are negative, the techniques developed may be used in other problems, ~~concerning a fast meson~~ and like the investigation of short life radioelements (B^{12} etc) under meson bombardment.

II) Application of the method of the radioactive indicators to the investigation of properties of

mesons, in particular ^{charge exchange} ~~concerning the~~ ~~charge exchange~~

When mesons interact with nuclei, disintegration takes one charge exchange (emission of a π^0) is possible. In both cases radioelements may be produced, according to the nucleus (For example



can be used a) To determine the process of formation, determination of the cross section for meson charge exchange for positive and negative mesons ("fast" + and - π)

b) To use the indicators

(radioelements) as meson indicators, in experiments on mesons.

of H^+ is only if a stable (versus ???) H^+ exists. In other words a negative results on H^+ is not significant.

However if the H^+ experiments are negative, the techniques developed may be used in other problems like the investigation of short life b radioelements (B^{12} etc) under meson bombardment.

Adolph

II) Application of the method of radioactive indicators to the investigation of properties of mesons:

When charged mesons interact with nuclei, stars and charge exchange (emission of a π^0) are possible. In both cases radioelements may be produced (For example $B^{10} + \pi^+ \rightarrow C^{11} + \pi^0$) These radioelements can be used a) To determine the cross section for meson charge exchange, for positive and negative mesons ("fast" + and - π 25 Mev or 50 MeV)

b) To use the radioelements as meson indicators in experiment on mesons.

Pagina 87: Siamo ad estate 1951 e Pontecorvo cerca di definire il programma per l'anno successivo.

double meson production. (This necessitates a chemical technique).



III ~~Supplement~~ Georgy Ivan Selivan
Investigation Properties of Interaction with
neutrons
- Development of techniques capable of detecting
electronically ~~of~~ mesons - ~~for the~~ Investigation
of π^+ production in hydrogen ^{neutron collisions} of neutrons -

~~III Investigation of the~~
The method consists in detecting the
 π^+ from the $2.2 \mu s \mu^+ - e^+$.

~~III~~ The n-p has special interest,
because until now there is no certain
evidence that charged mesons are intensively
formed in n-p collisions.

IV Direct detection of the meson beam (and -)
from the cyclotron ^{in the} with ~~chemical~~ ^{electronic} methods.
Application to an increase of π^+/π^- ratio.
Application to change of $\frac{\pi^+}{\pi^-}$ ratio with Z.
The method ^{will} use small phototubes or
scintillation counters in coincidence.

V Development of Cerenkov detectors
for the study of relativistic particles

double meson production (this necessitates a chemical technique).

III Georgy Ivan Selivan

-Development of techniques capable of detecting electronically mesons. Investigation of π^+ production in hydrogen and other elements by neutrons.

The method consists in detecting the π^+ from the $2.2 \mu s \mu^+ - e^+$.

The n - p has special interest because until now there is no certain evidence that charged mesons are intensively produced in n-p collisions

IV Direct detection of the meson beam (+ and -) in the cyclotron, with electronic methods ??? Counters.

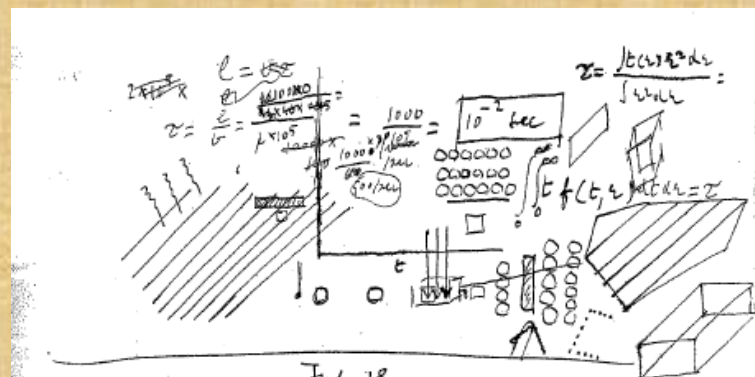
Application to the **measure of π^+/π^- ratio.**

Application to change of π^+/π^- ratio with Z.

The method will use small photoplates?(?) counters or scintillation counters in coincidence

V Development of Cerenkov detectors, for the study of relativistic particles.

Pagina 91: July 18, ultima data presente sul quaderno che contiene ancora 9 pagine. Si inizia ad attuare il programma stabilito con due esperimenti in contemporanea nelle sale sperimentali room 3 e room 2 (con fascio di neutroni).



July 18

- Box No 1, in Room 3 of cyclotron building -

Voltage: 830V. \downarrow Cyclotron does not work:

24x64 + 25 /min
24x64 + 50

Maximum* Cyclotron works with intensity

34x64 + 32
42x64 + 28
42x64 + 24
44x64 + 0

Cyclotron does not work

23x64 + 49

Room No 2, with neutron beam passing through 30x30 open window, and fenced shielding very bad with bricks:

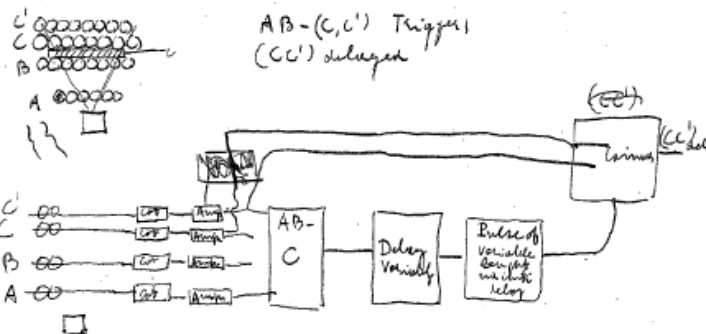
Near the hole close to detector room:

$\approx 70x64$ /min
 $\approx 60x64$ /min

Maximum intensity. nows "maximum" for the time being, i.e. $\approx \frac{1}{5} - \frac{1}{10}$ times less than usual, because the voltage on D must be small

Room No 2
At a distance of ≈ 100 cm from the neutron beam: 120x64 (But there was apparatus!)
- 2 hrs waiting to ask period time, about 2 hours, to measure carefully.
Conclusion: ~~nothing was found about the beam~~
since the beam was very wide, but (and the shielding of the big window poor, but the cyclotron intensity low, the background conditions are \approx like they will be:

- Beam of a key ≈ 10000 /min -
Calculations of beam on
- Experiment on Π^+



Background:

$(AB-C) \times (C') \times \tau$, small - , because

counting has been decreased by micron order, count AB-C and CC' are small and AB-C and CC' are respectively the actual particles from radiator.

Pontecorvo riprende a scrivere sul primo quaderno il 14 settembre 1951 iniziando dall'ultima pagina (la numero **100**) dopo aver capovolto il quaderno. Propone di fare l'esperimento di produzione di mesoni con un fascio di neutroni: **"Experiment on production of mesons by neutrons"**:

$\frac{1 \times 1538 \times 14}{8 \times 10^6} \times \frac{10^8 \times 14}{8 \times 10^6} \times 2 = \frac{1016}{10^6} = 1016 \times 10^{-6}$

$\pi^+ + n \rightarrow \pi^+ + p$
 $\pi^+ + n \rightarrow \pi^+ + \pi^0$
 $\pi^+ + n \rightarrow \pi^+ + p$
 $\pi^+ + n \rightarrow \pi^+ + p$
 $\pi^+ + n \rightarrow \pi^+ + p$
 $\pi^+ + n \rightarrow \pi^+ + p$

$2 N_1 N_2 Z = C$

061

14 September

Experiment on production of mesons by neutrons:

- π^0
It is necessary: 1) the "radiator" R
2) the "converter" C
3) the "absorber" A between the last counters
4) the absorber of γ radiation T

R \rightarrow The radiator must be a "sphere" of diameter = m.f.p. for γ . *exp. Diameter 10 cm approx.*

C \rightarrow The converter must be 1 cm Pb, area equal to the counter tray area.

A \rightarrow The absorber between counters must be 1 cm Al, area equal to the tray counter, small.

T \rightarrow Must be about 1 cm thin of Pb and 1 cm thin of Cu (to see that the collision is really γ), area equal to the tray counter.

The geometry as follows:

The detecting counters (five away) are the trouble key to increase the coincidence efficiency.

$\pi^0 \rightarrow \gamma + \gamma$

14 September

Experiment on production of mesons by neutrons:

- π^0
It is necessary:
 - the "radiator" R
 - the "converter" C
 - the "absorber" A between the 2 last counters
 - the absorber of γ radiation T

R \rightarrow The radiator must be a "sphere"

C \rightarrow The converter must be 1 cm Pb,....

A \rightarrow The absorber between counters

T \rightarrow Must be about 1 cm thin of Pb,....

The geometry as follows:

.....

.....

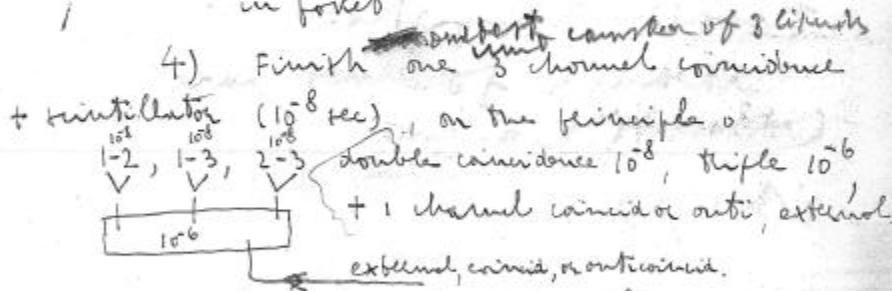
E continua a scrivere quello che oggi chiameremmo il "Technical Proposal" dell'esperimento...

A settembre 1951, meno di un anno dopo il suo arrivo a Dubna, Bruno Maximovich Pontecorvo è un rispettato **group leader di un gruppo di giovani fisici** (Vladimir, Anatol, Alex, Adolph and George Selivanov). In frequenti group meeting assegna ad ogni membro del gruppo il lavoro da fare secondo il programma stabilito e spesso suggerisce come realizzare certe misure

- 1) Vladimir: Finish work on H^4 in the present variant, + report. Help of Anatol, Alex. *Kirpa.*
- 2) Adolph: Finish work on mesons with radioactive indicators + report. Have β counters ready. Have α counter ready.

3) George

- 1) Finish work on duty factor + report.
- 2) Conclude on the work of production of π^0 in C by neutrons.
- 3) Initiate electronic detection of mesons in forest
- 4) Finish ~~one~~ ^{one} channel coincidence



Remember delay of 5×10^{-8} to measure accidental.

- 4) Anatol: finish work + write ² reports
 - a) angular
 - b) total work on secondary neutrons

Program

1) Measure the effective duty factor of the cyclotron as follows - From this:

- a) Measure the resolving time of the system (using continuous sources).
- b) Measure the accidental rates when the cyclotron works.

Experiment a: In the cyclotron building with a distance of two traps of about 10m, and with 2 sources, measure:

- 1) Single rates, and coincidence: $C-D-A-B$ min. vol. coinc.
- 2) Same with no source for collecting values in subsequent experiment.

Experiment b:

- 1) When cyclotron works, single rates and coincidence rates.

Measurement of α , duty factor.

- I) Take 2 single counters. +
- II) Put them on the beam, far away with from each other, and measure A B (AB)
- III) Verify that verify not the cyclotron is constant.
- IV) ~~In various conditions of~~ Measure also after shutdown, A, B (AB).
- V) Write all data relative to the cyclotron
- VI) If ^{the} experiments ^{are} reasonably reproducible, try various conditions of cyclotron.

L'attività di tutto il gruppo è molto ben documentata giornalmente

Workshop time - 180 Telescope
 *480 Protot.
660

It was decided to cancel * 660 for that the time for us will be:

180	} Telescope G.I
180	
100	} Other telescope.
100	
100	} small work
100	
<u>660</u>	} small chassi
	} Pent + Versum
	} prototipi.

Of these: \approx February 16

\approx 24^h television cylinders

- 1) ~~repeats~~
- 2) repeats
- 3) 2 cmolobn = }
- 4) Tommum
- 5) ~~Protot~~

27 February - Workshop time: 630.

2 telescope	: 400 hrs
Fe shield	: 25 hrs
chassi	: 30 hrs

15 Consider
 To test the various coincidences and anticoincidence efficiency, make the following experiments. Pro memoria:
 Big pink box: A 994
 Big black box: C' 1263 ~~1263~~
 Very wide tube box: A 885
 Other tube box: B. 885

- Ist experiment, testing an actual parameter

Fig 1

Measure $(ACBC')$

same but one channel

$$(ACBC') = 57x4+0 \quad 56x4+1$$

$$(ABC'-C) = 15x4+0 \quad 2x4+3$$

$$(ABC'-C) = 41x4+2 \quad 35x4+2$$

with no voltage on C.

II experiment

Fig 2

$$(AC'BC) = 113/2m$$

$$(ABC-C') = 53/2m$$

$$(ABC-C') = 140$$

with no voltage on C.

This is no good.

- Fig 2 - Repeat experiment after change of electronic.

$$(AC'BC) = 54x4+1 \quad 51x4+3$$

$$(ABC-C') = 15x4+0 \quad 15x4+3$$

$$(ABC-C') = 42x4+3 \quad 43x4+2$$

Dez ymnarov.

A	124	} x64	2m
B	46		
C	94		
D	115		
BCD-A	20		

Note: In this case the system was completed, but not the monitor! Take the time.

Dez ymnarov.

1	≤	c kolebrnypoll	
A		1400	} x64
B		514	
C		669	
D		656	
BCD-A		10x64+60	977/3m

Dez kolebrnypoll

2	≤	Dez kolebrnypoll	
A		1436	} x64
B		605	
C		693	
D		700	
BCD-A		5x64+26	1000/3m

Torlko c kolebrnypoll u Dez kolebrnypoll

3		883	} x64
A		421	
B		643	
D		651	
BCD-A		2x64+27	993/3m

Torlko c kolebrnypoll u c kolebrnypoll

4		847	} x64
A		571	
B		612	
D		615	
BCD-A		4x64+17	1007/3m

c kolebrnypoll

5	≤	1806	} x64
A		534	
B		694	
D		677	
BCD-A		6x64+33	982/3m

Tempo richiesto in officina meccanica per costruire il supporto dei rivelatori

Misure di efficienza di varie coincidenze e anticoincidenze dei rivelatori

La presa dati

Primi report interni sulla produzione di mesoni π

I risultati di tutti gli esperimenti fatti da Pontecorvo ed il suo gruppo nel periodo 1951-1954 al ciclotrone di Dubna furono pubblicati in russo come Report Interni del Laboratorio.

In questi primi esperimenti venne studiata la produzione di mesoni π , carichi e neutri, con protoni e neutroni su bersaglio di protoni e nuclei complessi.

La produzione del π^0 con fascio di neutroni su protoni e nuclei complessi fu studiata per la prima volta nel mondo. (B.M.Pontecorvo, G.I.Selivanov, RINP, 1951) and (B.M.Pontecorvo, G.I.Selivanov, RINP, 1952; Dokl. Acad. Nauk SSSR, 102, 253 (1955)).

- Production of neutral mesons π^0 by neutrons -

A) Introduction B) Apparatus C) Absolute experiment in Carb D) Production measurements Discussion in relation to production of mesons E) Polarization measurements F) Discussion ^{of the} conclusions - ^{of the} ~~production~~ ~~of~~ ~~mesons~~

Introduction

A considerable amount of data have been published in the last years on the production of ~~charged~~ mesons by protons. The ~~highest~~ production of ~~charged~~ mesons by neutrons has been so far only the object of a short communication and the production of neutral mesons by neutrons to ~~fast lead and~~ has been observed. The following table summarizing the present-day information on this subject.

Table I

It is clear from this table that production of charged and neutral mesons in ~~simple~~ collisions has not yet been observed and ~~only~~ in complex nuclei. The ~~same~~ production of neutral mesons by neutrons has not yet been observed. For this reason, because of the absence of data on this subject, it was ~~not~~ ~~possible~~ ~~to~~ ~~present~~ ~~some~~ ~~interests~~. In the present work we report experiments we have made utilizing the neutrons from the ~~synchrotron~~ of our laboratory, we have investigated (and observed, for the first time), the production of neutral mesons in Hydrogen and complex nuclei by neutrons. The ~~results~~ ~~of~~ ~~the~~ ~~production~~ ~~of~~ ~~neutral~~

УТВЕРЖДАЮ*

ДИРЕКТОР ФИЗИКО-МАТЕМАТИЧЕСКОГО ЦЕНТРА
И.А.Абрамзон (И.А.АБРАМЗОН)
 25.09.1952г.

НАУЧНО-ИССЛЕДОВАТЕЛЬСКИЙ ЦЕНТР Ф 62

О Т Ч Е Т
ОБРАЗОВАНИЕ π^0 -МЕЗОНОВ В (n-p) И (n-d) СТОЛКОВЕНИЯХ.

Начальник сектора Ф 62
 профессор (Б.М.Понтекорво)

Исполнители:
 профессор (Б.М.Понтекорво)
 Ст. инженер (Селиванов Г.И.)

1952 г.

25 September 1952

REPORT
Production of π^0 mesons in (n-p) and (n-d) collisions

Section leader
 Professor (B.M.Pontecorvo)

Executors:
 Professor (B.M.Pontecorvo)
 Engineer (Selivanov G.I.)

Fisico sperimentale ma anche fisico teorico

Dalle pagine di questi quaderni emerge la figura di bravo fisico sperimentale che coordina le attività del suo gruppo con competenza e grande rigore scientifico.

Gli esperimenti sull'interazione π -nucleone che in questi anni Pontecorvo realizza col suo gruppo al ciclotrone di Dubna sono certamente molto importanti per capire, almeno da un punto di vista fenomenologico, le interazioni forti.

Con questi esperimenti si conferma che il protone e il neutrone per quanto riguarda le interazioni forti non sono due particelle diverse ma sono la stessa particella in due stati diversi di un nuovo numero quantico chiamato spin isotopico.

Tuttavia l'interesse scientifico di Pontecorvo va ben oltre questi esperimenti, e molte delle sue riflessioni di questi anni riguardano ancora le interazioni deboli.

L'interazione debole aveva da sempre affascinato Pontecorvo, alla cui comprensione aveva già dato contributi fondamentali.

Interessantissimo è ciò che scrive a pagina 8! !

(primi di Novembre del 1950)

On the transformations of mesons

8

The Σ meson has a long life $\approx 10^{-9}$ sec, and is supposed to decay into $\pi^+ + \pi^+ + \pi^+$. If this is so, it must be concluded that Σ does not interact with nuclei, because, if the Σ interacts with nuclei, then the rate of the ~~disintegration~~ ^{disintegration} would be very fast. (through the interaction with nucleons of the vacuum). Let us suppose that it does not interact strongly. Since it is strongly produced, it must be produced as a decay product of a strongly interacting meson M . But this M then would decay into π quicker than in Σ . So there is a contradiction between the evidence of a strong interacting particle and its long lifetime. This contradiction, of course, is resolved if the strongly ^{strongly} particle is produced in pairs. So from the very fact that a) Σ mesons have a long life, b) that they are present in abundance, we can conclude that there are ~~mesons~~ ^{mesons} (not necessarily the Σ mesons) which are strongly produced in pairs. ^{Incidentally these considerations explain the fact that until present day cyclotron no other mesons than π mesons have been produced.} A consistent picture until now would be:

$\mu \rightarrow e + 2\nu$
 $\pi \rightarrow \mu + \nu$
 $\Sigma^+ = K = V^+ \rightarrow \begin{cases} \mu^+ + 2\nu \\ \mu^+ + \pi^+ + \pi^+ \\ \mu^+ + \pi^0 \end{cases}$

~~.....~~
 $V_{\text{light}} \rightarrow \pi^+ + \pi^+ \text{ or } \pi^+ + \mu^- ?$
 $V_{\text{heavy}} \rightarrow \rho + \pi$

$\mu \rightarrow e + \nu + \nu$

On the transformations of mesons -

The τ meson has a long life $\approx 10^{-9}$ sec, and is supposed to decay into $\pi^+ + \pi^+ + \pi^+$. If this is so, it must be concluded that τ does not interact with nuclei, because, if the τ interacts with nucleons then the rate of the disintegration would be very fast. (through the interaction with nucleons of the vacuum)

Let us suppose that it does not interact strongly. Since is strongly produced, it must produced as a decay product of a strongly interacting meson M . But this M then would decay into π quicker than in τ . So there is a contradiction between the existence of a strong interacting particle and his long lifetime. This contradiction, of course, is resolved if the strongly interacting particle is produced in pairs. (*) So from the very fact that a) τ mesons have a long life, b) that they are present in abundance, - we can conclude that there are mesons (not necessarily the τ mesons) which are strongly produced in pairs.

(incidentally these considerations explain the fact that until present day cyclotron no other mesons than π mesons have been produced.)

A consistent picture until now would be:

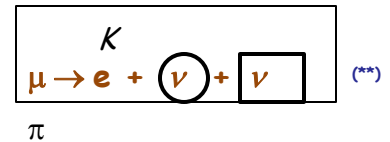
$\mu \rightarrow e + 2\nu$

$\pi \rightarrow \mu + \nu$

$\tau^+ = K = V^+ \rightarrow \begin{cases} \mu^+ + 2\nu \\ \mu^+ + \pi^+ + \pi^+ \\ \mu^+ + \pi^0 \end{cases}$

$V_{\text{light}} \rightarrow \pi^+ + \mu^+ \text{ or } \pi^+ + \mu^- ?$

$V_{\text{heavy}} \rightarrow \rho + \pi$



(*) a fine 1950 senza la nozione della stranezza, è necessaria una profonda intuizione per proporre che una produzione in coppia risolve la contraddizione.

(**) forse solo una coincidenza! Due righe prima scrive $\mu \rightarrow e + 2\nu$ mentre qui scrive $\mu \rightarrow e + \nu + \nu$ indicando i due neutrini con due diversi segni.

Due profonde intuizioni in una singola pagina ?!

Forse c'è una seconda intuizione geniale a pagina 8 !!

(primi di Novembre del 1950)

On the transformations of mesons

8

The Σ meson has a long life $\approx 10^{-9}$ sec, and is supposed to decay into $\pi^+ + \pi^+ + \pi^+$. If this is so, it must be concluded that Σ does not interact with nuclei, because, if the Σ interacts with nuclei, then the rate of the disintegration would be very fast. (through the interaction with nucleons of the vacuum). Let us suppose that it does not interact strongly. Since it is strongly produced, it must produced as a decay product of a strongly interacting meson M . But this M then would decay into π quicker than in Σ . So there is a contradiction between the evidence of a strong interacting particle, and its long lifetime. This contradiction, of course, is resolved if the strongly interacting particle is produced in pairs. So from the very fact that a) Σ mesons have a long life, b) that they are present in abundance, we can conclude that there are mesons (not necessarily Σ mesons) which are strongly produced in pairs. A consistent picture until now would be:

$\mu \rightarrow e + 2\nu$
 $\pi \rightarrow \mu + \nu$
 $\Sigma^+ = K = V^+ \rightarrow \begin{cases} \mu + 2\nu \\ \pi^+ + \pi^+ + \pi^+ \\ \pi^+ + \pi^0 \end{cases}$

$V_{light} \rightarrow \pi^+ + \mu^+ \text{ or } \pi^+ + \mu^- ?$
 $V_{heavy} \rightarrow p + \pi$

$V_{light} \Rightarrow \pi^+ + \mu^+ \text{ or } \pi^+ + \mu^- ?$
 $V_{heavy} \Rightarrow p + \pi$

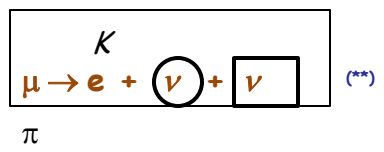
$\mu \rightarrow e + \nu + \bar{\nu}$

Other mesons have been produced.

On the transformations of mesons -

The τ meson has a long life $\approx 10^{-9}$ sec, and is supposed to decay into $\pi^+ + \pi^+ + \pi^+$. If this is so, it must be concluded that τ does not interact with nuclei, because, if the τ interacts with nucleons then the rate of the disintegration would be very fast. (through the interaction with nucleons of the vacuum). Let us suppose that it does not interact strongly. Since is strongly produced, it must produced as a decay product of a strongly interacting meson M . But this M then would decay into π quicker than in τ . So there is a contradiction between the existence of a strong interacting particle and his long lifetime. This contradiction, of course, is resolved if the strongly interacting particle is produced in pairs. (*) So from the very fact that a) τ mesons have a long life, b) that they are present in abundance, - we can conclude that there are mesons (not necessarily the τ mesons) which are strongly produced in pairs. (incidentally these considerations explain the fact that until present day cyclotron no other mesons than π mesons have been produced.) A consistent picture until now would be:

$$\begin{aligned} \mu &\rightarrow e + 2\nu \\ \pi &\rightarrow \mu + \nu \\ \tau^+ = K = V^+ &\rightarrow \begin{cases} \mu + 2\nu \\ \mu^+ + \pi^+ + \pi^- \\ \mu^+ + \pi^0 \end{cases} \\ V_{light} &\rightarrow \pi^+ + \mu^+ \text{ or } \pi^+ + \mu^- ? \\ V_{heavy} &\rightarrow p + \pi \end{aligned}$$



(*) a fine 1950 senza la nozione della stranezza, è necessaria una profonda intuizione per proporre che una produzione in coppia risolve la contraddizione.
 (**) forse solo una coincidenza! Due righe prima scrive $\mu \rightarrow e + 2\nu$ mentre qui scrive $\mu \rightarrow e + \nu + \bar{\nu}$ indicando i due neutrini con due diversi segni.
 Due profonde intuizioni in una singola pagina ?!

Sul quaderno si trova un'altra pagina (la 76) interessantissima !!

Questa pagina è stata scritta tra il 25 dicembre 1951 e il 30 gennaio 1952.

Appendix

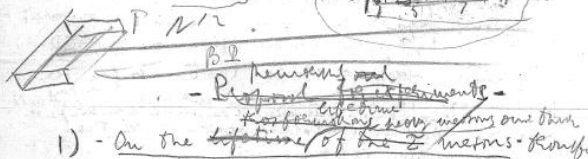
Appendix for the determination of a search for a stable state of H^+ (Hecht & Vay).

Appendix The Experiment

The concept for the experiment is to produce a beam of H^+ stable ions in the emission of β particles. The experiment planned is to make a beam of H^+ ions in the emission of β particles. The aim with emission of β particles of a 20 MeV. The apparatus consist of a β particles consist of 3-counter in coincidence.

Problems - Future work -

- 1) - Production of mesons in nuclear
- 2) - Interaction of mesons with nuclei
- 3) - Production of N^{12} in stars and in H^+



- 1) - On the lifetime of the τ mesons - Kowalew
- 2) $\tau \rightarrow$ experiment

Remarks and

Proposal for experiments -

- 1) - On the lifetime transformations lifetime of the τ mesons heavy mesons and their transformation -
- 2) $\tau \rightarrow$ experiment

β or the charge symmetry - On the charge symmetry -

A. Alex. - Observations

In the course of this year several remarks on proposed experiments were made in the 62 group, of which it is possible to mention some.

- 1) At the seminaire ~~a method~~ was discussed ~~to solve~~ the problem of the detection of free neutrons, i.e. of the detection of neutrons which is not connected with the act of a β disintegration (like in the classical experiment of Leipunski). The conclusion is that such possibility is not too far from present day facilities. A short report on this subject was written.
- 2) ~~Lifetime of τ mesons~~ - ~~Proposed experiment on τ mesons~~. In photographic plates it was observed τ

(3) Lifetime etc

(4) ~~On the charge symmetry hypothesis - A discussion~~

c 10^{16} Km (*)

$Cl^{37} + \nu \rightarrow Ar^{37} + e$

3) On the charge symmetry - On the charge symmetry

A. Alex. -

Observations

In the course of this year several remarks or proposed experiments were made in the 62 group, of which it is possible to mention some.

Neutrino -

1) At the seminaire ~~a method~~ was discussed ~~in-rele~~ the problem of the detection of free neutrons, i.e. of a detection of neutrino, a method which is not connected with the act of a β disintegration (like in the classical experiment of Leipunski). The conclusion is that such possibility is not too far from present day facilities. A short report on this subject was written.

(2) Lifetime of τ mes Heavy mesons - Possible experiment on τ meson.

In photographic plates it was observed τ

(3) Lifetime etc.

(4) On the charge symmetry hypothesis - A discussion

(*) H. Bethe and R. Peierls in Nature 133, 532 (07 April 1934) avevano valutato un limite superiore per la probabilità per il neutrino di interagire con la materia solida e scrivevano: "For an (neutrino) energy of $2 \cdot 3 \times 10^6$ volts... $\sigma < 10^{-44}$ cm² (corresponding to a penetrating power of 10^{16} Km in solid matter) It is therefore absolutely impossible to observe process of this kind with neutrons created in nuclear transformation."

Come rivelare il neutrino

Immagino che Pontecorvo quando scrive, alla fine del 1951, nell'angolo in alto a destra della pagina 76 del quaderno



stia valutando il flusso di neutrini e la quantità di Cloro necessaria a rivelare questa particella così elusiva da poter attraversare 10^{16} Km di materia solida senza interagire !

Già alla fine del 1951 Pontecorvo pensa di essere in grado di fare l'esperimento Cloro/Argon per rivelare finalmente il neutrino.

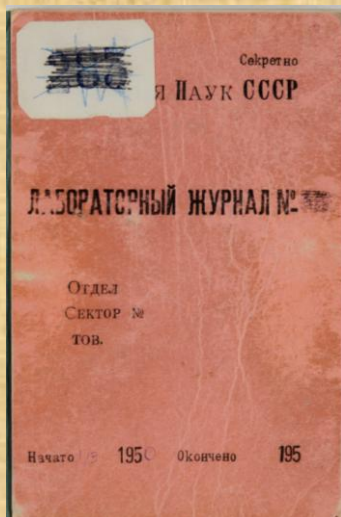
A handwritten note in blue ink on a piece of paper. The text reads: "like the classical experiment of Lavoisier... The conclusion is that such possibility is not too far from present day facilities, a report on this subject was written".

Sarebbe veramente interessante trovare questo **"short report"** per sapere come e dove Pontecorvo pensava che fosse possibile per lui fare l'esperimento in Russia. Sfortunatamente questa possibilità non si concretizzò mai; forse perché non aveva a disposizione una miniera o una caverna sufficientemente grande e profonda da provare a rivelare il flusso di neutrini solari o forse semplicemente perché, come testimonia il fisico russo S.S.Gershtein, gli era negato l'accesso ai reattori nucleari che, come aveva scritto nei suoi articoli, sembravano essere una sorgente di neutrini più promettente di quella del sole per fare l'esperimento.

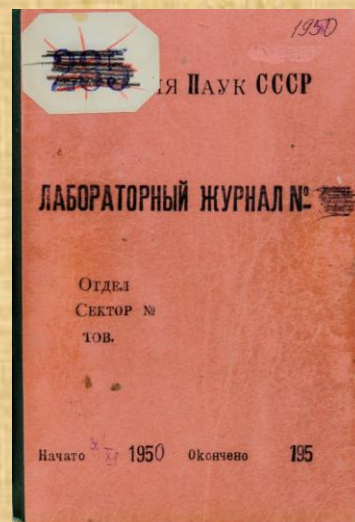


Sognando di rivelare i neutrini solari !
di Misha Bilenky

Conclusioni



01/11/1950 — ≤ 30/11/1950
14/09/1951 — ≥ 24/03/1952

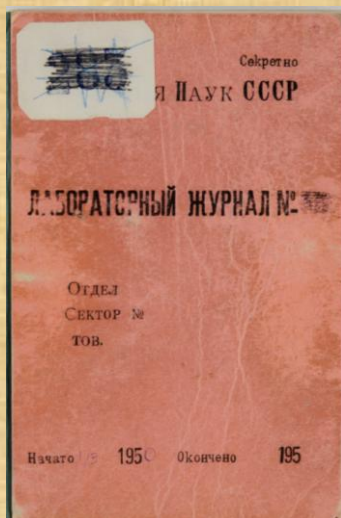


30/11/1950 — ≥ 18/07/1951

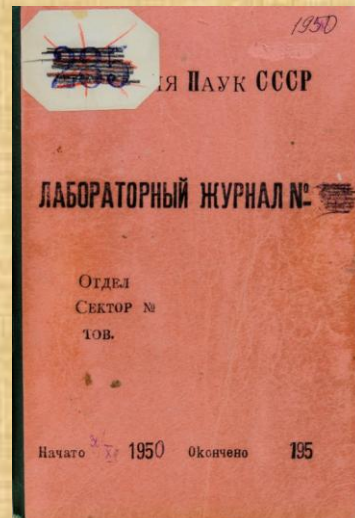
Attraverso la preziosa documentazione di questi due quaderni inediti (che molto spesso Pontecorvo usa soltanto come un "brogliaccio" scrivendoci in modo disordinato e senza un particolare ordine logico numeri, formule e grafici, forse solo per fissare i pensieri e le idee che gli turbinano nella mente) è stato possibile ricostruire il lavoro sperimentale, gli interessi scientifici e anche alcune grandi intuizioni che questo grande scienziato ha avuto durante il suo primo anno e mezzo di permanenza in Russia. Questi due documenti sono particolarmente interessanti anche perché spesso sono stati sollevati dubbi su quale sia stata la vera attività scientifica di Bruno Pontecorvo nei primi anni della sua permanenza in Russia.

In questi quaderni c'è la conferma più evidente che Bruno Maximovich Pontecorvo non abbia mai lavorato né contribuito alla realizzazione della bomba atomica russa ma abbia solo fatto ricerca di base in fisica delle particelle elementari.

Ringraziamenti



01/11/1950 — ≤ 30/11/1950
14/09/1951 — ≥ 24/03/1952



30/11/1950 — ≥ 18/07/1951

I nostri più sinceri ringraziamenti vanno a Gil Pontecorvo, il figlio maggiore di Bruno, che ci ha gentilmente concesso la possibilità di consultare questi documenti inediti. Un ringraziamento particolare ad Elena Volterrani per averci aiutato in ogni fase di questo lavoro, dal ritrovamento dei quaderni alla loro lettura. Ringraziamo Vincenzo Cavasinni e Marco Maria Massai che si stanno prodigando per la realizzazione di un centro Pontecorvo presso il Dipartimento di fisica di Pisa, per il supporto e l'incoraggiamento che ci hanno dato. Ringraziamo poi xxxx per averci tradotto dal alcune pagine scritte in russo che si trovano su questi quaderni. Ringraziamo infine Bruno Sereni e Antonio D'Agnelli dell'Ufficio Documentazione dell'Università di Pisa per l'accurato lavoro fatto nel fotografare in alta risoluzione tutte le pagine di questi quaderni.

Grazie dell'attenzione