

Experimental techniques in high-energy nuclear and particle physics

“Dottorato di Ricerca in Ingegneria dell’Informazione”

LECTURE 6a.

Tracking at LHC

Prof. Rino Castaldi

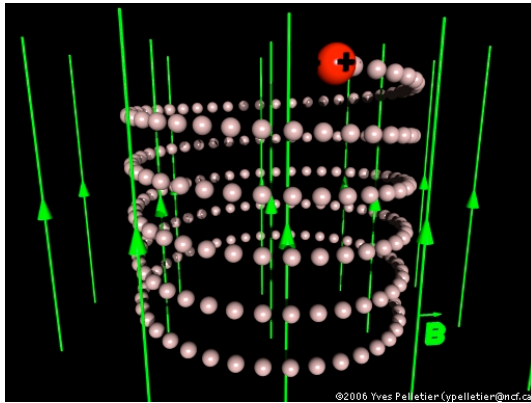
INFN-Pisa

rino.castaldi@pi.infn.it

Tracking: why

At hadron colliders the challenging aim is the full reconstruction of the events produced in the interaction under study. Therefore primary goals are:

- ✓ reconstruct the trajectories ("tracking") of charged particles and measure their momenta



Most common case: in a solenoidal uniform magnetic field the Lorentz force

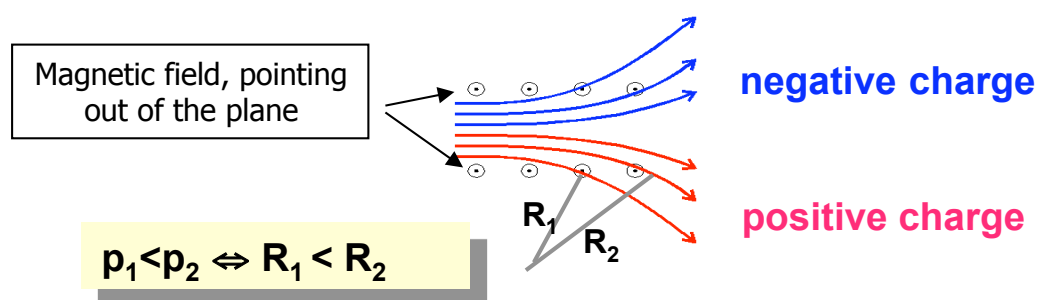
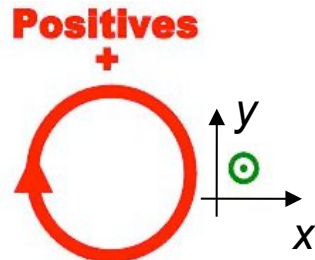
$$\vec{F} = \frac{d\vec{p}}{dt} = q\vec{E} + q(\vec{v} \times \vec{B})$$

induce charged particles to follow a helicoidal path:

- describe circles in the transverse plane
- move uniformly along the magnetic field direction

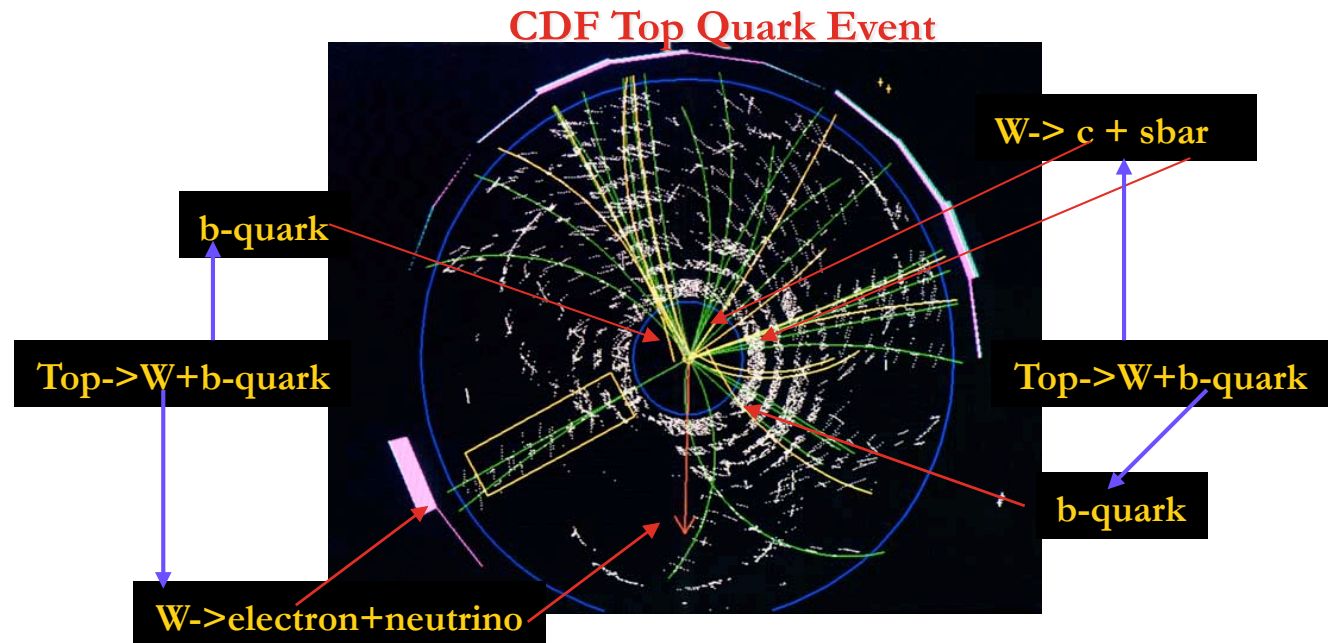
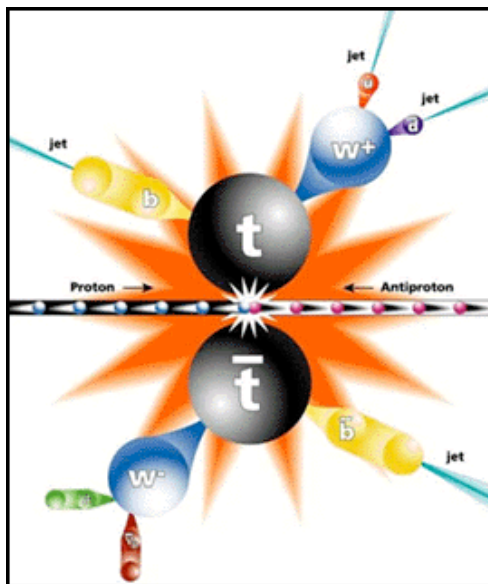
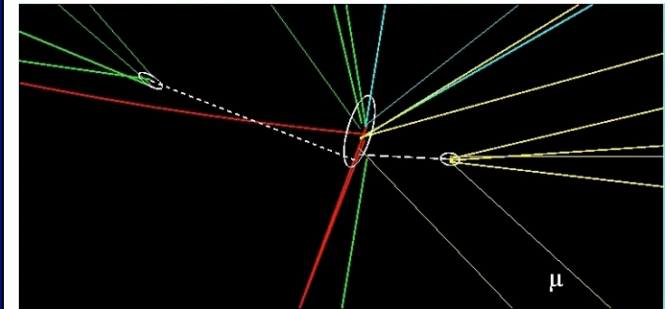
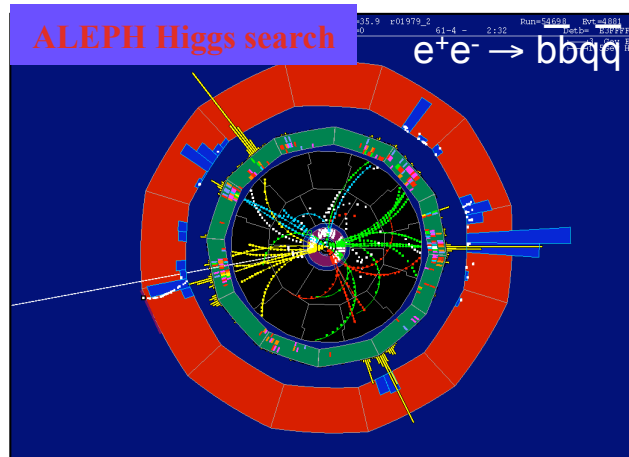
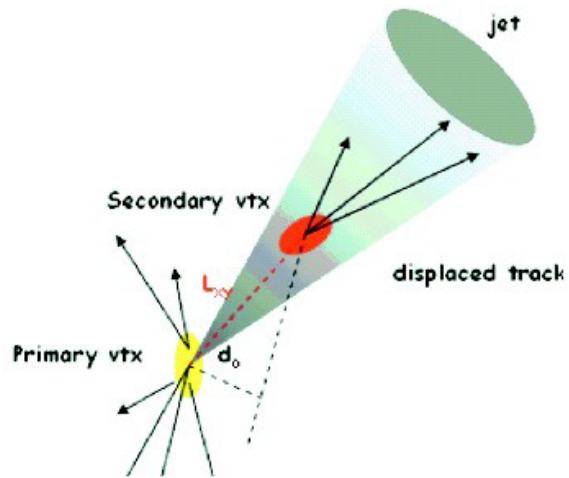
$$p_T(\text{GeV}) = 0.3 B(\text{T}) R(\text{m})$$

- ✓ identify the sign of the charge



Tracking: why

✓ reconstruct the primary and secondary vertices of the interaction
(at LHC with large pile-up of events in the same bunch crossing !)



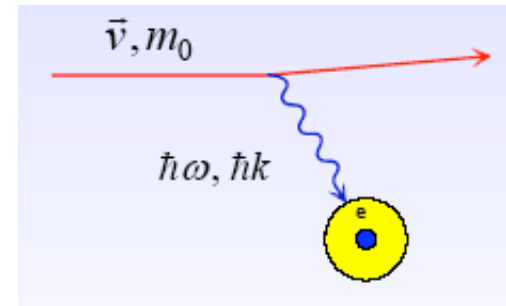
Tracking: a real challenge at LHC

- Tracking at LHC is a very complex procedure due to the high track density. It needs specific implementation adapted to the detector type and geometry
- Precise and efficient detector modules are required to measure where the particle crossed the module
- Fast and radiation hard detectors and electronics are needed
- Track reconstruction requires specific software implementation:
 - track finding (pattern recognition)
 - estimation of track parameters (fitting)
- Precise alignment of detector modules is a prerequisite for efficient tracking

Tracking: how

measurable signals occur via the interaction of particles with the detector material.

Dominant interaction is due to the coulomb interactions with the atomic electrons of the detector.

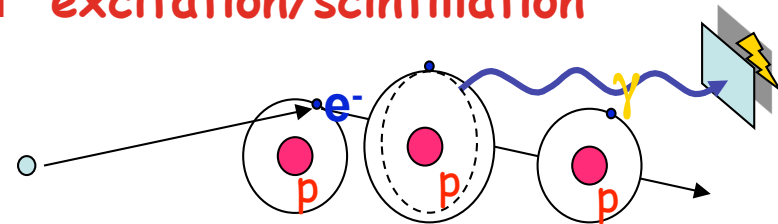
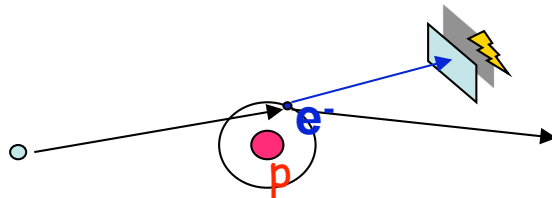


Depending on the $\hbar\omega$ value we may have:

□ ionization

or

□ excitation/scintillation



Ionization and excitation of atomic electrons in matter are the most common processes and allow to build precise tracking detectors .

Ionization: the Bethe-Bloch formula

$$\left\langle \frac{dE}{dx} \right\rangle = -4\pi N_A r_e^2 m_e c^2 z^2 \frac{Z}{A} \frac{1}{\beta^2} \left[\frac{1}{2} \ln \frac{2m_e c^2 \gamma^2 \beta^2}{I^2} T^{\max} - \beta^2 - \frac{\delta}{2} \right]$$

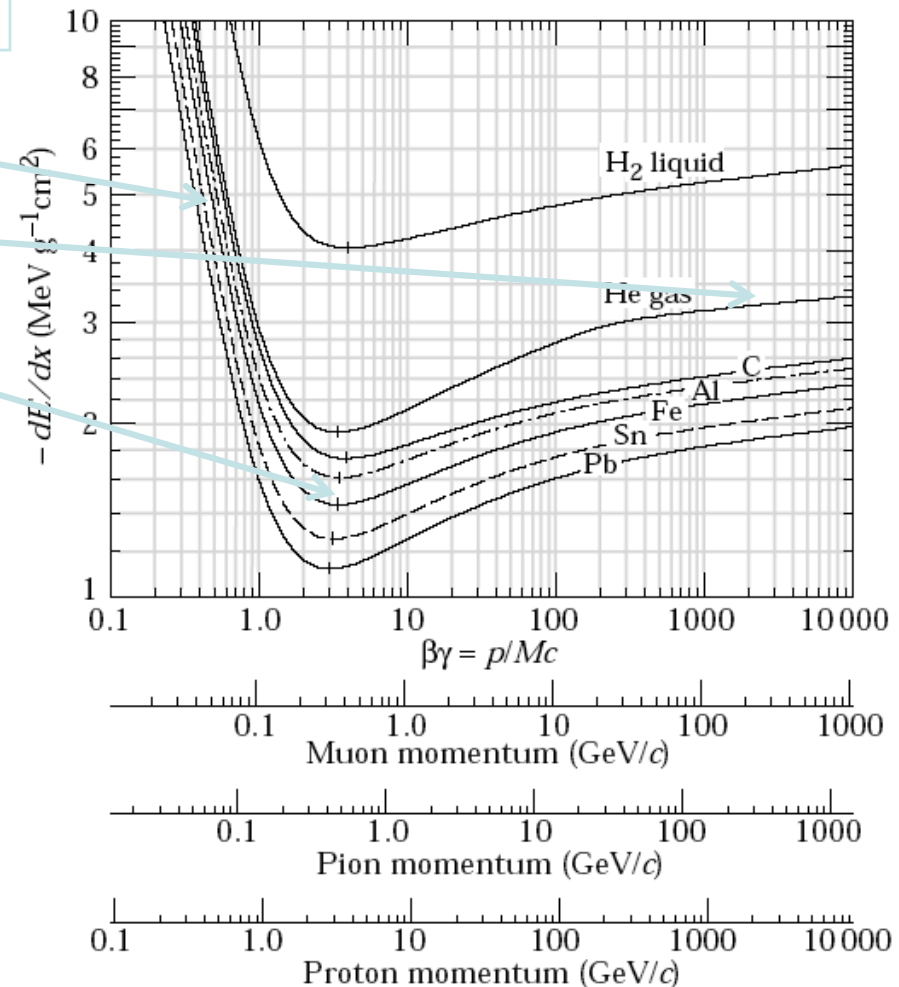
$$T_{\max} \approx 2m_e c^2 \beta^2 \gamma^2 \quad \left[-\frac{dE}{dx} \approx K q^2 \frac{Z}{A \beta^2} \left[\ln \frac{2m_e c^2 \beta^2 \gamma^2}{I^2} - \beta^2 \right] \right]$$

Characterized by:

- a fall off at low energy $\sim 1/\beta^2$
- a relativistic rise $\sim \ln \beta\gamma$
- a minimum at $\beta\gamma \approx 3$
- depends only on $\beta\gamma$ not on m

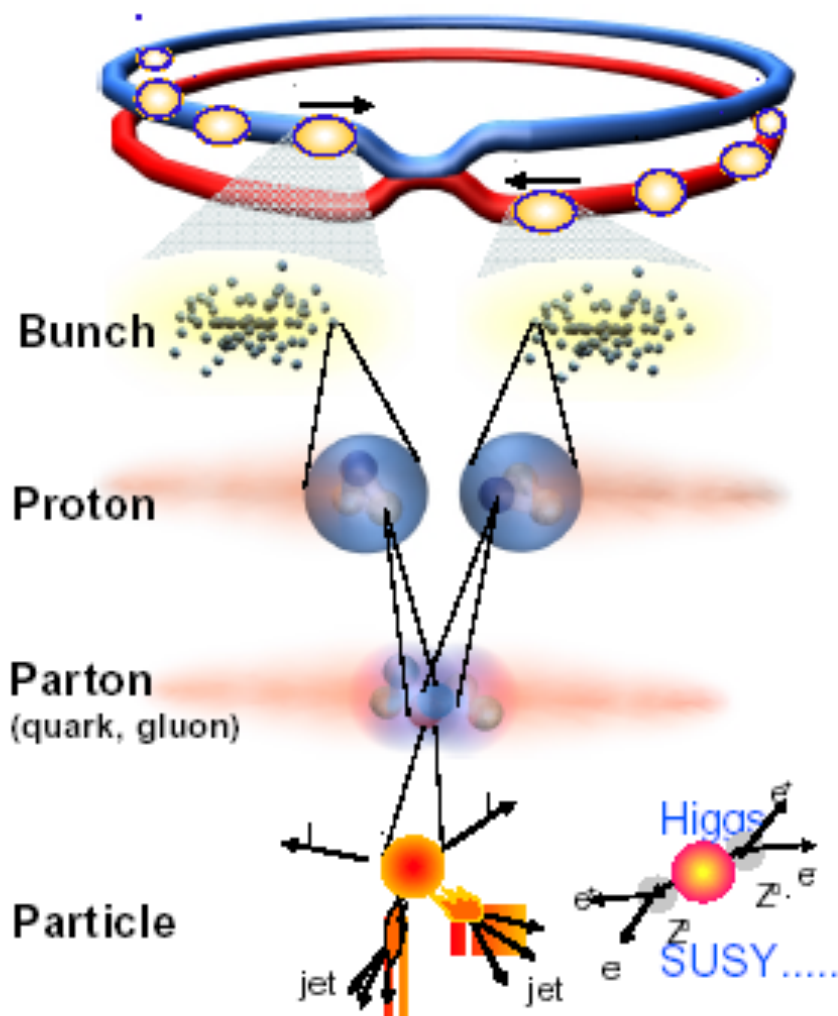
High energy charged particles lose energy **slowly** in material due to ionization leaving tracks as they pass (For $Z \approx 0.5A$ at $\beta\gamma \approx 3$ $1/\rho dE/dx \approx 1.4 \text{ MeV cm}^2/\text{g}$)

→ many kinds of tracking detectors can be done !



LHC: a very hostile environment for tracking
(high event rate, high multiplicity of charged tracks, high radiation flux)

Collisions at LHC



Proton-Proton	2835 bunch/beam
Protons/bunch	10^{11}
Beam energy	7 TeV (7×10^{12} eV)
Luminosity	$10^{34} \text{ cm}^{-2} \text{ s}^{-1}$

Crossing rate 40 MHz

Collisions rate $\approx 10^7 - 10^9$ Hz

New physics rate $\approx .00001$ Hz

Event selection:
1 in 10,000,000,000,000

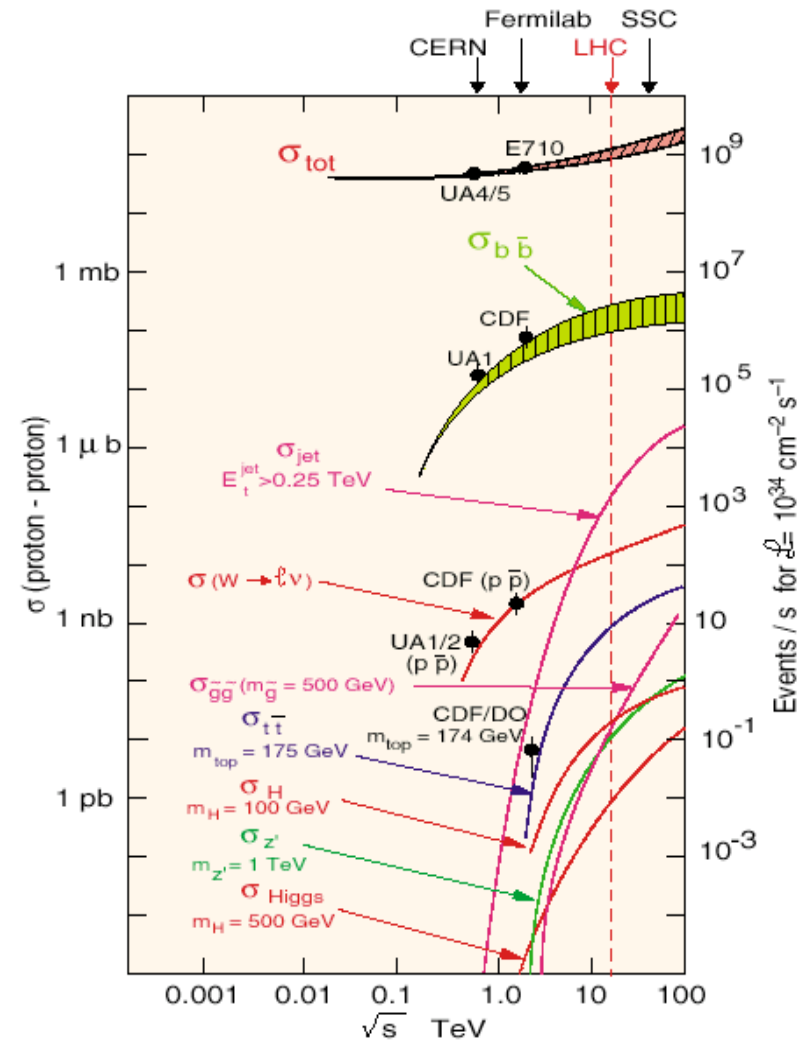
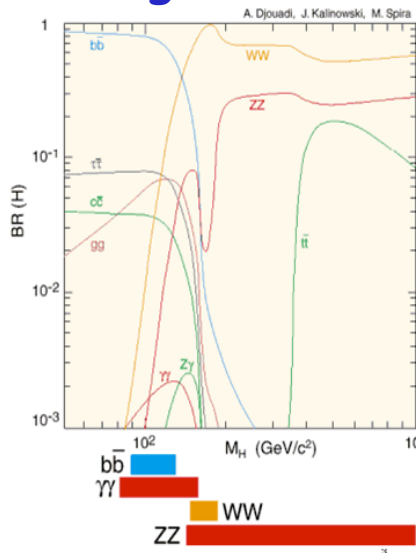


Signal and background $\rightarrow \mathcal{L}=10^{34} \text{ cm}^{-2}\text{s}^{-1}$

Cross sections for various physics processes vary over many orders of magnitude

Higgs ($600 \text{ GeV}/c^2$): $1 \text{ pb} @ 10^{34} \rightarrow 10^{-2} \text{ Hz}$
 Higgs ($100 \text{ GeV}/c^2$): $10 \text{ pb} @ 10^{34} \rightarrow 0.1 \text{ Hz}$
 $t \bar{t}$ production: $\rightarrow 10 \text{ Hz}$
 $W \rightarrow \ell \nu$: $\rightarrow 10^2 \text{ Hz}$
 Inelastic: $\rightarrow 10^9 \text{ Hz}$

Selection needed: $1:10^{10-11}$
Before branching fractions...



\Rightarrow Needle in a Hay Stack

pp cross-sections and minimum bias

of interactions/crossing:

Interactions/s:

$$\text{Lum} = 10^{34} \text{ cm}^{-2}\text{s}^{-1} = 10^7 \text{ mb}^{-1}\text{Hz}$$

$$\sigma_{\text{inel}}(\text{pp}) = 80 \text{ mb}$$

$$\text{Interaction Rate, } R = 8 \times 10^8 \text{ Hz}$$

Events/beam crossing:

$$\Delta t = 25 \text{ ns} = 2.5 \times 10^{-8} \text{ s}$$

$$\text{Interactions/crossing} = 20$$

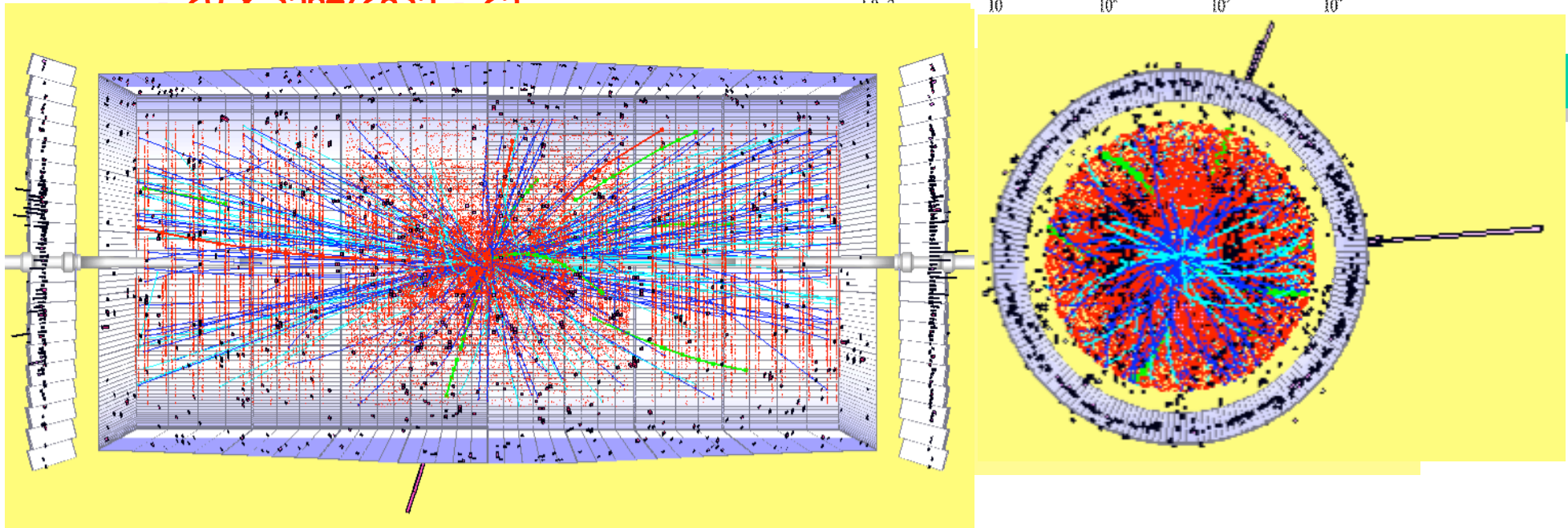
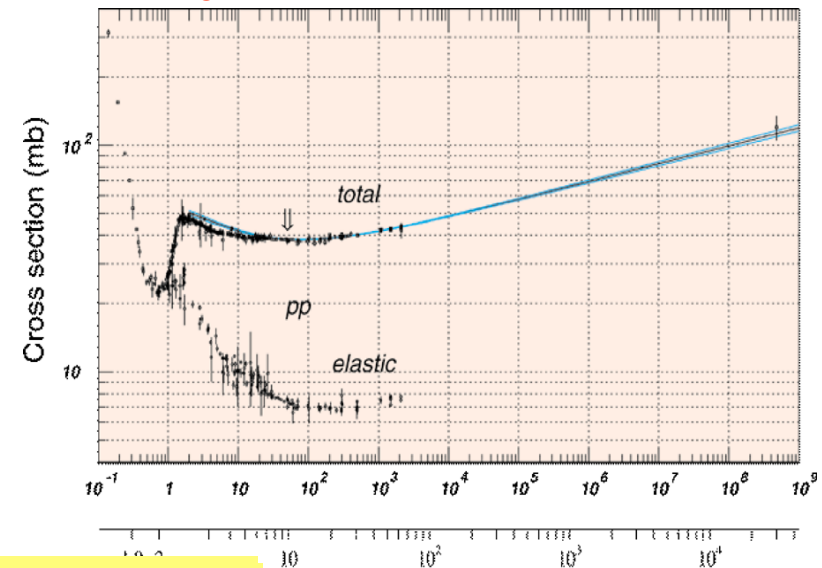
Not all p bunches are full

2835 out of 3564 only

Interactions/"active" crossing =

$$= 20 \times 3564 / 2835 = 25$$

$$\sigma_{\text{inel}}(\text{pp}) \approx 80 \text{ mb} @ 14 \text{ TeV}$$



Impact on detector design

LHC detectors must have fast response

Otherwise will integrate over many bunch crossings

→ large “pile-up”

Typical response time : 20-50 ns

→ integrate over 1-2 bunch crossings

→ pile-up of 25-50 min-bias

→ very challenging readout electronics

LHC detectors must be highly granular

Minimize probability that pile-up particles be in the same detector element as interesting object

→ large number of electronic channels

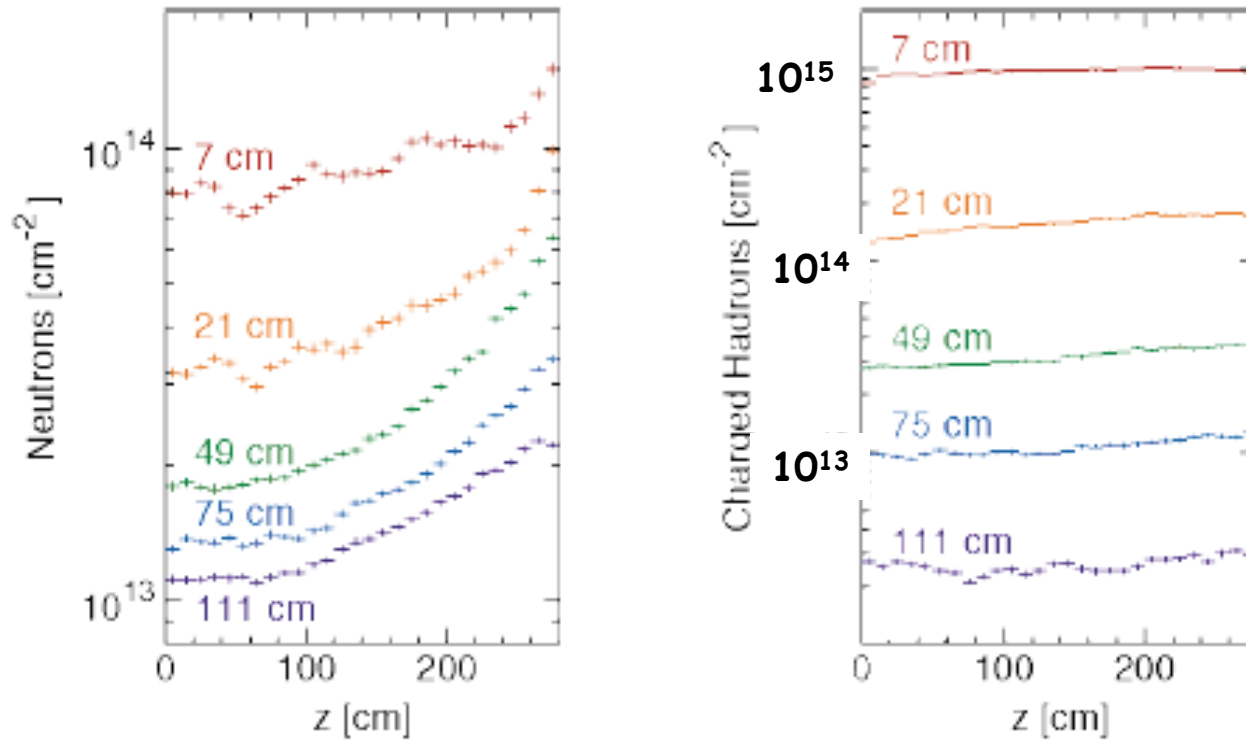
→ high cost

LHC detectors must be radiation resistant:

high flux of particles from pp collisions

→ high radiation environment

Radiation environment at the LHC




Expected particle fluences for the silicon detector inner layers in CMS integrated over 10 years as a function of the distance from the vertex point and for various radii.

Left: neutrons


Right: charged hadrons

Evian, March 1992: first meeting on experimental program at LHC
Evian, November 2009: more than 17 years later ready to start....

Towards the LHC Experimental Programme
General Meeting on LHC Physics & Detectors
 5-8 March 1992
 Evian-les-Bains, France



Expressions of Interest	
The Ascot detector at the LHC P. Norton (Rutherford-Appleton Laboratory).....	8
CMS : a compact solenoidal detector for LHC M. Della Negra (CERN) & J-C. Lottin (DAPNIA, CEN-Saclay)	9
EAGLE : Experiment for Accurate Gamma, Lepton and Energy measurements P. Jenni (CERN).....	11
L3 detector upgrade for LHC : The Extended L3 Collaboration S.C.C. Ting (MIT)	13
An LHC collider Beauty experiment for CP-violation measurements P. Schlein (UCLA).....	15
Measurement of CP-violation in B-decays using an LHC extracted beam : The LHB Collaboration G. Carboni, Pisa.....	16
A study of CP violation in B-meson decays using a gas jet at LHC T. Nakada (PSI).....	18
Neutrino physics at LHC K. Winter (CERN).....	19
A neutrino experiment at LHC F. Vannucci (Paris)	20
A dedicated heavy ion experiment at the LHC J. Schukraft (CERN).....	21
A feasibility study of using DELPHI as a detector for heavy ion collisions at LHC G. Jariskog (Lund).....	23
A heavy ion experiment with CMS at LHC L. Ramello (Turin).....	25



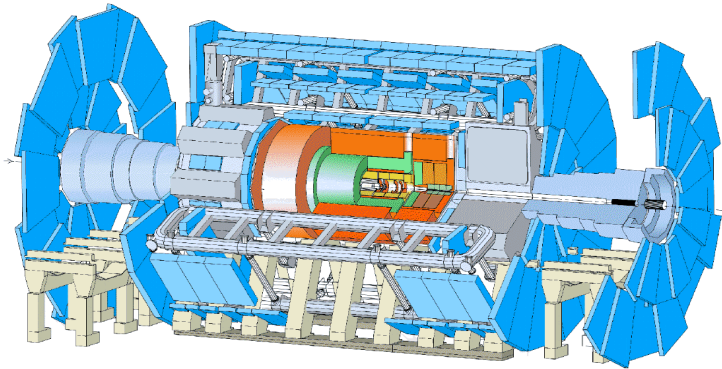
Hadron Collider Physics Symposium
Evian, France
16 – 20 November, 2009



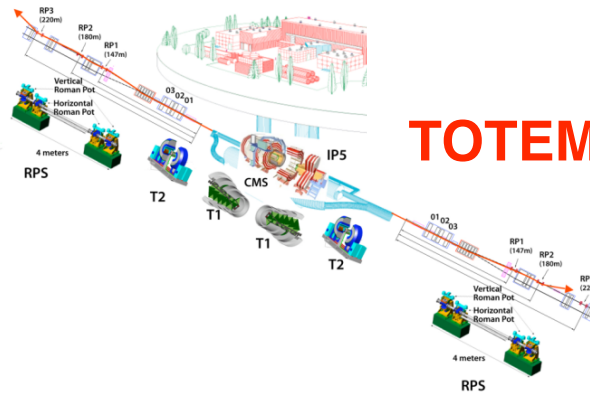
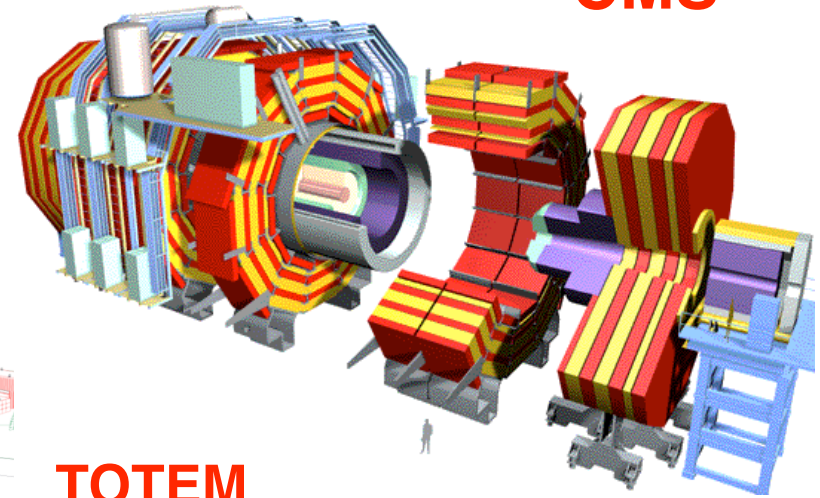
**ATLAS, CMS, ALICE,
 LHCb, TOTEM, LHCf**

The LHC Detectors

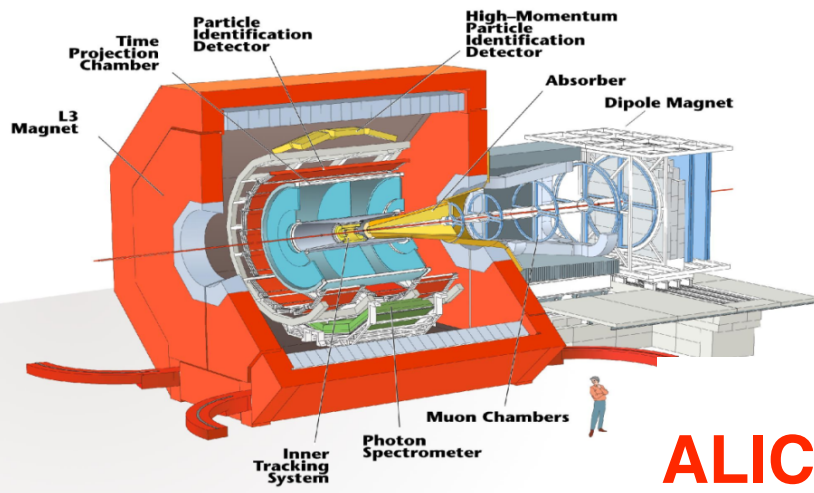
ATLAS



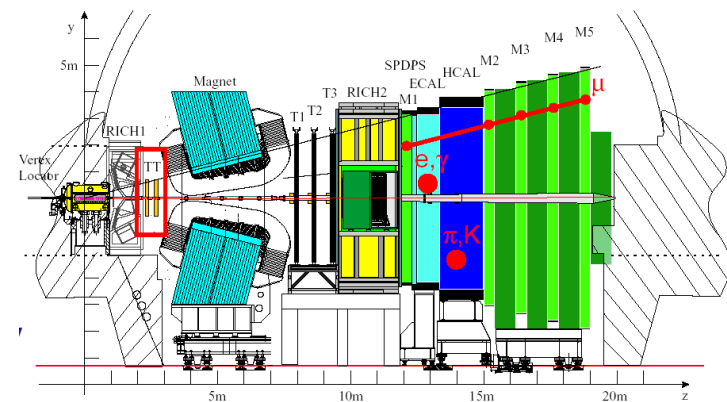
CMS



TOTEM



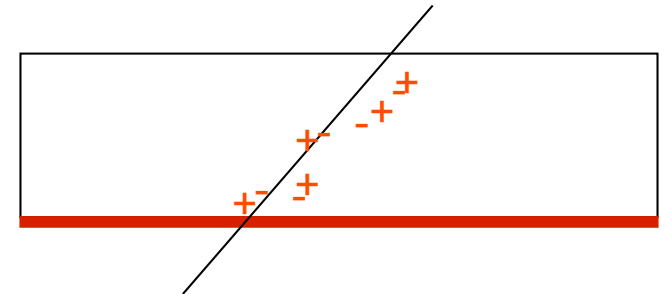
ALICE



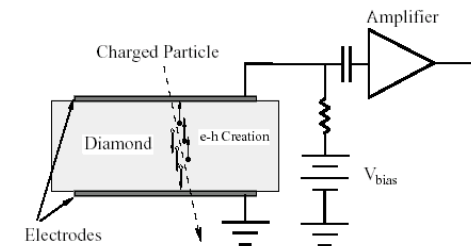
LHCb

Tracking Detectors

Charged particles crossing a material lose energy by ionizing (and exciting) atoms and thus leaving along their path a trace of electron-ion pairs in gases and liquids and electron-hole pairs in solids.



Measurable electronics signals can be induced by the charges produced in this way and can be read by dedicated electronics



□ In solid state detectors the charges produced by the ionization due to the incoming particle are sufficient to provide a measurable signal.

□ In gas detectors the charges produced by the primary ionization due to the incoming particle need amplification in order to provide a measurable signal.

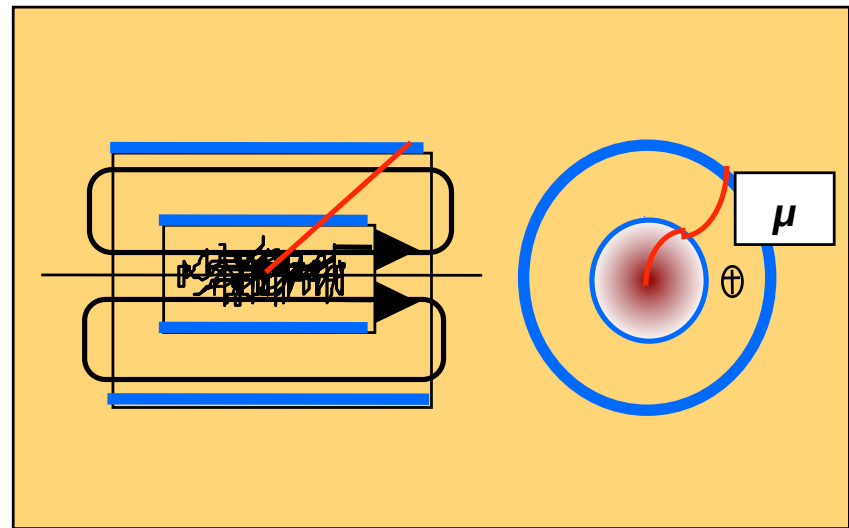
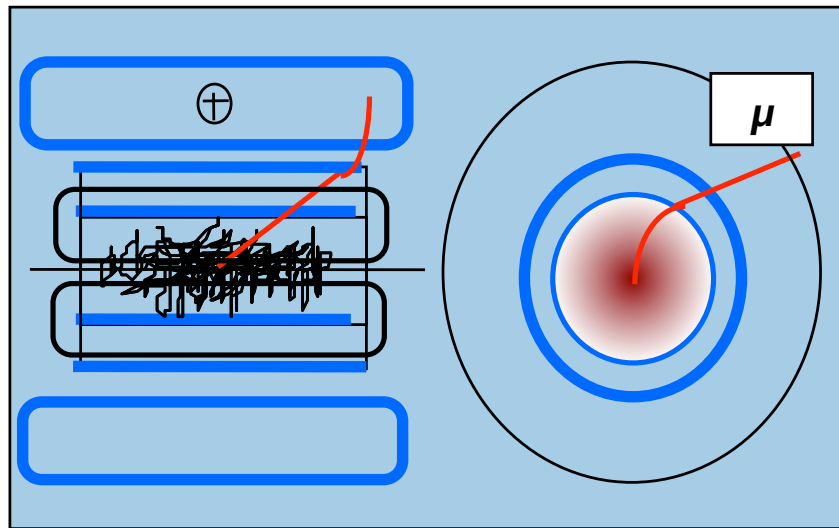
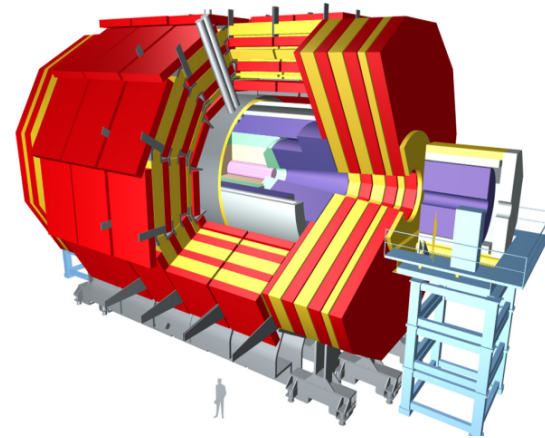
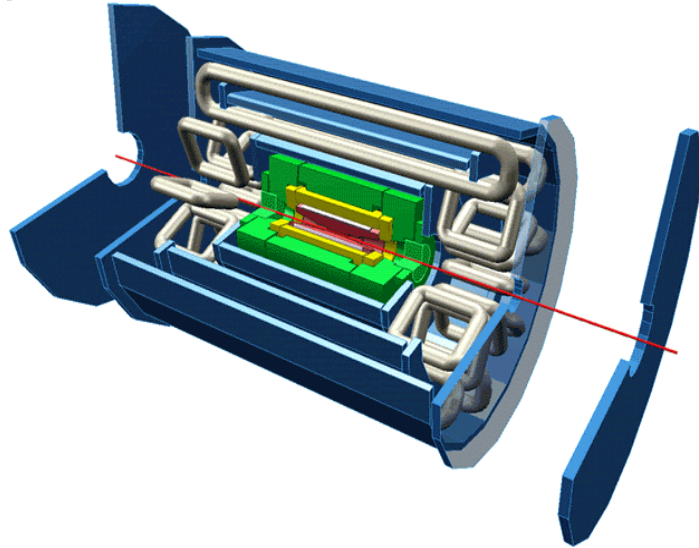
- Mean (most probable) energy loss: 116 (78) keV for 300 μ m Si thickness
- 3.6 eV to create an e-h pair
 \Rightarrow 72 e-h/ μ m (mean)
 \Rightarrow 108 e-h/ μ m (most probable)
- Mean charge (300 μ m Si)
 \approx 22000 e \approx 3.6 fC

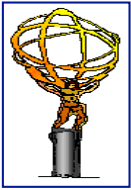
GAS)	Helium	Argon	Xenon	CH ₄	DME
dE/ dx (keV/ cm)	0.32	2.4	6.7	1.5	3.9
<n>(ion-pair/ cm)	5.9	29	44	16	55

Two general-purpose pp detectors with different strategies

ATLAS A Toroidal LHC Apparatus

CMS C ompact M uon Solenoid





Different Strategies....

2.3m x 5.3m Solenoid ~ 2 Tesla Field

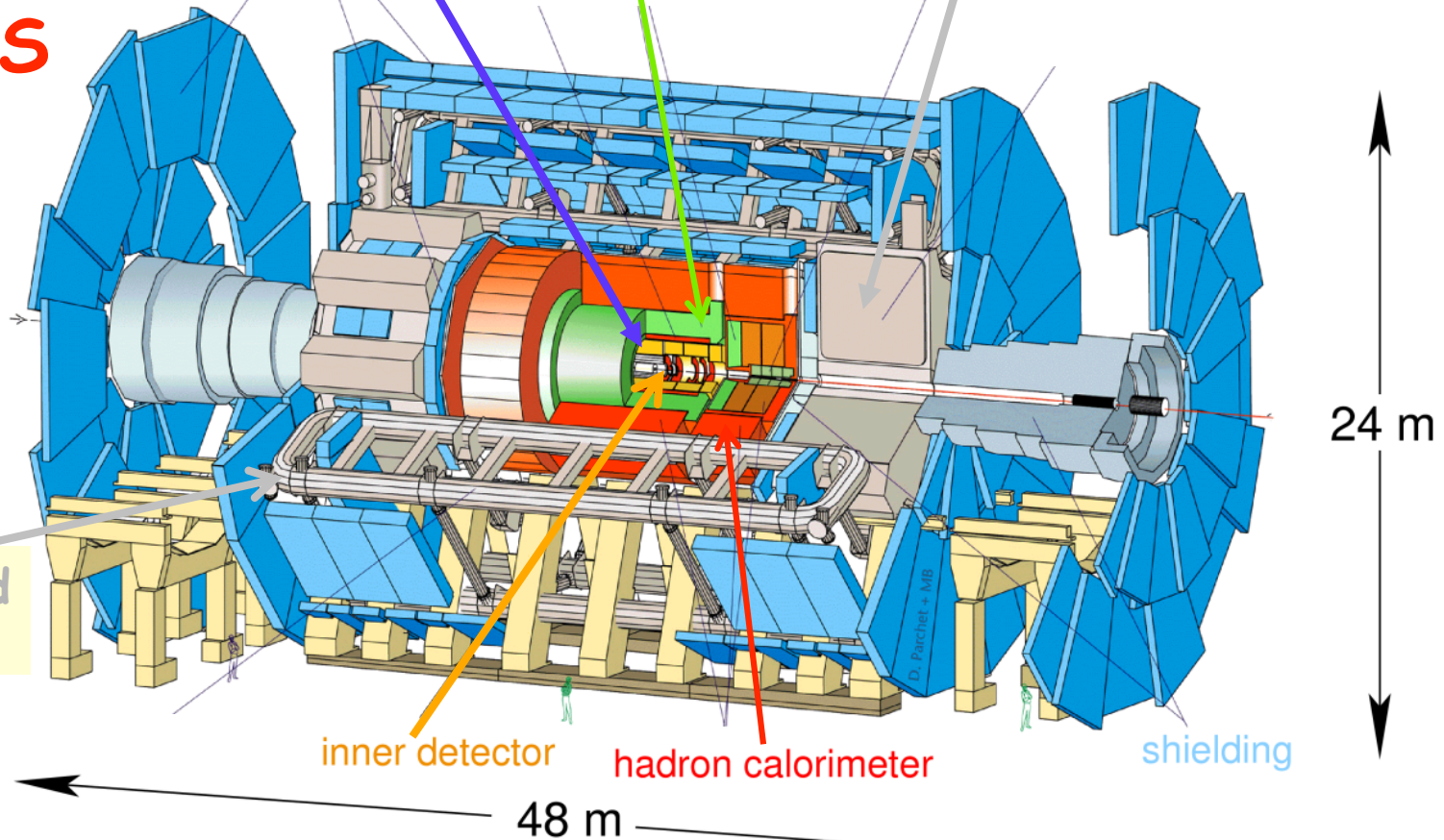
End Cap Toroid

muon detectors

electromagnetic calorimeters

forward calorimeters

ATLAS



Barrel Toroid
 $\langle B \rangle = 0.4 \text{ T}$

inner detector

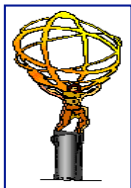
hadron calorimeter

shielding

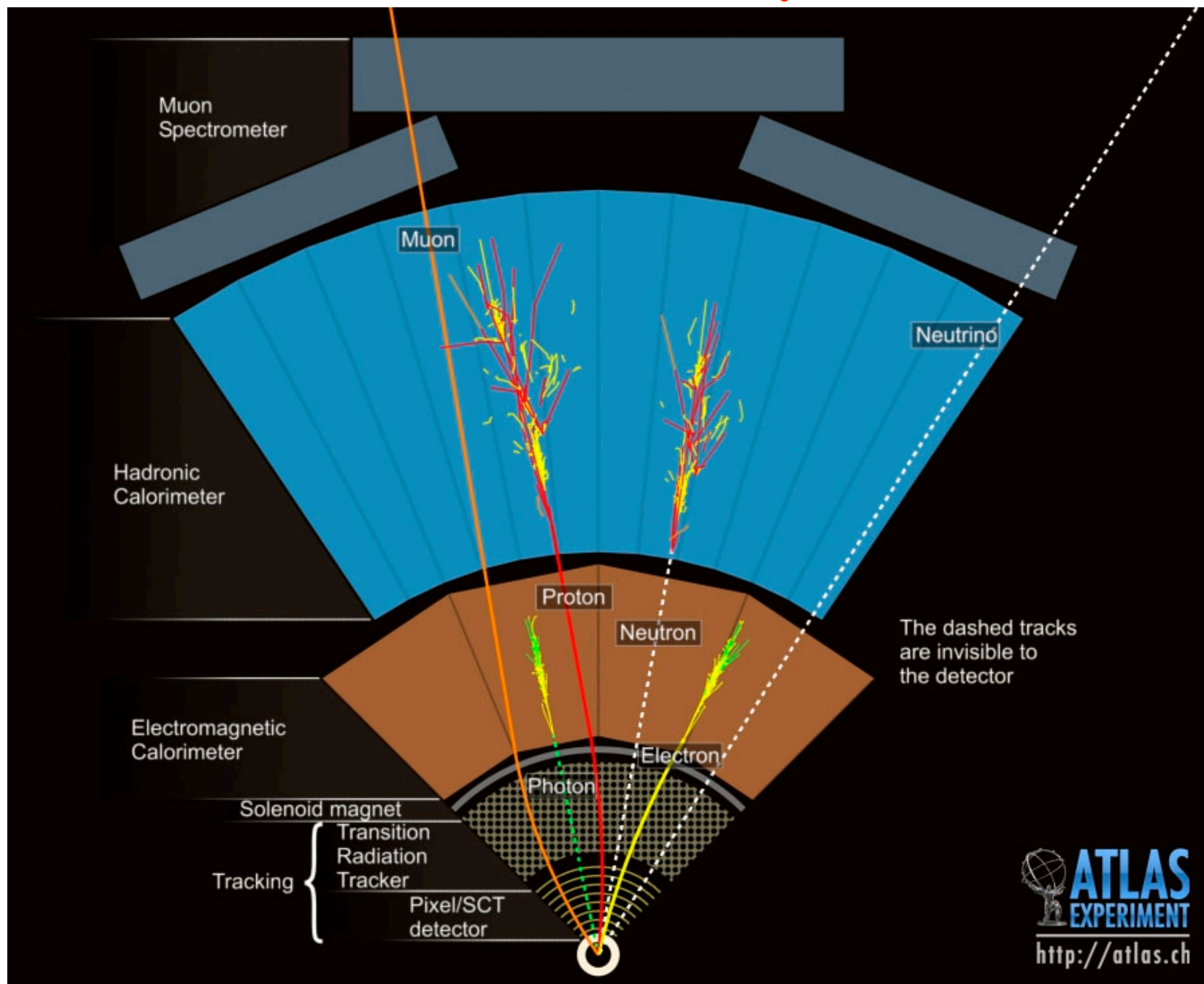
48 m

24 m

48m Long, 24m Diameter, 7'000 Ton Detector



Transverse slice of ATLAS





Different Strategies...

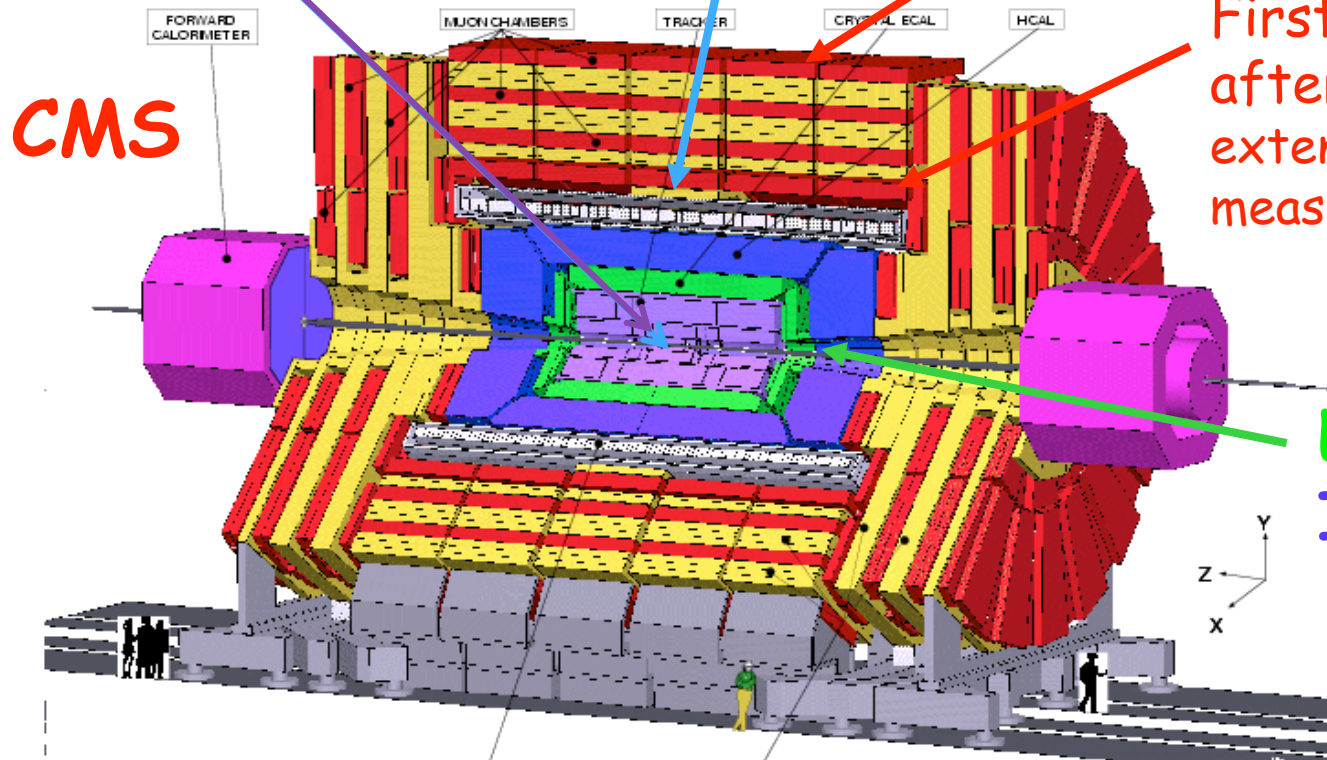
13m x 6m Solenoid: 4 Tesla Field

Tracking up to $\eta \sim 2.4$

Muon system in return yoke

First muon chamber just after solenoid \rightarrow extended lever arm for p_T measurement

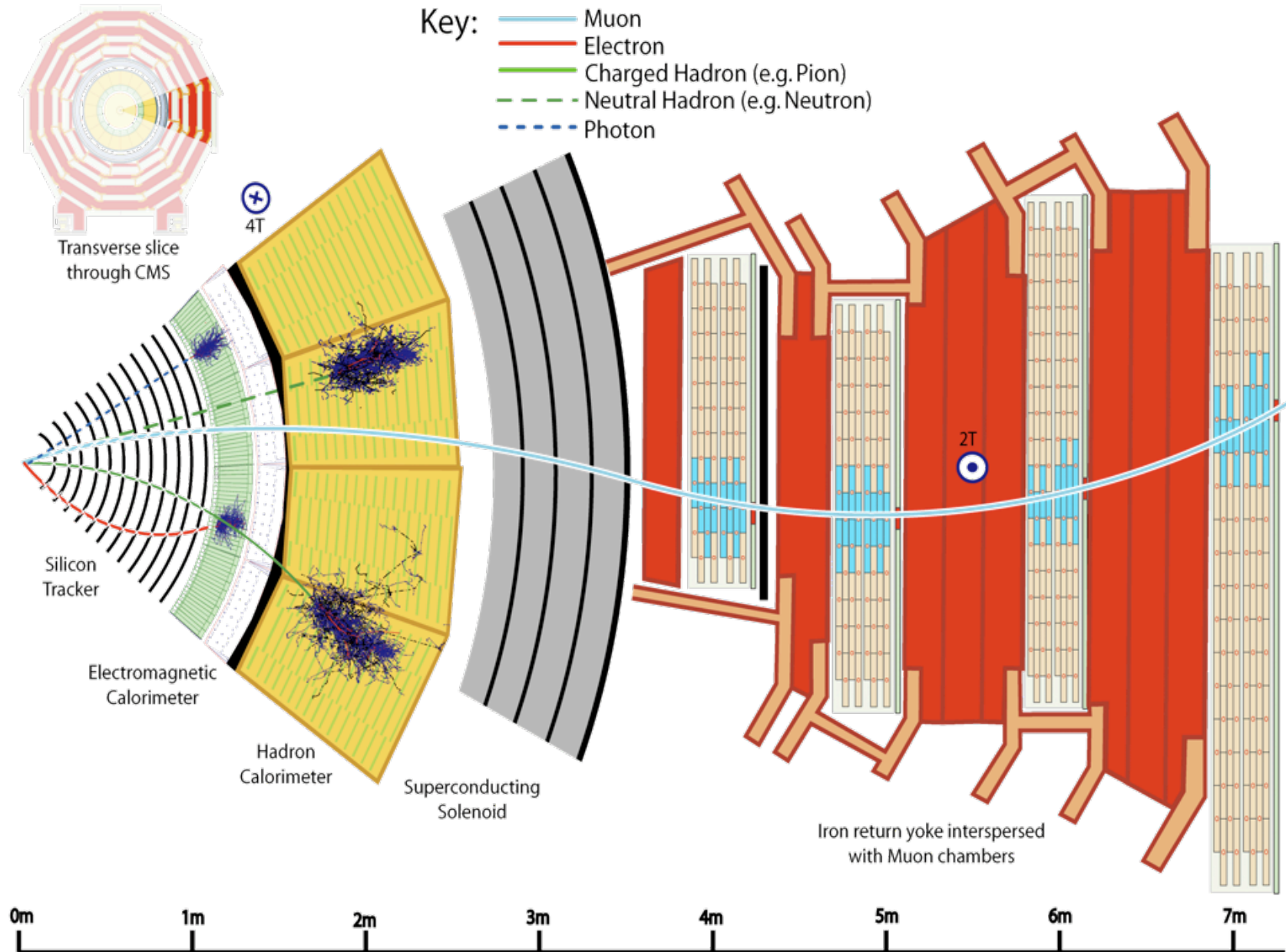
ECAL & HCAL
Inside solenoid



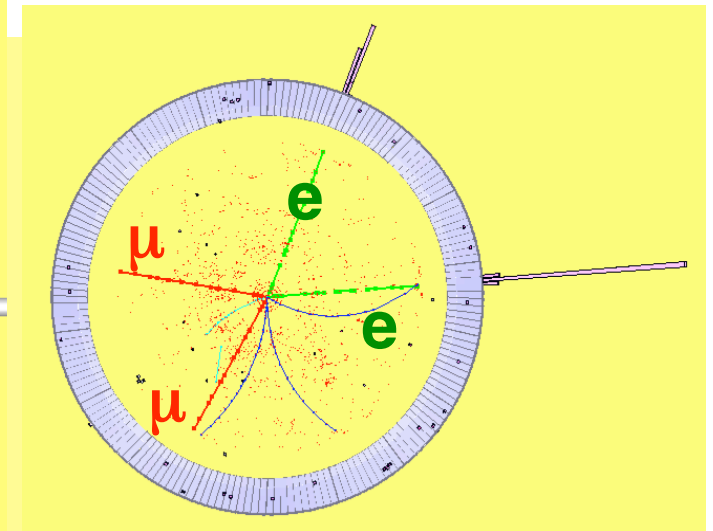
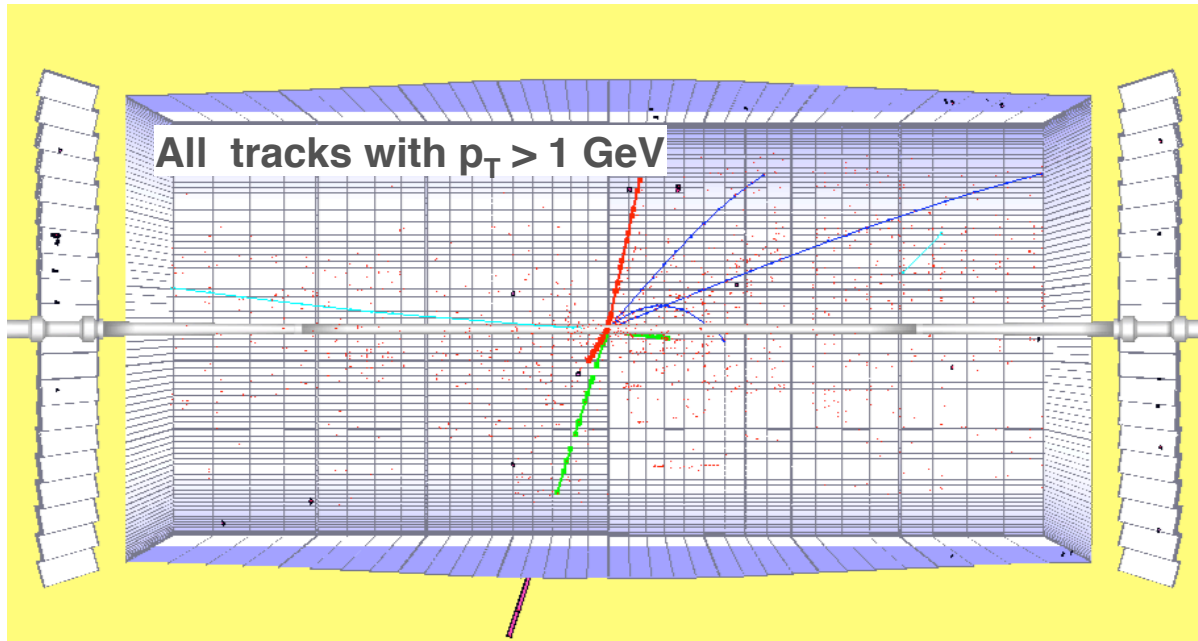
22m Long, 15m Diameter, 14'000 Ton Detector



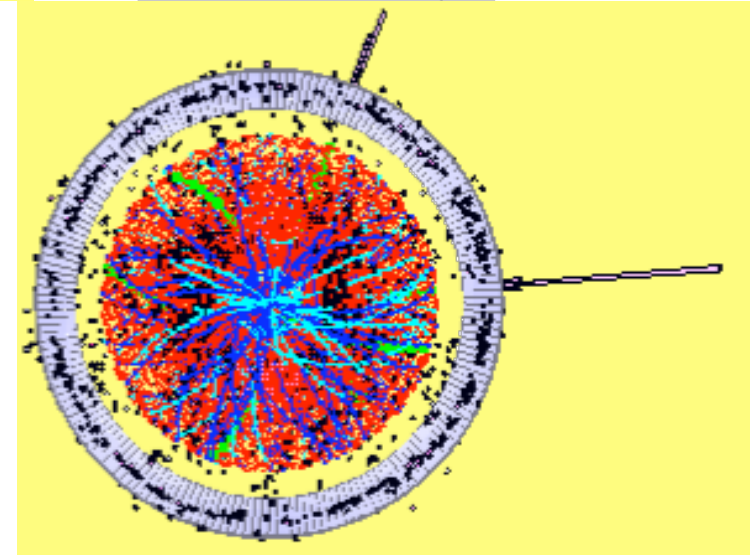
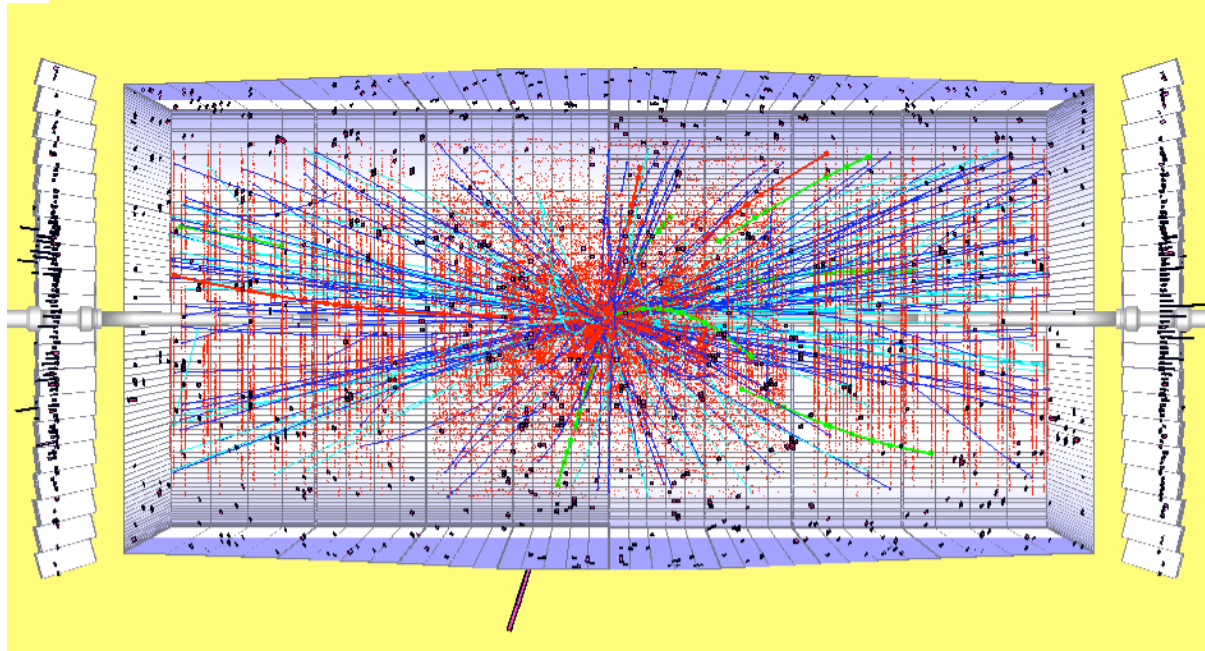
Transverse slice of CMS



What a luminosity of $10^{34}\text{cm}^{-2}\text{s}^{-1}$ means on Tracking ...



$H \rightarrow ZZ^* \rightarrow 2e2\mu$



Tracking Requirements

Precise and efficient tracking at LHC is needed to extract physics signals from competitive but reducible background signals:

- **Mass cut: precise invariant mass determination from precise momentum measurement of the decay products**
- **Isolation cut: high track reconstruction efficiency in dense environment (pile-up, jets...)**
- **Long-lifetime particles tagging: precise and efficient measurement of the track impact parameter and secondary vertices.**

The approach

- ❑ Silicon microstrip detectors allow a very good point resolution (10-30 microns) that coupled to large lever arms m in solenoidal field of 2-4T would allow an adequate momentum resolution, good impact parameter resolution for b-tagging and excellent measurement of the charge up to 1TeV and beyond.
 - ❑ Single bunch crossing resolution is feasible in silicon (collection time $< 10\text{ns}$) with fast read-out electronics.
 - ❑ The real challenge is pattern recognition for track reconstruction: the high density of tracks typical of the inner regions of high luminosity hadronic colliders can be tackled with extreme segmentations both in r - ϕ and r - z : pixel detectors and silicon microstrip modules.
 - ❑ Silicon detectors are well radiation resistant ($\Phi > 10^{14} \text{ n/cm}^2$) and can be produced in large scale at rather low cost.
- ⇒ **Pixel detector** (in the radial region $5\text{cm} < r < 20\text{cm}$: $10^{14}\text{cm}^{-2} < \Phi < 10^{15}\text{cm}^{-2}$)
- ⇒ **Silicon microstrip** (in the radial region $r > 20\text{cm}$: $10^{13}\text{cm}^{-2} < \Phi < 10^{14}\text{cm}^{-2}$)

Two different strategies

This approach has been followed very aggressively by **CMS** who decided to build the first full silicon tracker in HEP

- more challenge in terms of technology and costs
- higher performance particularly in pattern recognition

ATLAS has adopted a more traditional approach based on a hybrid tracker: pixel and silicon microstrip detectors in the innermost part and a large gaseous detectors in the outer part (straw tubes).



ATLAS Inner Detector

- **Transition Radiation Tracker (TRT)**

- 4 mm diameter straw tubes
- 351 k channels
- resolution 130 μm
- polypropylene/polyethylene as transition radiation material:
electron id $0.5 \text{ GeV} < E < 150 \text{ GeV}$

- **SemiConductor Tracker (SCT)**

- 4088 modules
- 80 μm strips (40 mrad stereo)
- 6 M channels
- resolution 17 $\mu\text{m} \times 580 \mu\text{m}$

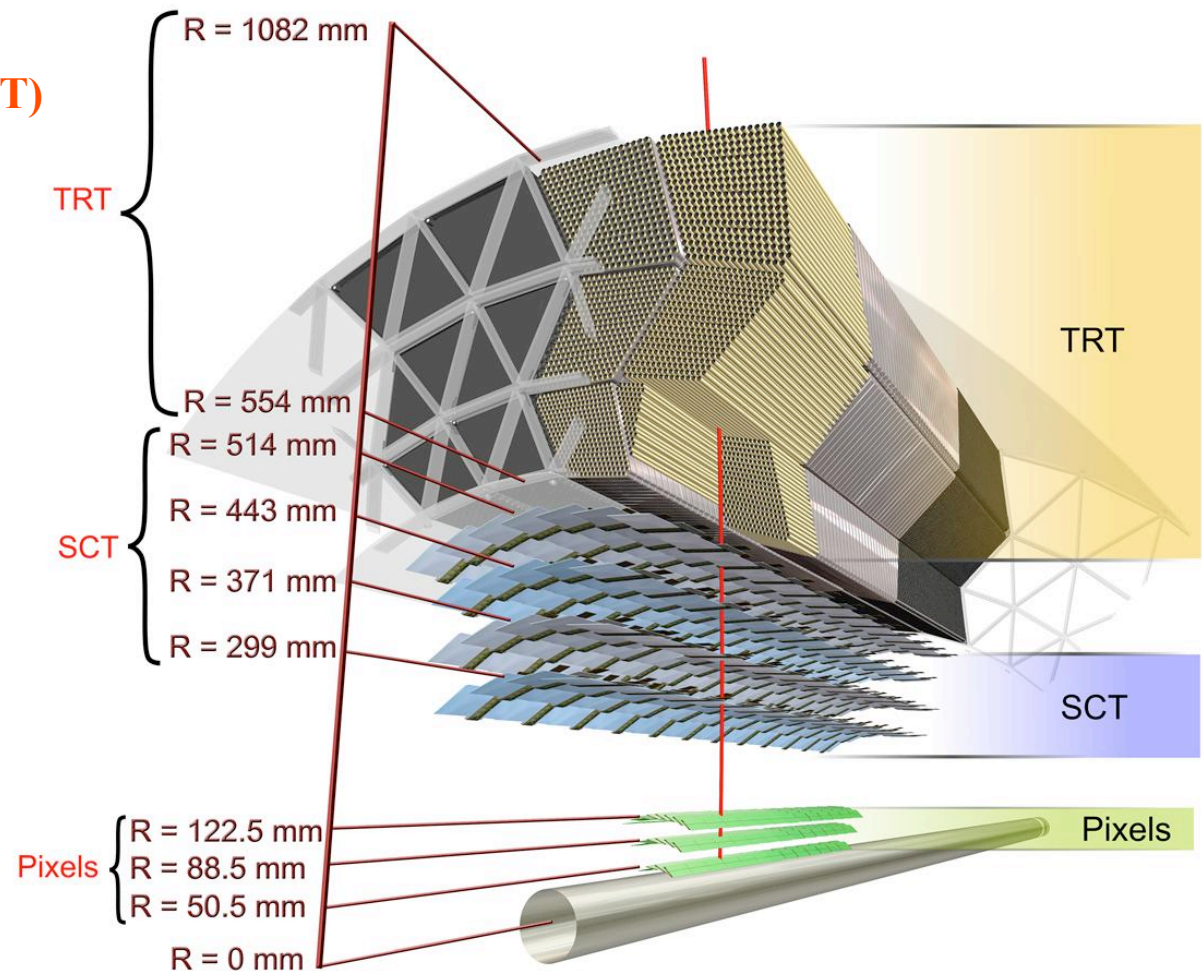
- **Pixel**

- 1744 modules of 46080 pixels
- mostly 50 $\mu\text{m} \times 400 \mu\text{m}$
- 80 M channels
- resolution 10 $\mu\text{m} \times 110 \mu\text{m}$

- **7-points silicon (pixels + strips) tracker ($|\eta| < 2.5$)**

- **straw tubes quasi-continuous tracker (36 points +electron id) (TRT) ($|\eta| < 2$).**

- **2 T solenoidal magnetic field**



- Momentum resolution:
 $\sigma(p_t) / p_t = 0.05\% p_t (\text{GeV} / c) \oplus 1\%$

- IP resolution:
 $\sigma(d_0) = 10 \mu\text{m} \oplus 140 \mu\text{m} / p_t (\text{GeV} / c)$

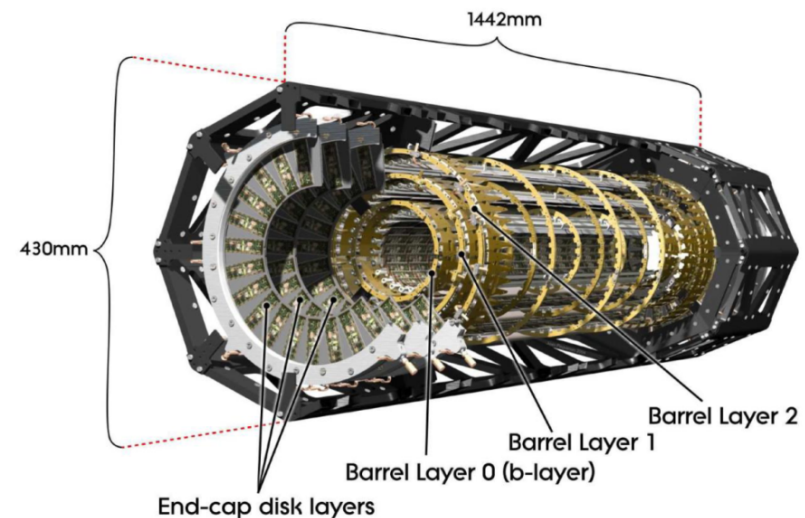
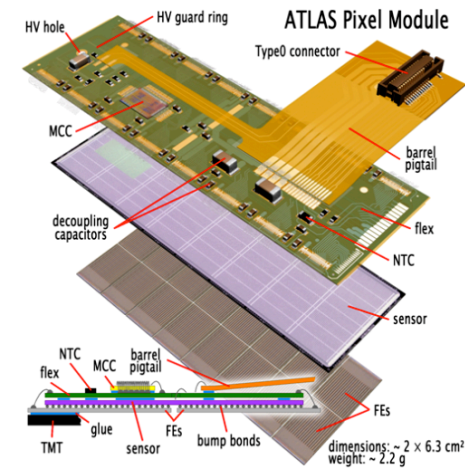


The ATLAS Pixel Detector

- **Requirements:**
 - Position resolution in $r\Phi$ -direction $< 15\mu\text{m}$
 - 3 track points for $|\eta| < 2.5$
 - Time resolution $< 25\text{ ns}$
 - Hit detection efficiency $> 97\%$
- **3 barrel layers**
 - $r = 5.05\text{ cm}, 8.85\text{ cm}, 12.25\text{ cm}$
- **3 pairs of Forward/Backward disks**
 - $z = 49.5\text{ cm}, 58.0\text{ cm}, 65.0\text{ cm}$
- **Pixel size:**
 - $50\text{ mm} \times 400\text{ mm}$ & $50\text{ mm} \times 600\text{ mm}$
- **$\sim 2.0\text{ m}^2$ of sensitive area with 8×10^7 ch**
- **Modules are the basic building elements**
 - 1456 in the barrel + 288 in the endcaps
 - Active area $16.4\text{ mm} \times 60.8\text{ mm}$
 - Radiation tolerance $500\text{ kGy} / 10^{15}\text{ } 1\text{ MeVn}_{\text{eq}}\text{cm}^{-2}$

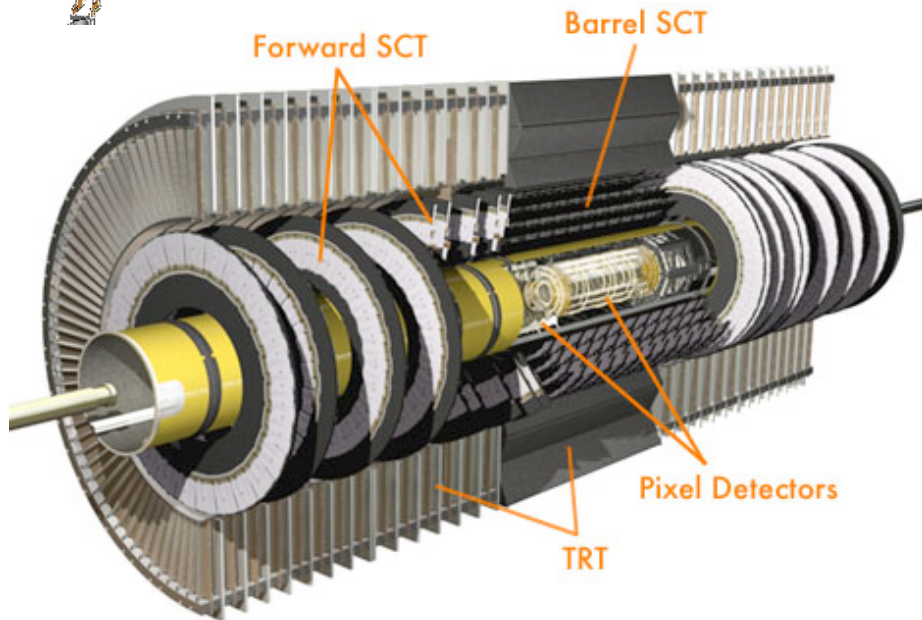
The same module is used in the barrel and in the disks:

- staves (13 modules along the beam axis) for the barrel.
- sectors (6 modules on a two-sided octant) for the disks.





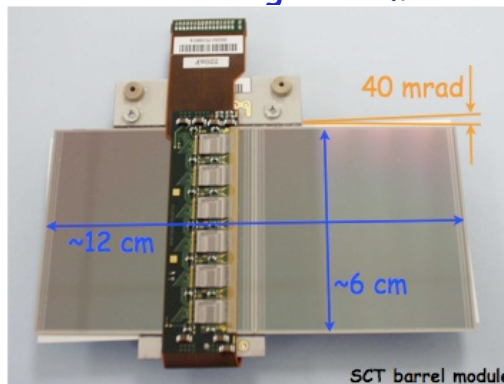
The ATLAS Semiconductor Tracker (SCT)



The SemiConductor Tracker (SCT) is a silicon strip detector. It's organized in 4 layers barrel, built with 2112 modules and two 9 disks end-caps, made of 1976 modules. The total number of strips is $6.3 \cdot 10^6$.

The barrel module consists of four single sided p-on-n strip detectors:

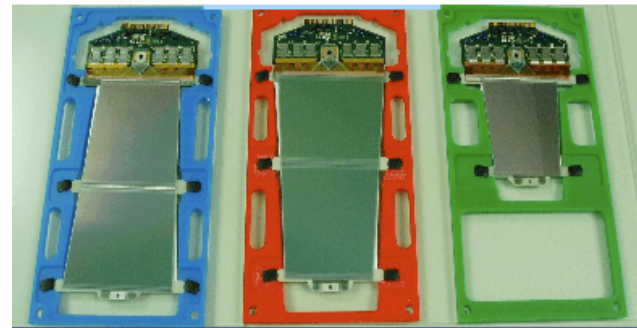
- Pitch 80 mm
- Strip length 120 mm
- Stereo angle 40 mrad



SCT barrel module

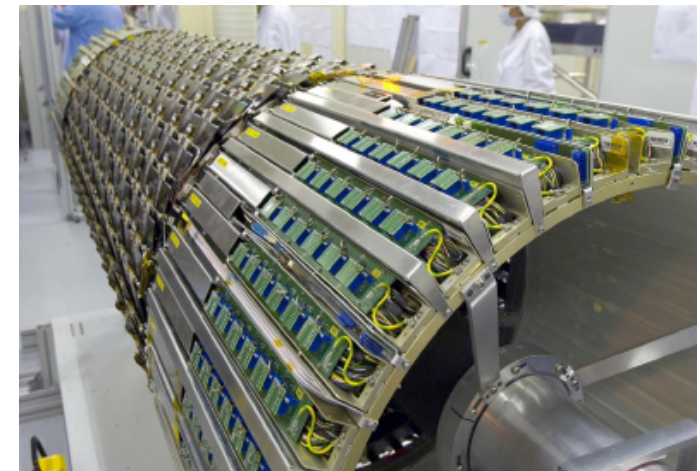
The end-caps are built with three different modules:

- Pitch 57-95 mm
- Strip length 55-120 mm



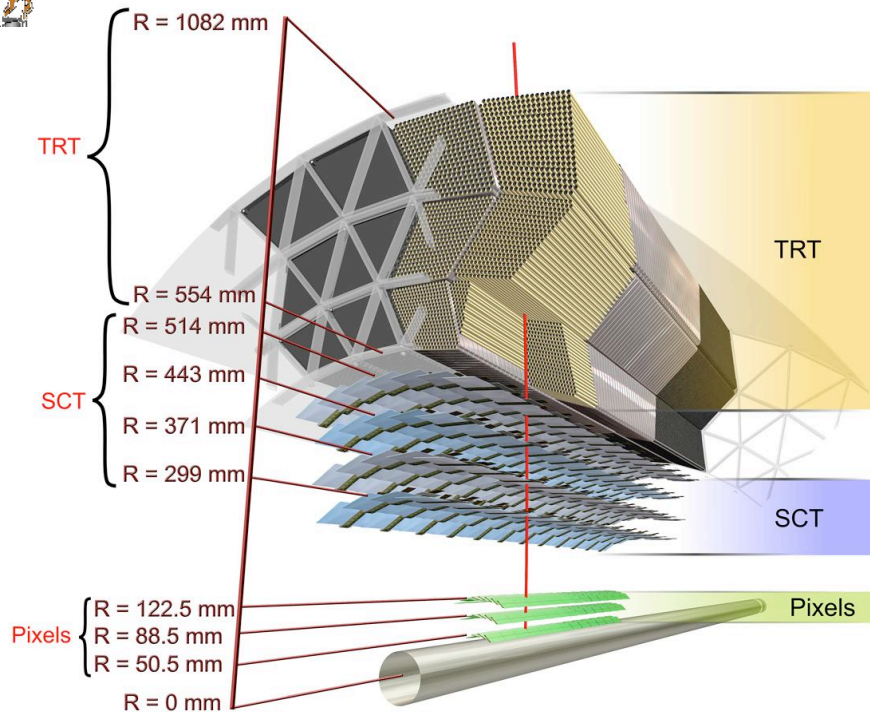
Physics requirements

- Resolution (x* y): $17 \mu\text{m} * 580 \mu\text{m}$
- Alignment tolerances (x*y): $12 \mu\text{m} * 50 \mu\text{m}$
- Noise occupancy: $< 5 * 10^{-4}$
- Efficiency: $> 99\%$
- Radiation : $\sim 2 * 10^{14} \text{ Neq cm}^{-2}$ over 10 years





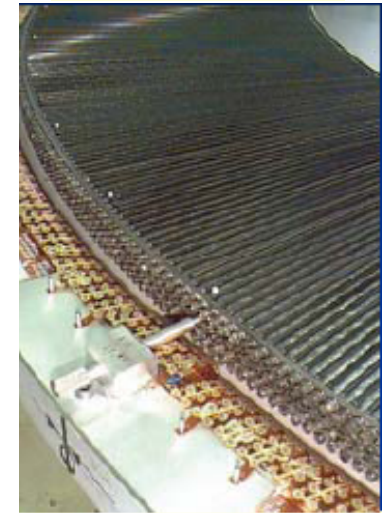
The ATLAS Transition Radiation Tracker (TRT)



TRT: 4 mm straw tubes, arranged in 2 · 160 disks (end-cap) and 73 layers (barrel), 40 K channels, acceptance $|\eta| < 2$

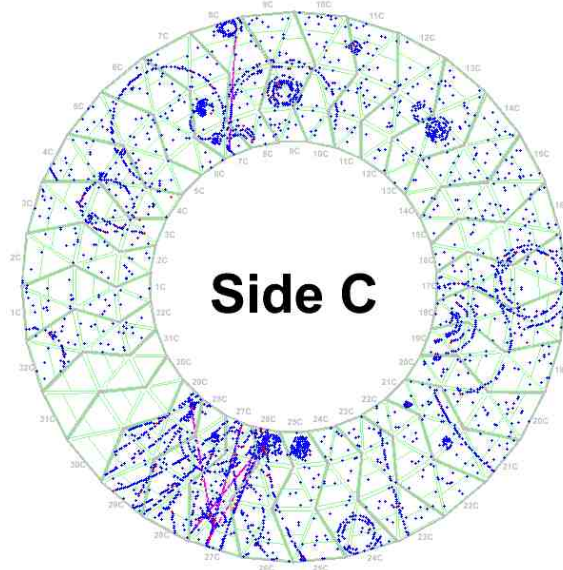
TR (polypropylene foils/fiber): pion-electron separation (TR γ 's convert into e 's in Xe)

TRT end-cap during assembly



Hits (Barrel):
 $3(\text{Pixel}) + 4(\text{SCT}) + \langle 36 \rangle (\text{TRT})$
 $\sigma/p_T \sim 5 \times 10^{-4} p_T \oplus 0.01$

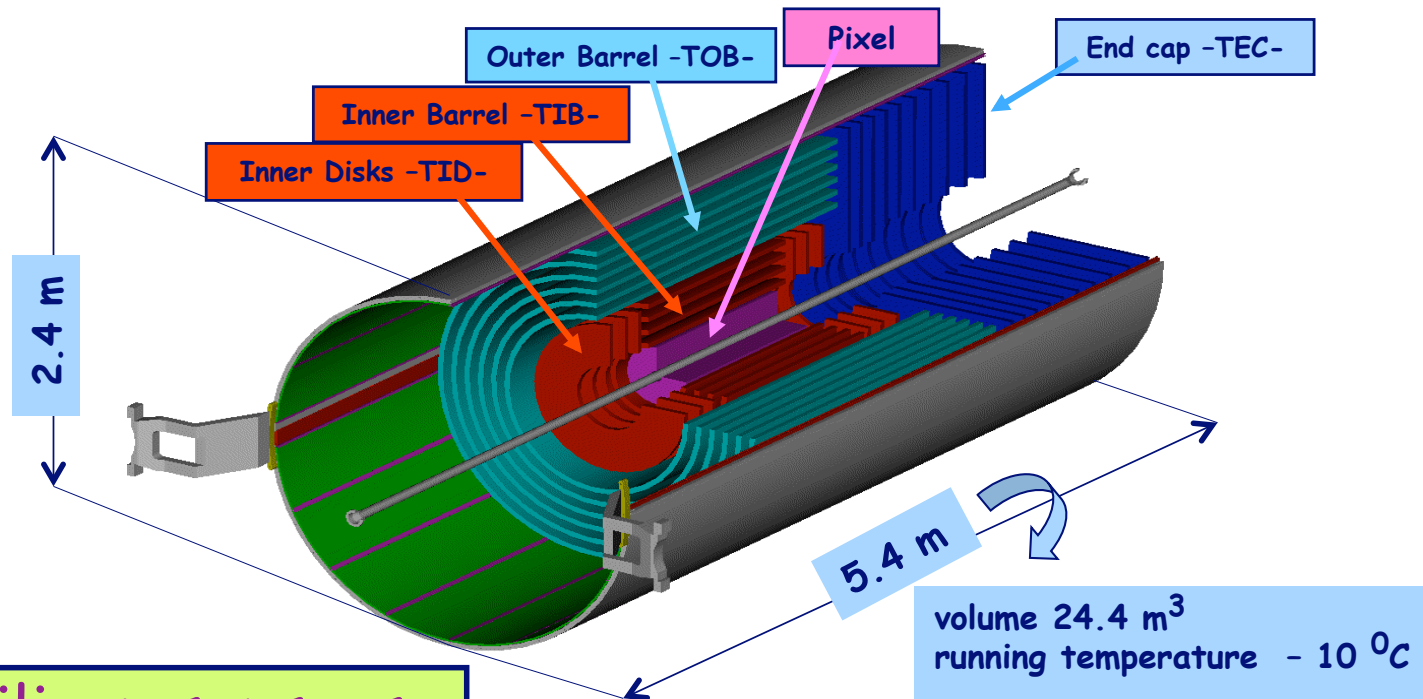
Cosmic ray event in TRT



$\frac{1}{4}$ of TRT barrel during integration



CMS has chosen an all-silicon configuration



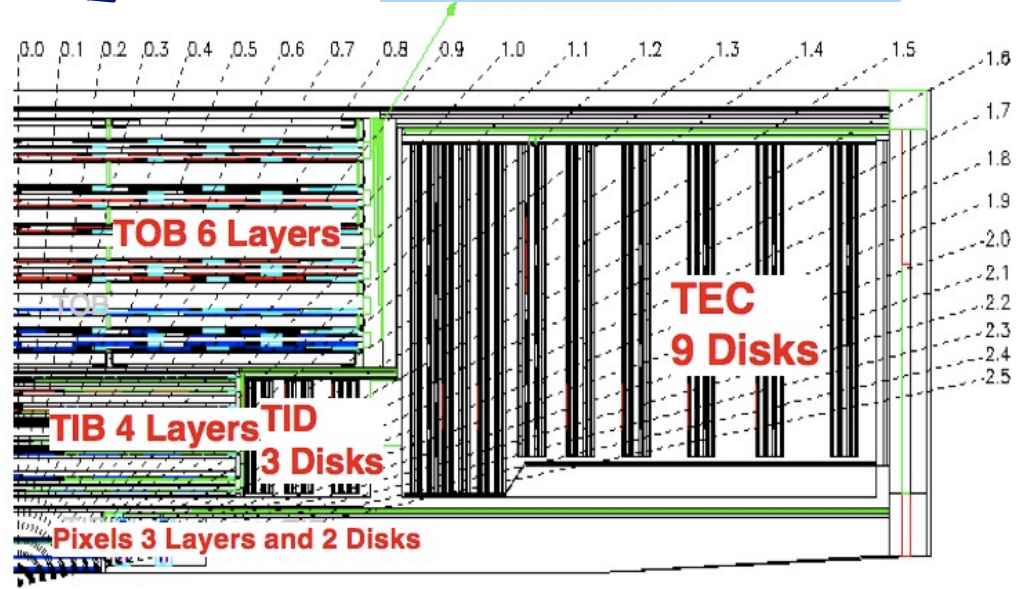
207m² of silicon sensors
10.6 million silicon strips
65.9 million pixels ~ 1.1m²

Strips:

Pitch: 80 μm to 180 μm
Hit Resolution: 20 μm to 50 μm

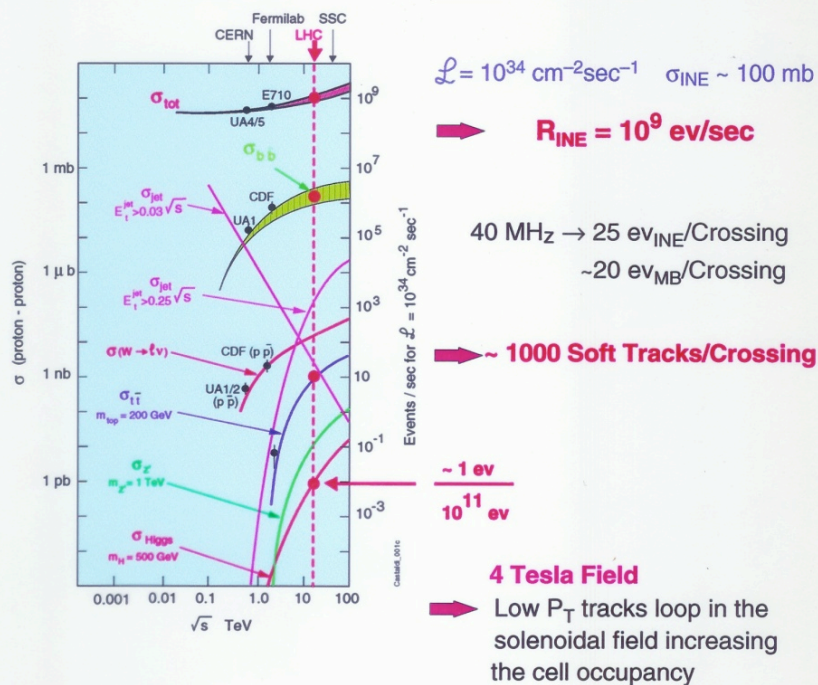
Pixels:

100 μm x 150 μm
 $r\phi$ and z resolution: ~15 μm



CMS Tracking: very old slides presented to LHCC (towards a full silicon tracker)

LHC ENVIRONMENT



- ⇒ • Fast Detector and Electronics: **Bunch Crossing Ident.**
- Radiation Hard Electronics and Detectors: **Mrad. Doses**
- Very Difficult Pattern Recognition:
 Efficient track finding requires small cell occupancies
 $\Rightarrow \geq 10^7$ channels are needed
- ⇒ High precision fast detectors with high granularity
 like
Pixels + Si-Strip + MSGC

Rino_001



Detector characteristics

- Good quality detectors available in 6'' technology from at least two producers each with over-capacity to produce all the silicon wafers in less than 2.5 years
- Longer strips and thicker detectors can be used in the outer Tracker
 - The strip capacitance (per unit length), (and therefore electronics noise), is almost independent of the sensor thickness: (It is also ~ independent of the pitch and depends only on the ratio w/p of strip width to pitch)

$$C_{400\mu\text{m}}(6'') = C_{300\mu\text{m}}(4''/6'') = C(w/p)$$
 - Charge collection:

$$Q(400 \mu\text{m}) = 1.3 \times Q(300 \mu\text{m})$$
 - Rad-hardness:

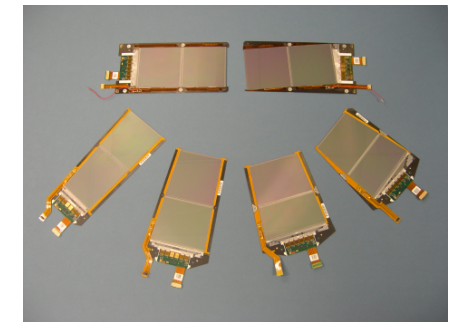
$$\langle 100 \rangle 6'' = \langle 100 \rangle 4''$$

$$V_{\text{bdwn}} > 500 \text{ V for } w/p > 0.2$$

LHCC 8/3/2000, R. Castaldi

1

The CMS Silicon Strip Tracker

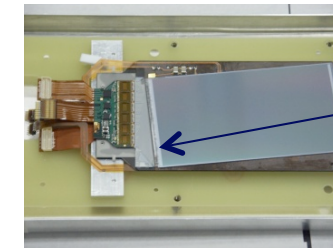


Outer Barrel (TOB): 6 layers

- Thick sensors (500 μm)
- Long strips

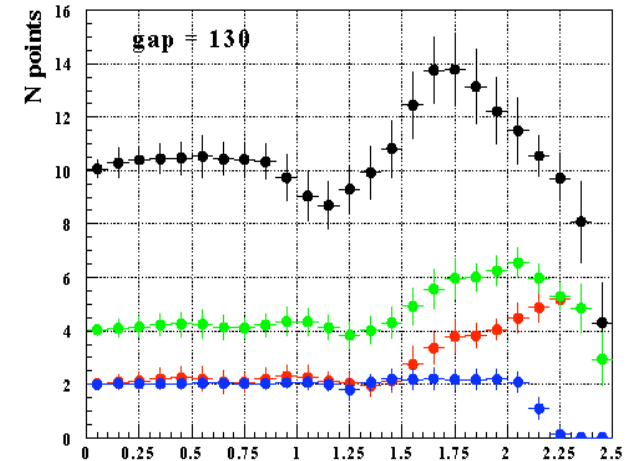
Endcap (TEC): 9 Disk pairs

- $r < 60$ cm thin sensors
- $r > 60$ cm thick sensors

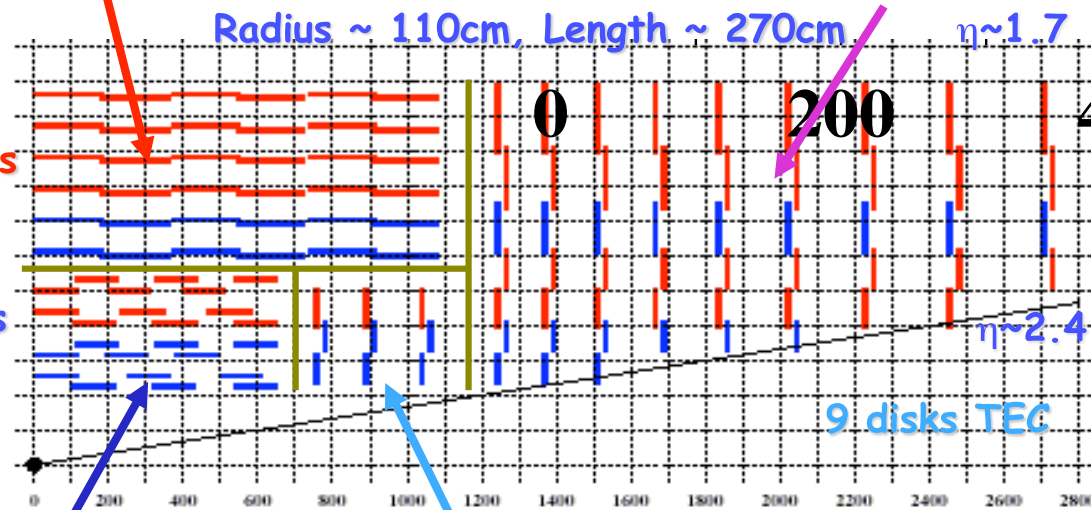


small angle tilt (100mrad) to provide second coordinate measurement

9'648'128 strips
 75'376 APV chips
 6'136 Thin sensors
 18'192 Thick sensors
 ~17'000 modules
 ~25'000'000 Bonds



Black: total number of hits
 Green: double-sided hits
 Red: ds hits - thin detectors
 Blue: ds hits - thick detectors



Inner Barrel (TIB): 4 layers

- Thin sensors (320 μm)
- Short strips

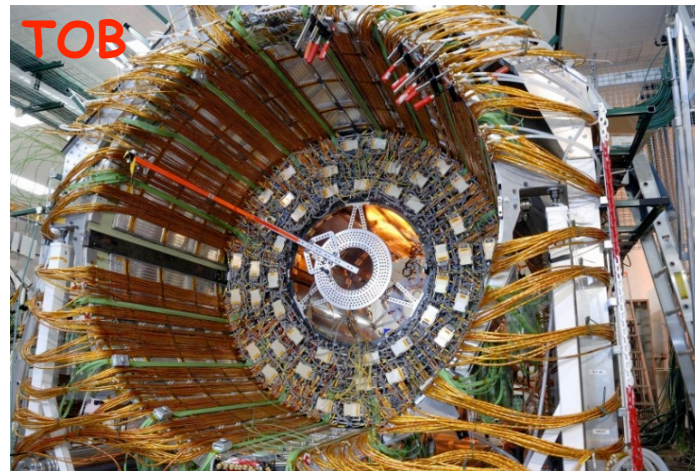
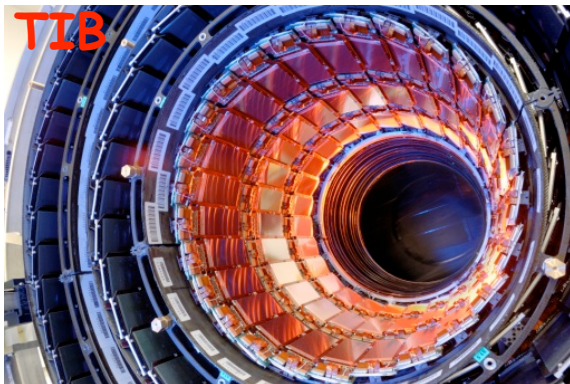
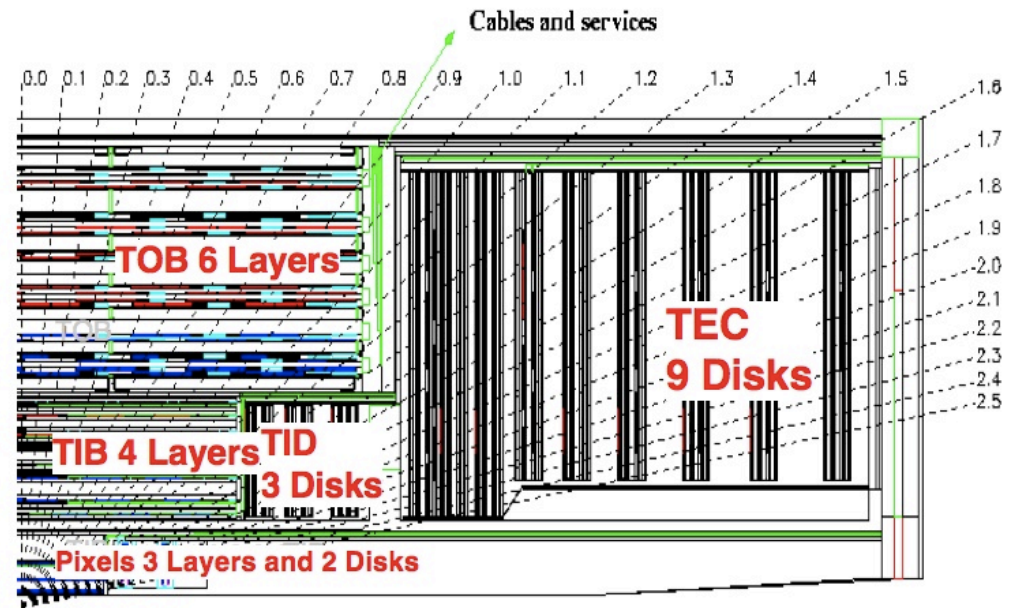
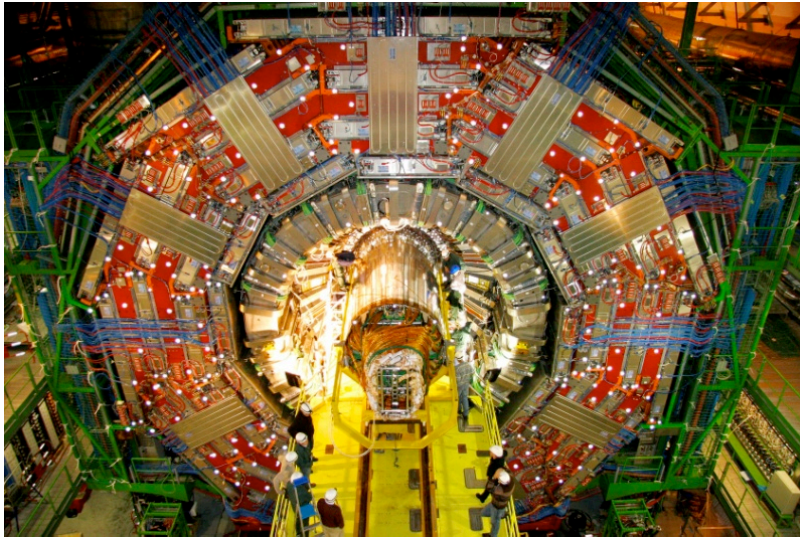
Inner Disks (TID): 3 Disk pairs

- Thin sensors

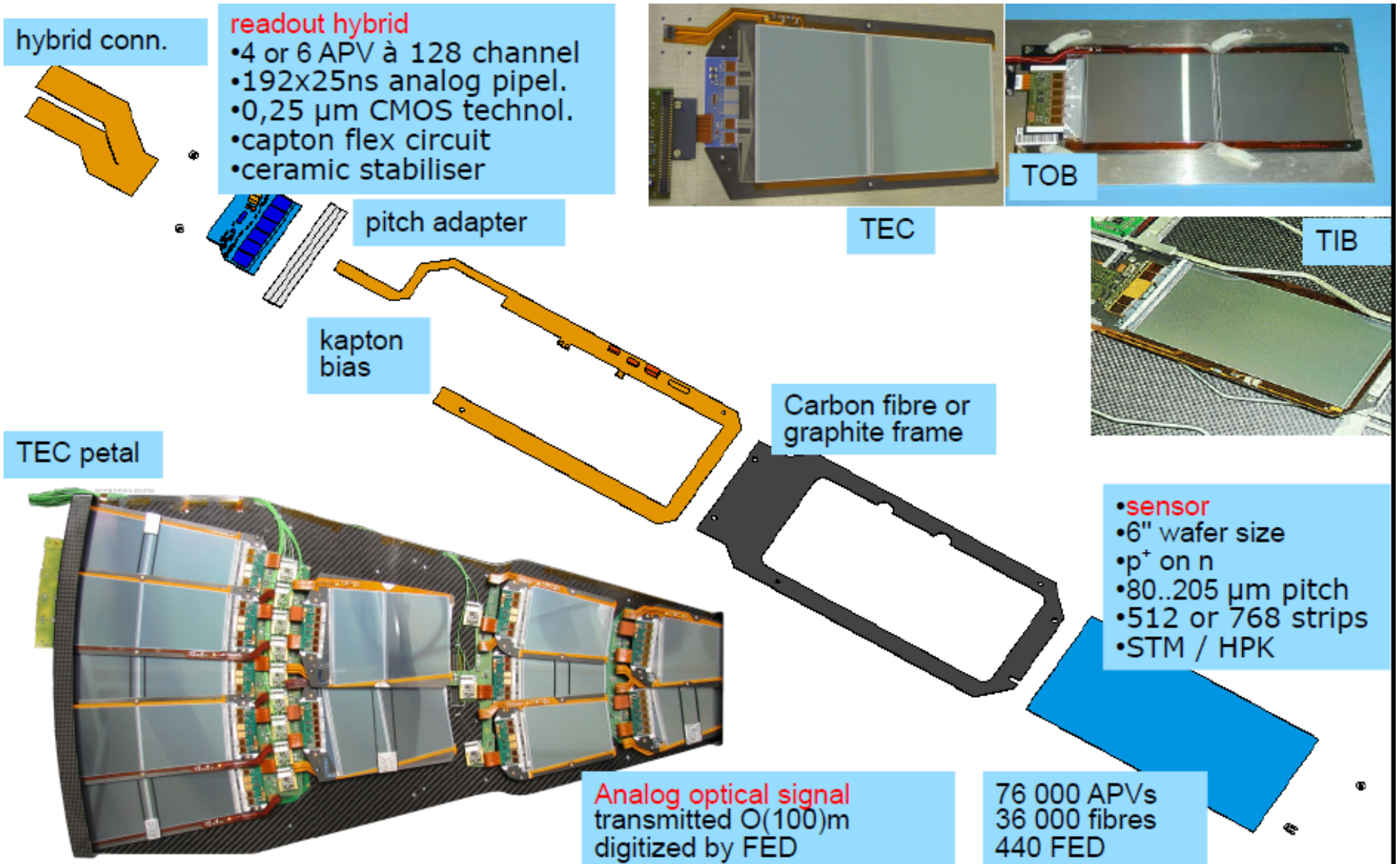
Strip length ranges from 10 cm in the inner layers to 20 cm in the outer layers.

Pitch ranges from 80 μm in the inner layers to near 200 μm in the outer layers

The CMS Silicon Strip Tracker



CMS modules of Silicon Strip Tracker

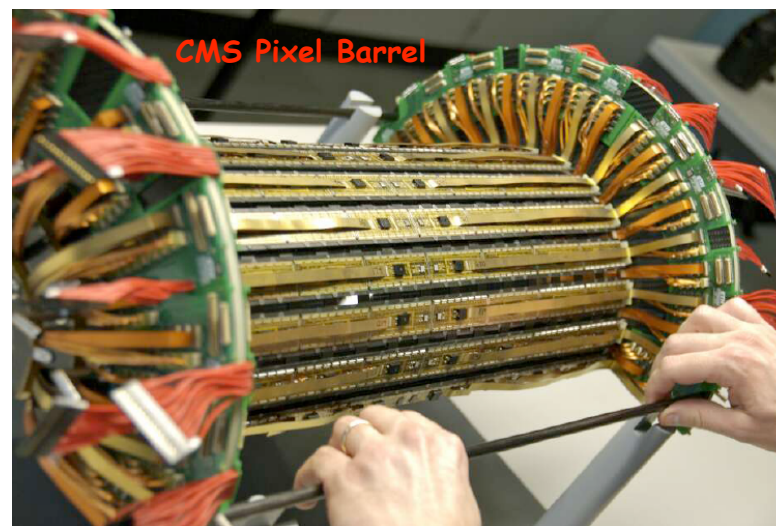
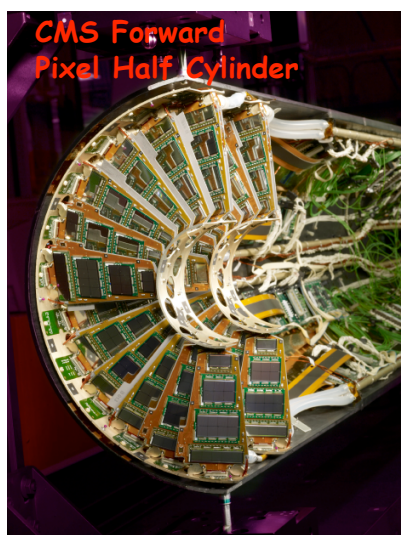
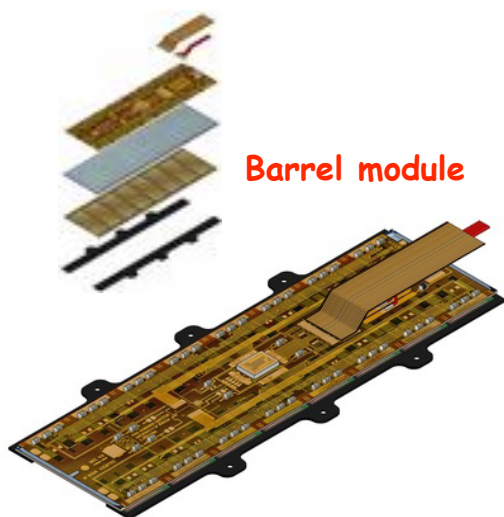
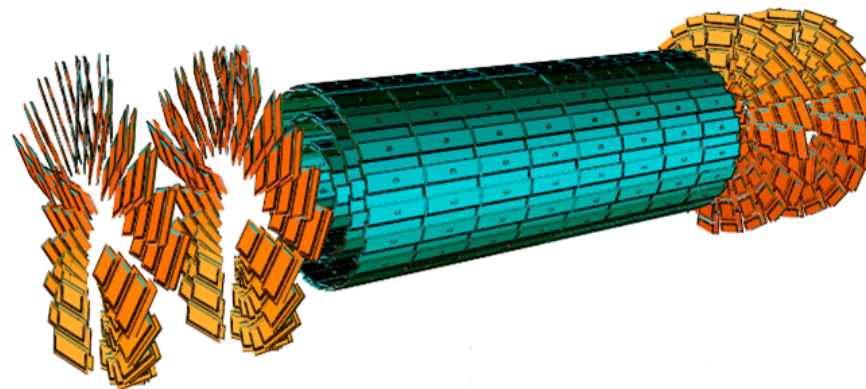


The CMS Pixel Detector

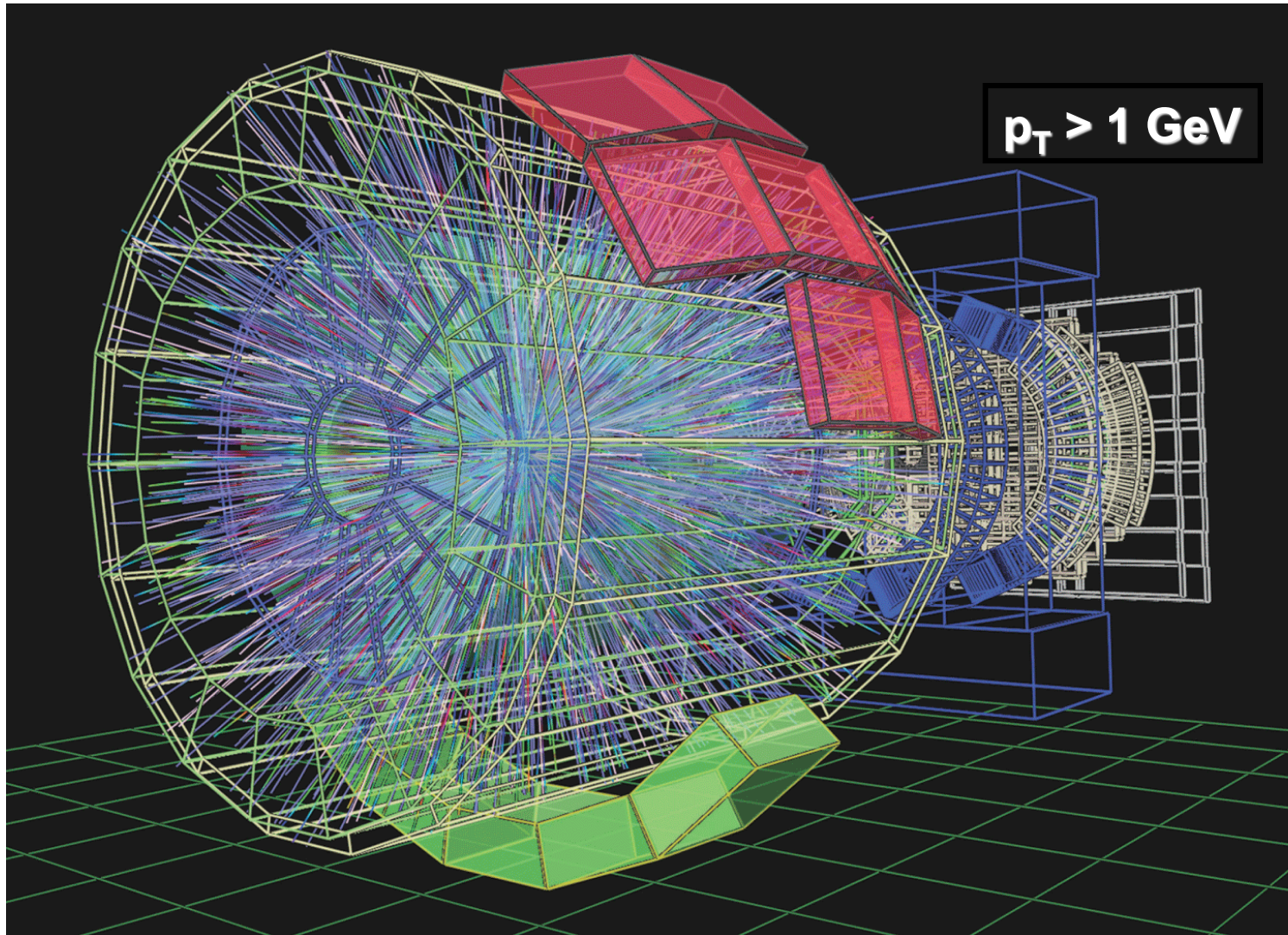
- 3 Barrel layers at 4.3, 7.2, 11.0 cm
- 2 Forward Disks
- 3-hit coverage for tracks $|\eta| < 2.5$

Total Area: $0.78 + 0.28 \text{ m}^2$
66 Million Pixels

Sensors: n on n Silicon 265 - 270 μm
150 x 100 μm pixels $\sigma(z) \sim \sigma(r\phi) \sim 15\mu\text{m}$
Bump-bonded to PSI 46 Read Out Chips (analog readout)



...and in ALICE tracking will not be easier...



An ALICE PbPb Event

$N_{ch}(-0.5 < \eta < 0.5) = 8000 !!$

ALICE

L3 Solenoid Magnet
 $B=0.5T$



HMPID
TRD
TPC
PMD
ITS
TOF
PHOS
FMD

MUON SPECTR.

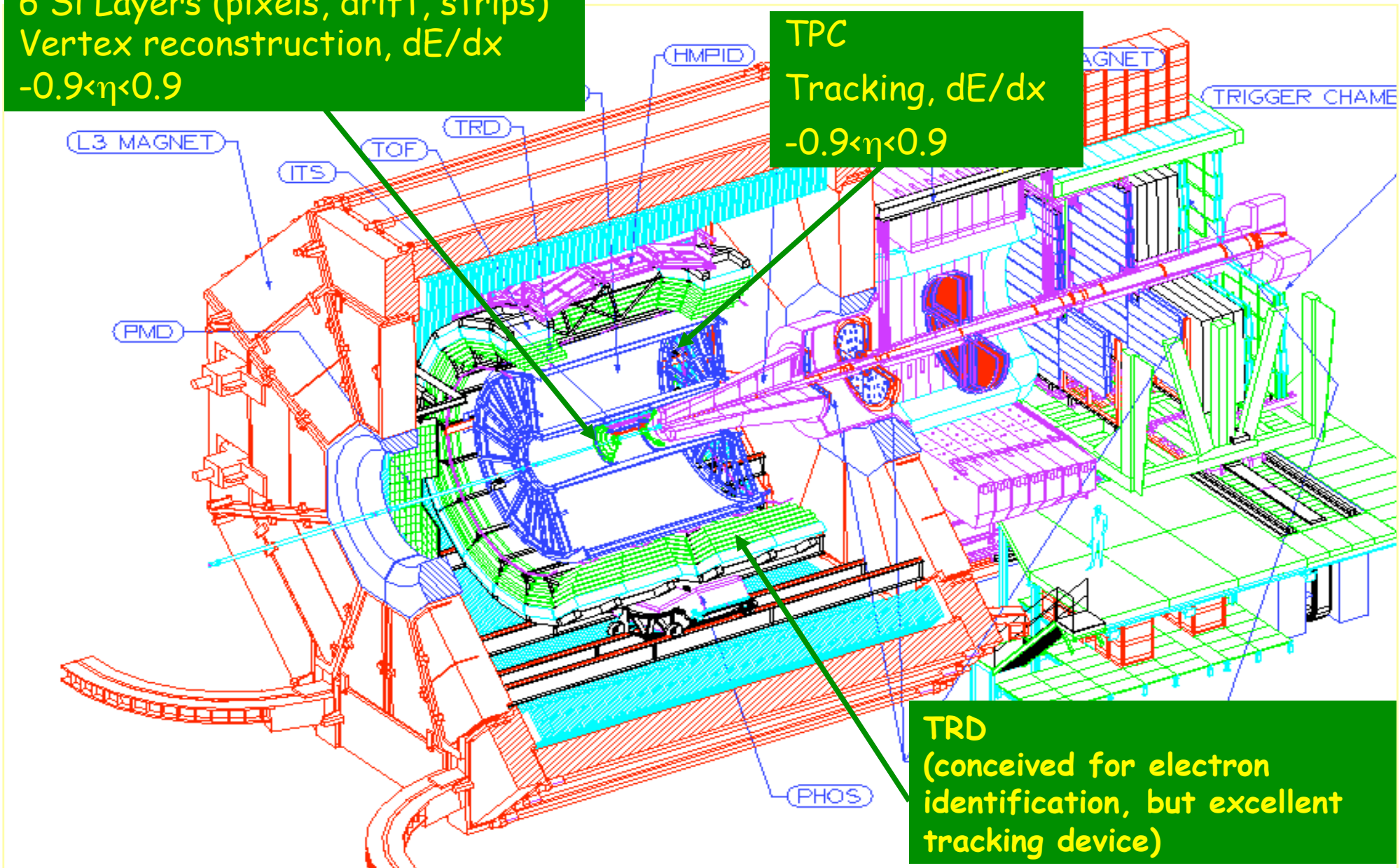
Size: 16 x 26 meters
Weight: 10,000 tons

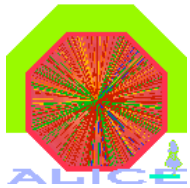
ALICE Tracking Layout

Inner Tracking System (ITS):
6 Si Layers (pixels, drift, strips)
Vertex reconstruction, dE/dx
 $-0.9 < \eta < 0.9$

TPC
Tracking, dE/dx
 $-0.9 < \eta < 0.9$

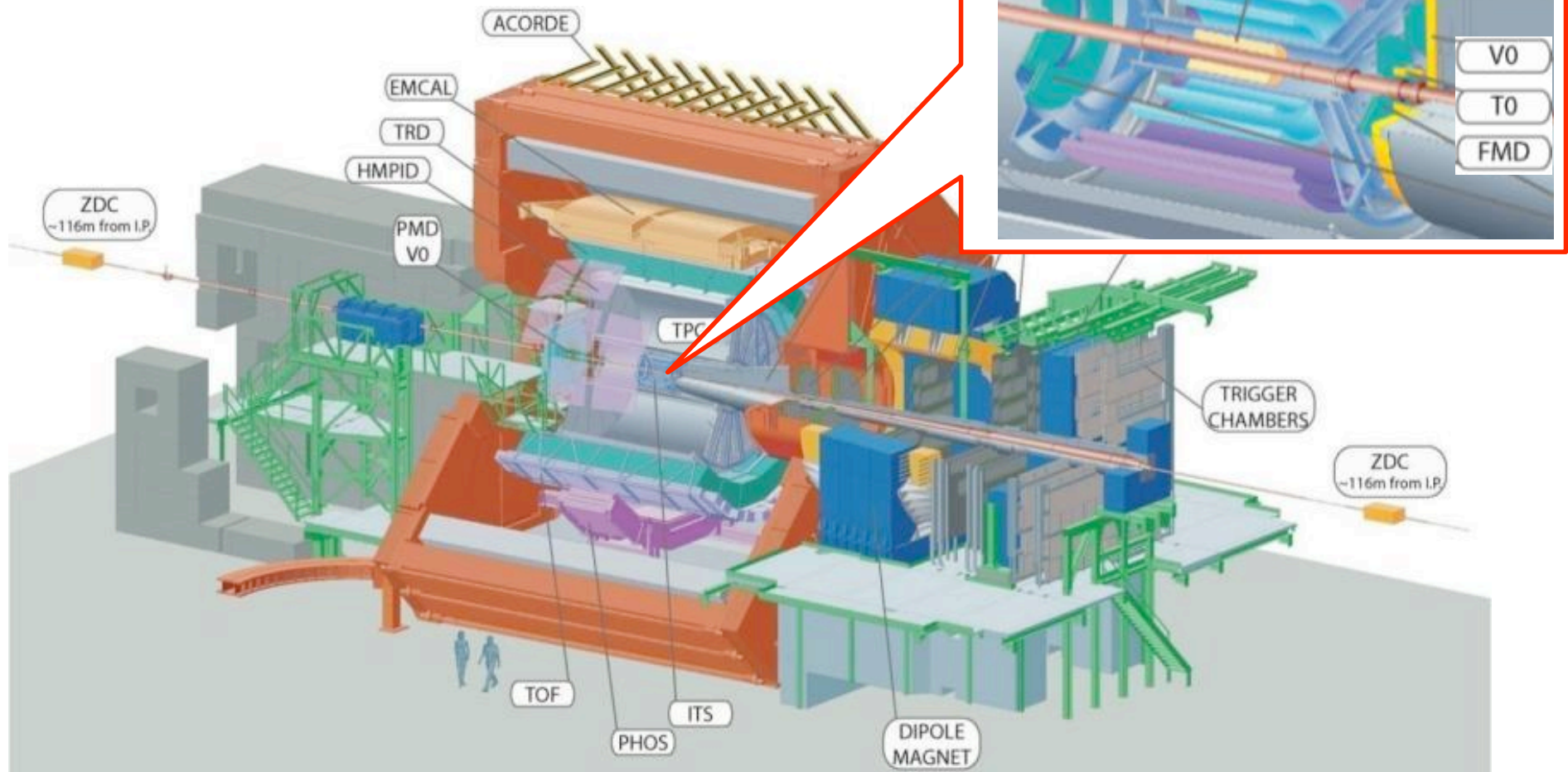
TRD
(conceived for electron
identification, but excellent
tracking device)

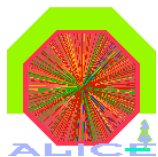




ALICE Inner Tracking System

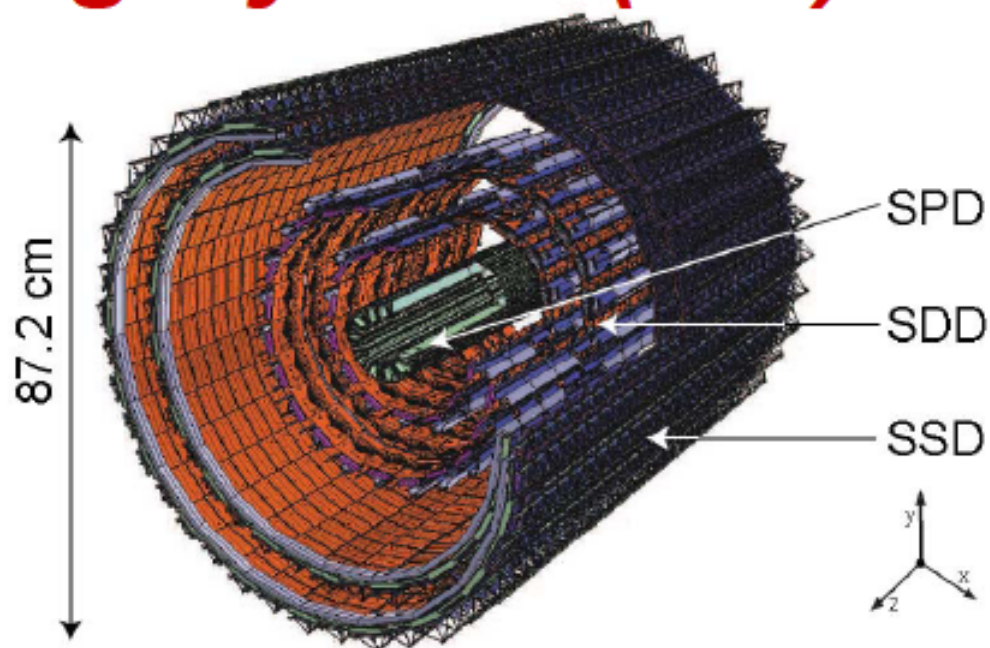
- 6 barrel layer
- 3 different silicon detector technologies,
2 layers each, as seen by produced particles:
 - Pixels (SPD), Drift (SDD), double-side Strips (SSD)





Inner Tracking System (ITS)

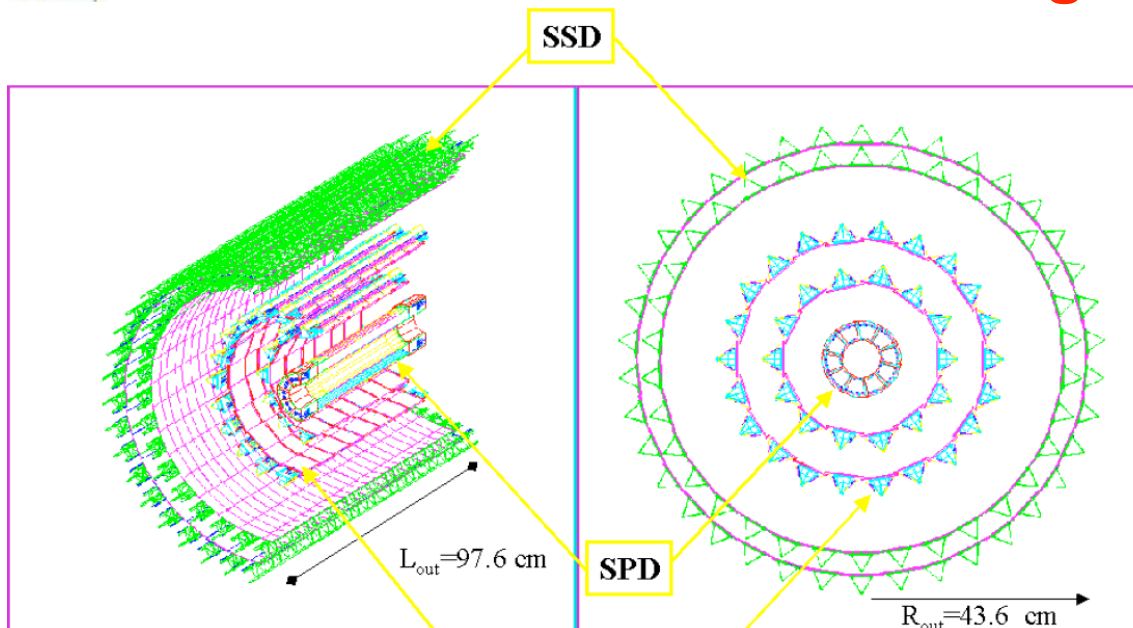
- Six layers of silicon detectors
 - ⇒ Coverage: $|\eta| < 0.9$
- Three technologies
 - ⇒ Pixels (SPD)
 - ⇒ Drift (SDD)
 - ⇒ Double-sided Strips (SSD)
- Design goals
 - ⇒ Optimal resolution for primary vertex and track impact parameter
 - ✓ Minimize distance of innermost layer from beam axis ($\langle r \rangle \approx 3.9$ cm) and material budget
 - ⇒ Maximum occupancy (central PbPb) < few %
 - ⇒ 2D devices in all the layers
 - ⇒ dE/dx information in the 4 outermost layers for particle ID in $1/\beta^2$ region



Layer	Det. Type	Radius (cm)	Length (cm)	Resolution (μm)	
				$r\phi$	Z
1	SPD	3.9	28.2	12	100
2	SPD	7.6	28.2	12	100
3	SDD	15.0	44.4	35	25
4	SDD	23.9	59.4	35	25
5	SSD	38.0	86.2	20	830
6	SSD	43.0	97.8	20	830



Inner Tracking System



Silicon Pixel Detector (SPD):

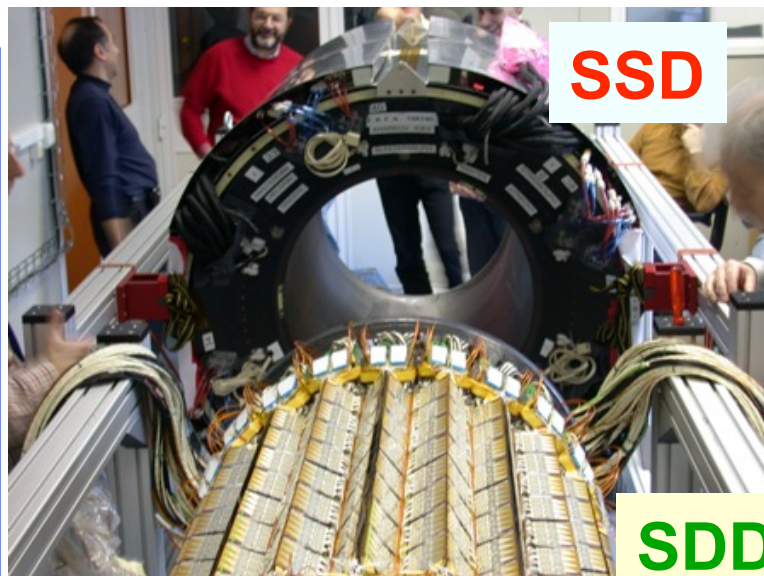
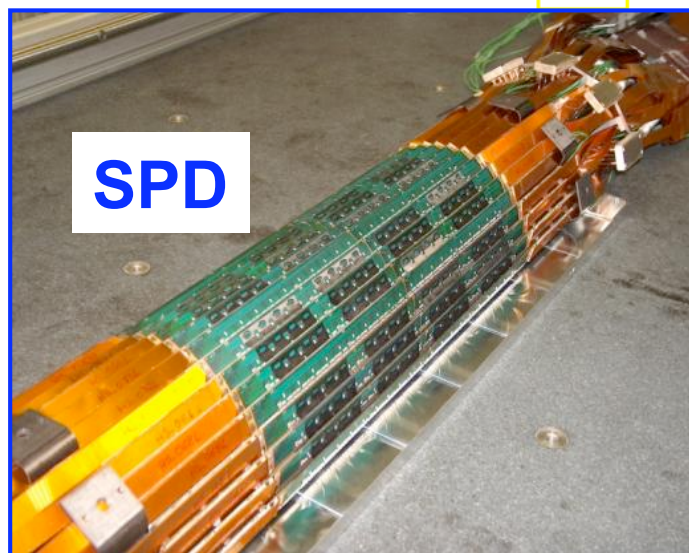
- ~10M channels
- 240 sensitive vol. (60 ladders)

Silicon Drift Detector (SDD):

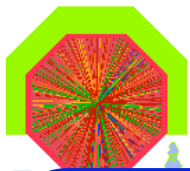
- ~133k channels
- 260 sensitive vol. (36 ladders)

Silicon Strip Detector (SSD):

- ~2.6M channels
- 1698 sensitive vol. (72 ladders)

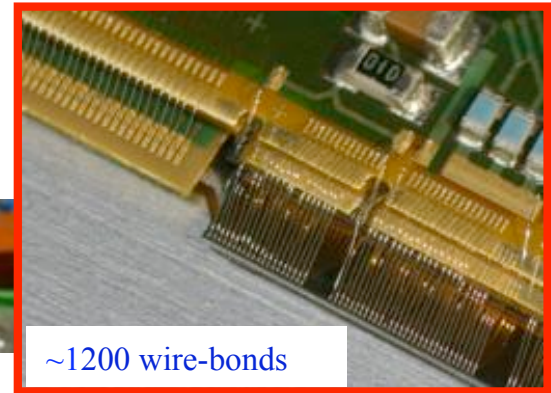
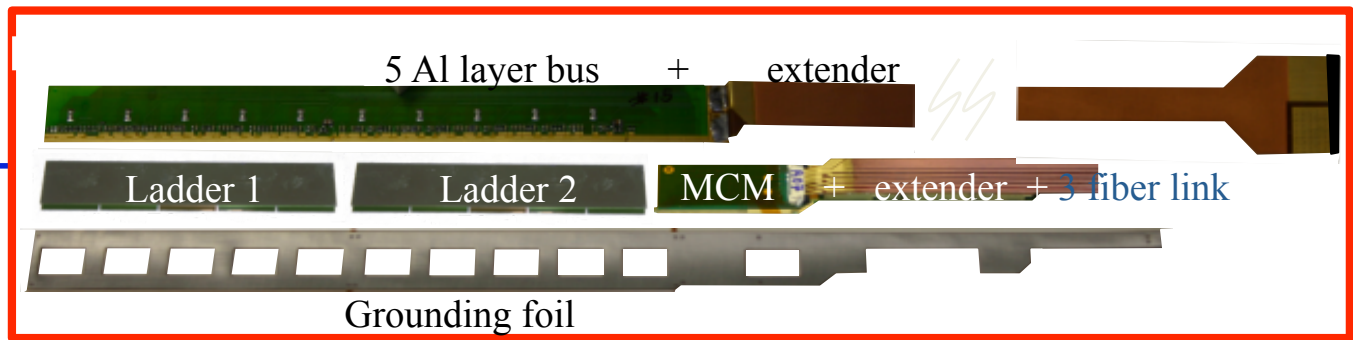


ITS total:
2198 alignable sensitive volumes
→ 13188 d.o.f.

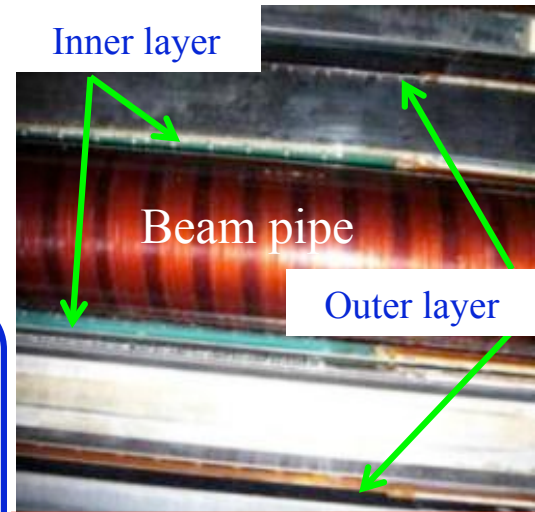
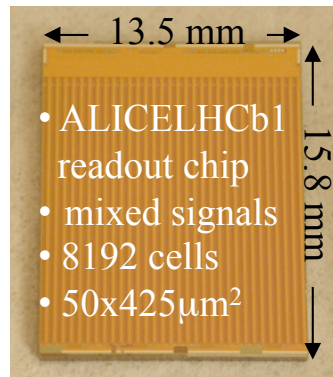


Pixel

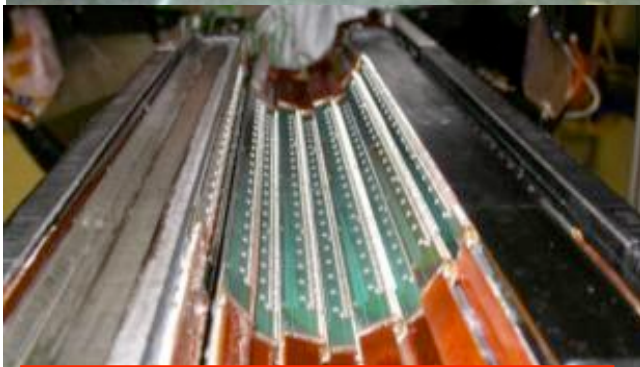
- 2 layer barrel
- Total surface: $\sim 0.24\text{m}^2$
- Power consumption $\sim 1.4\text{kW}$
- Evaporative cooling C_4F_{10}
- Operating at room temperature
- Material budget per layer $\sim 1\% X_0$



Outer surface: 80 half-staves



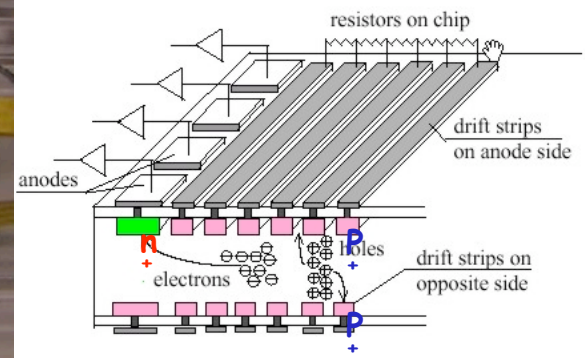
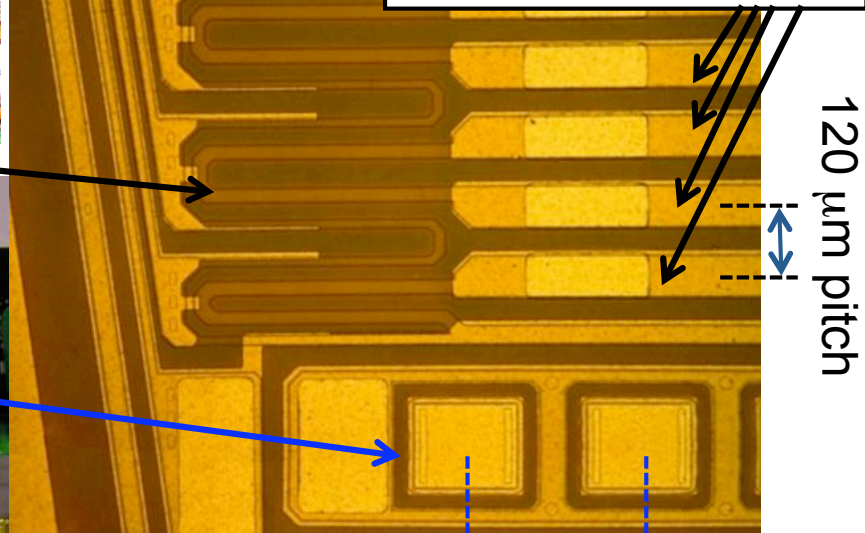
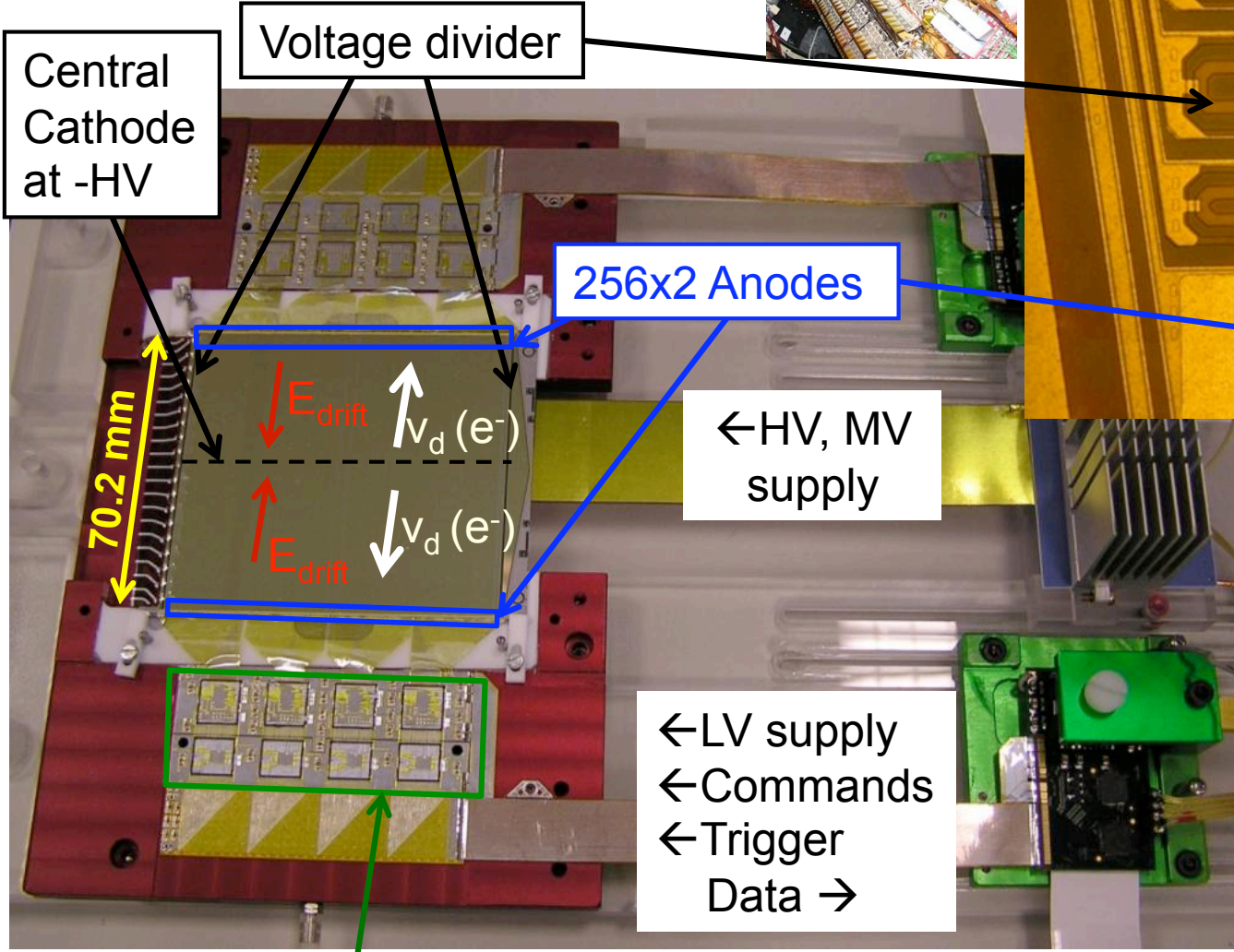
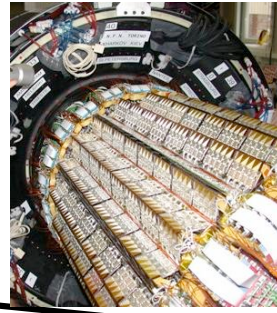
Minimum distance inner layer-beam pipe $\approx 5\text{ mm}$



Inner surface: 40 half-staves

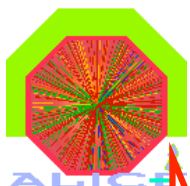
- Unique L0 trigger capability**
- Prompt FastOR signal from each chip
 - Extract and *synchronize* 1200 FastOR signals from the 120 half-staves
 - User defined programmable algorithms

Silicon Drift Detector

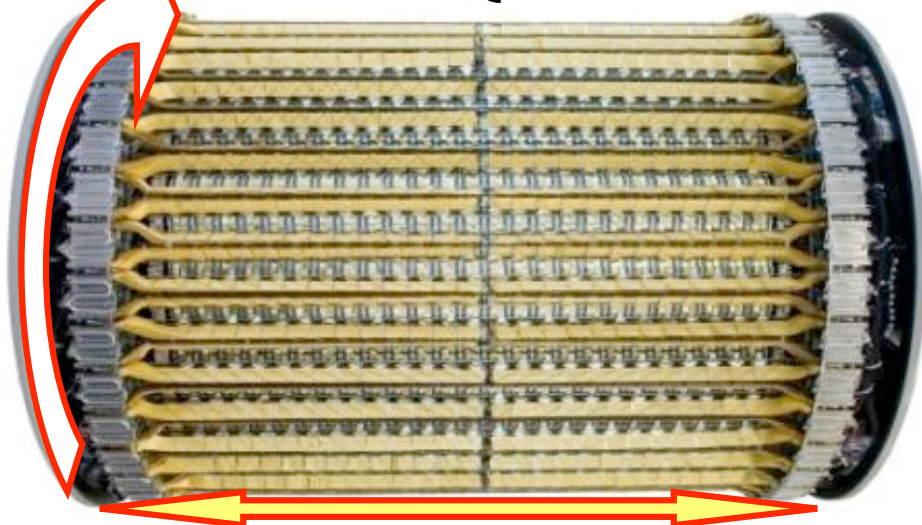


Hybrid with front-end electronics (4 pairs of ASICs)

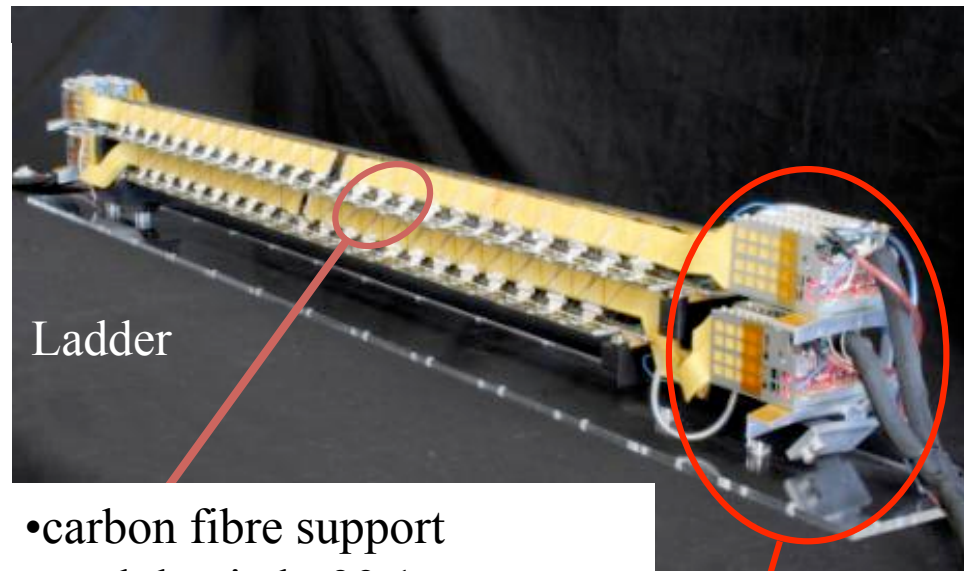
double-side Strip



$r\phi$ - overlap: $\begin{cases} \text{L5: 34 ladders} \\ \text{L6: 38 ladders} \end{cases}$



z - overlap: $\begin{cases} \text{L5: 22 modules} \\ \text{L6: 25 modules} \end{cases}$



Ladder

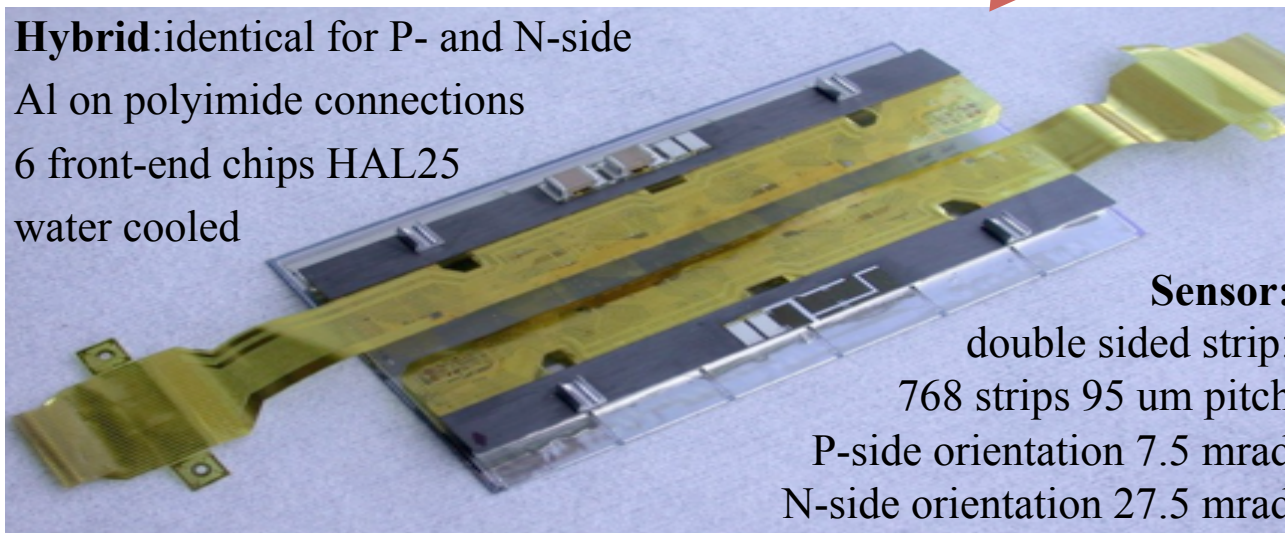
- carbon fibre support
- module pitch: 39.1 mm
- Al on polyimide ladder cables

Hybrid: identical for P- and N-side

Al on polyimide connections

6 front-end chips HAL25

water cooled



Sensor:

double sided strip:

768 strips 95 μm pitch

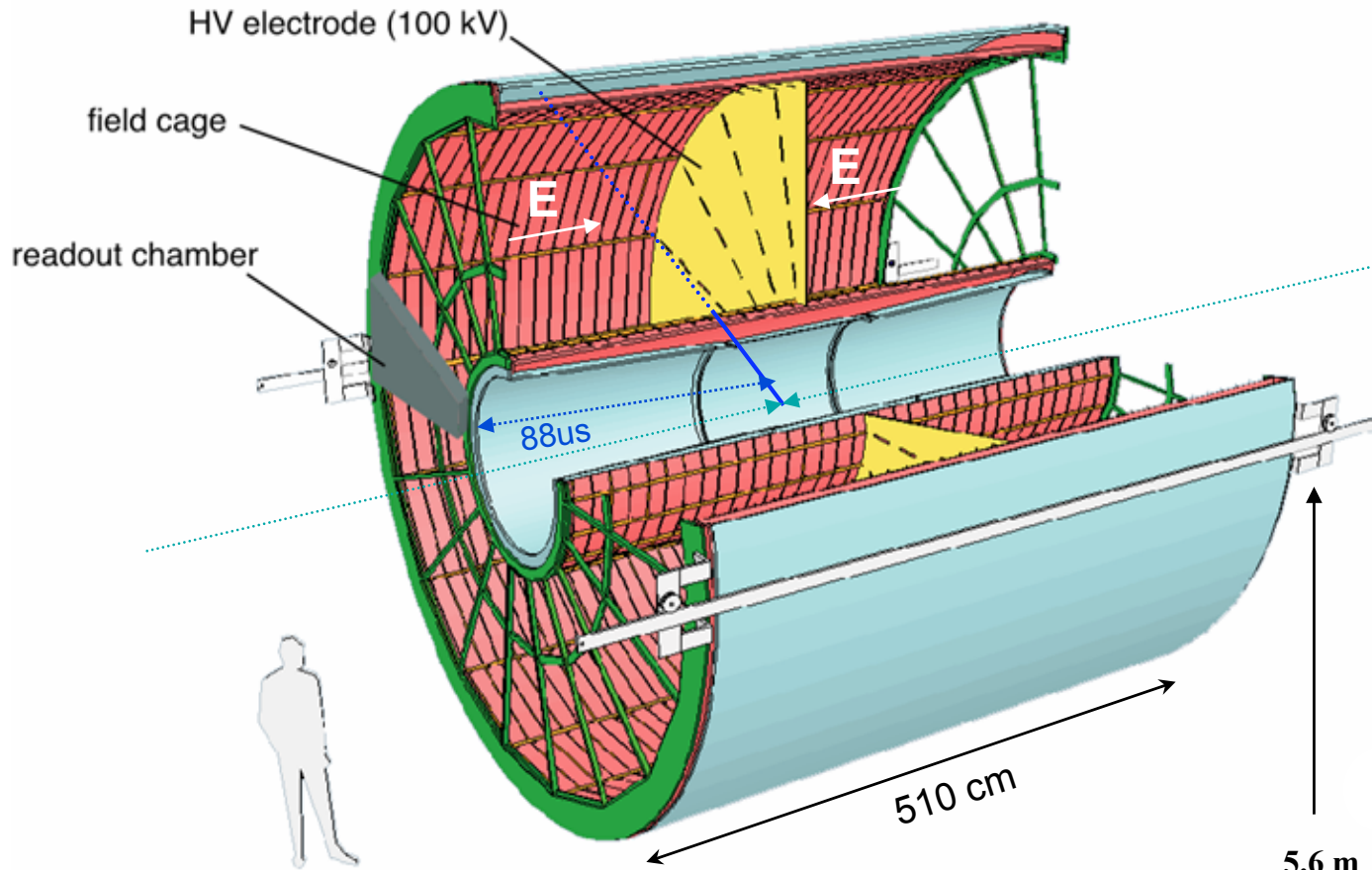
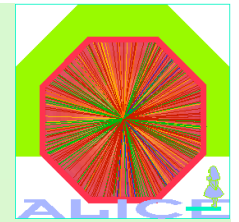
P-side orientation 7.5 mrad

N-side orientation 27.5 mrad



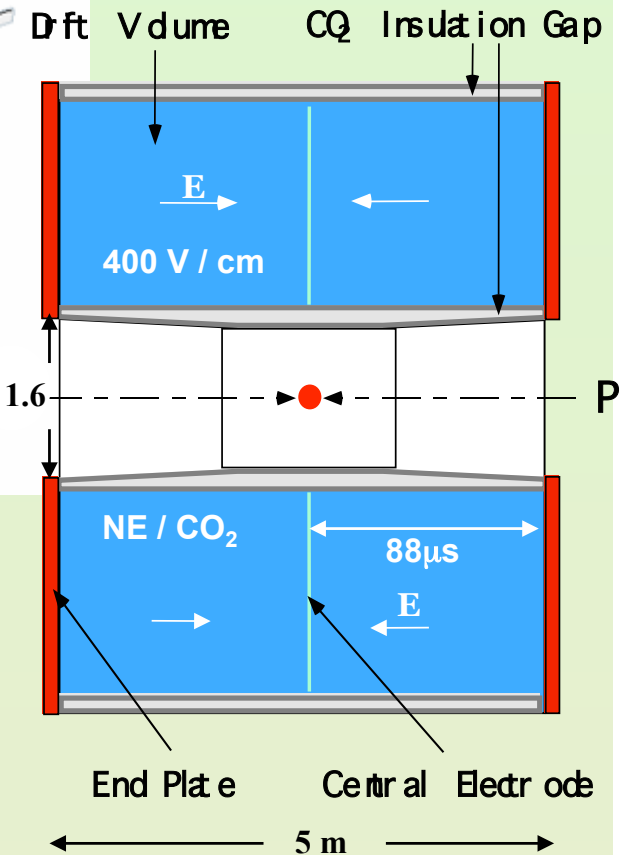
End ladder electronics

TPC layout

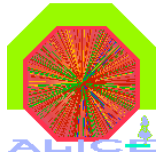


GAS VOLUME
88 m³

DRIFT GAS
90% Ne - 10%CO₂

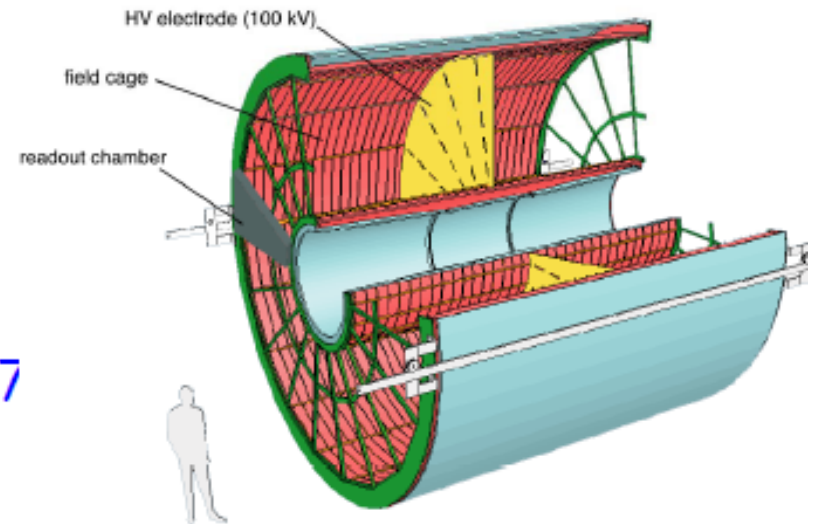


Readout plane segmentation
18 trapezoidal sectors
each covering 20 degrees in azimuth
570k channels

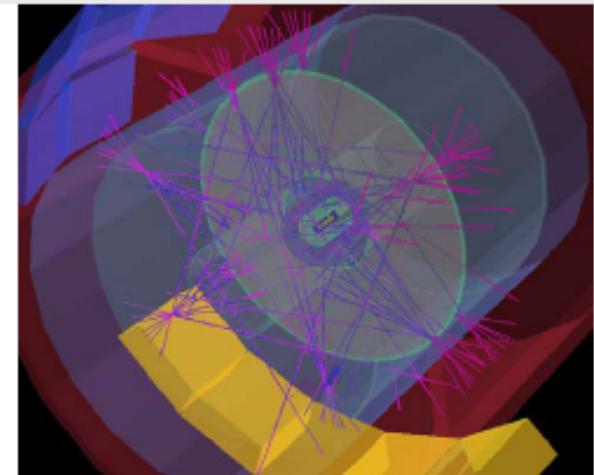


Time Projection Chamber (TPC)

- **Characteristics:**
 - ⇒ 85 m³ - NeC₂O₂N₂ gas mixture
 - ⇒ 557,568 readout channels
 - ⇒ Maximum drift time = 92 μs
 - ⇒ Many (>90) 3D points (+dE/dx) per track
- Installation in ALICE since 2007
- Running continuously from May to October 2008 and since August 2009
- **Calibration:**
 - ⇒ >750 million events (cosmics, krypton, and laser) recorded, with and without B
 - ⇒ First round of calibrations (dE/dx, momentum, alignment, gain) completed before p-p collisions

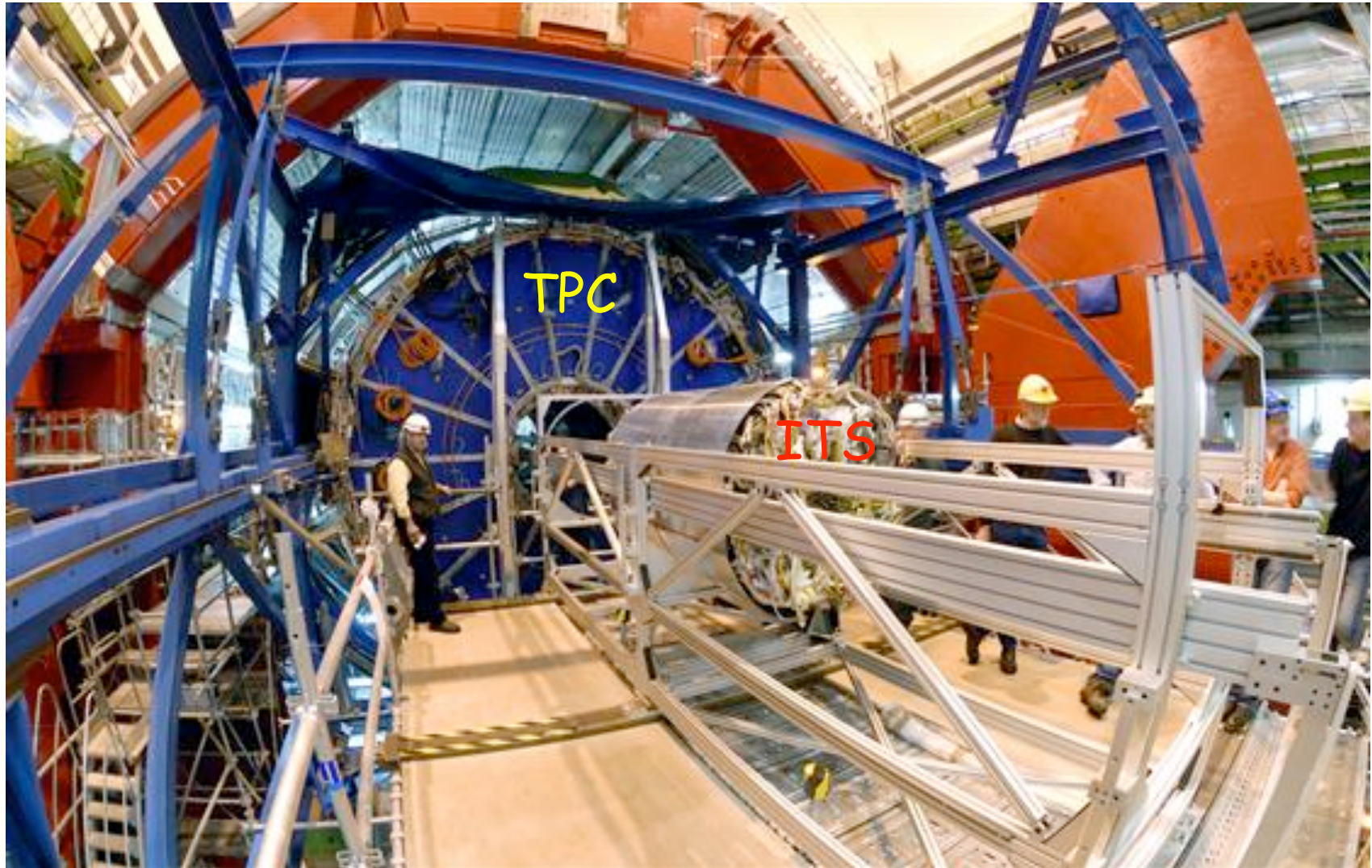


Laser event



TPC

Largest TPC ever: $R=2.5$ m, length=5m



MORE
SLIDES

The LHCb Detector

Muon System

RICH Detectors

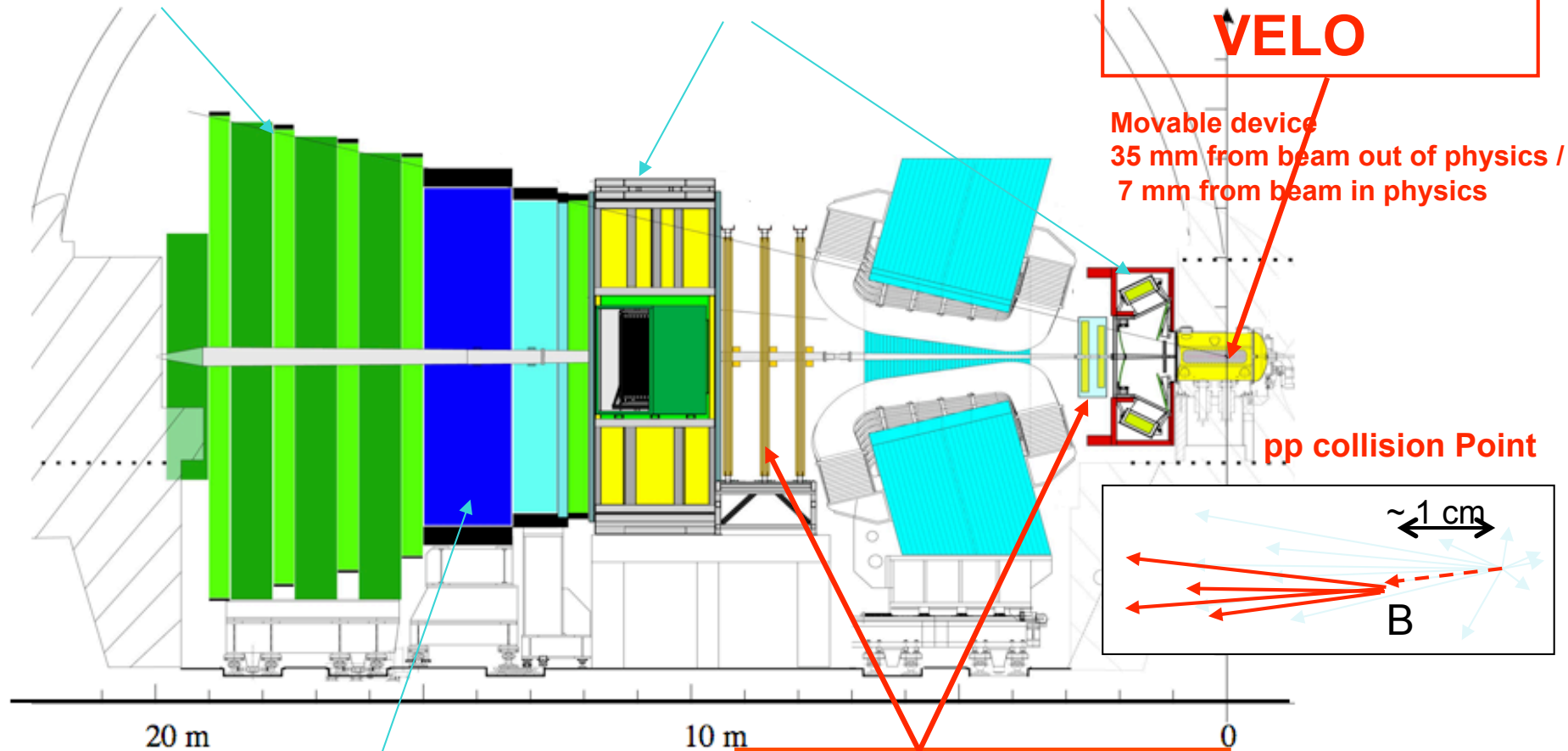
Vertex Locator
VELO

Movable device
35 mm from beam out of physics /
7 mm from beam in physics

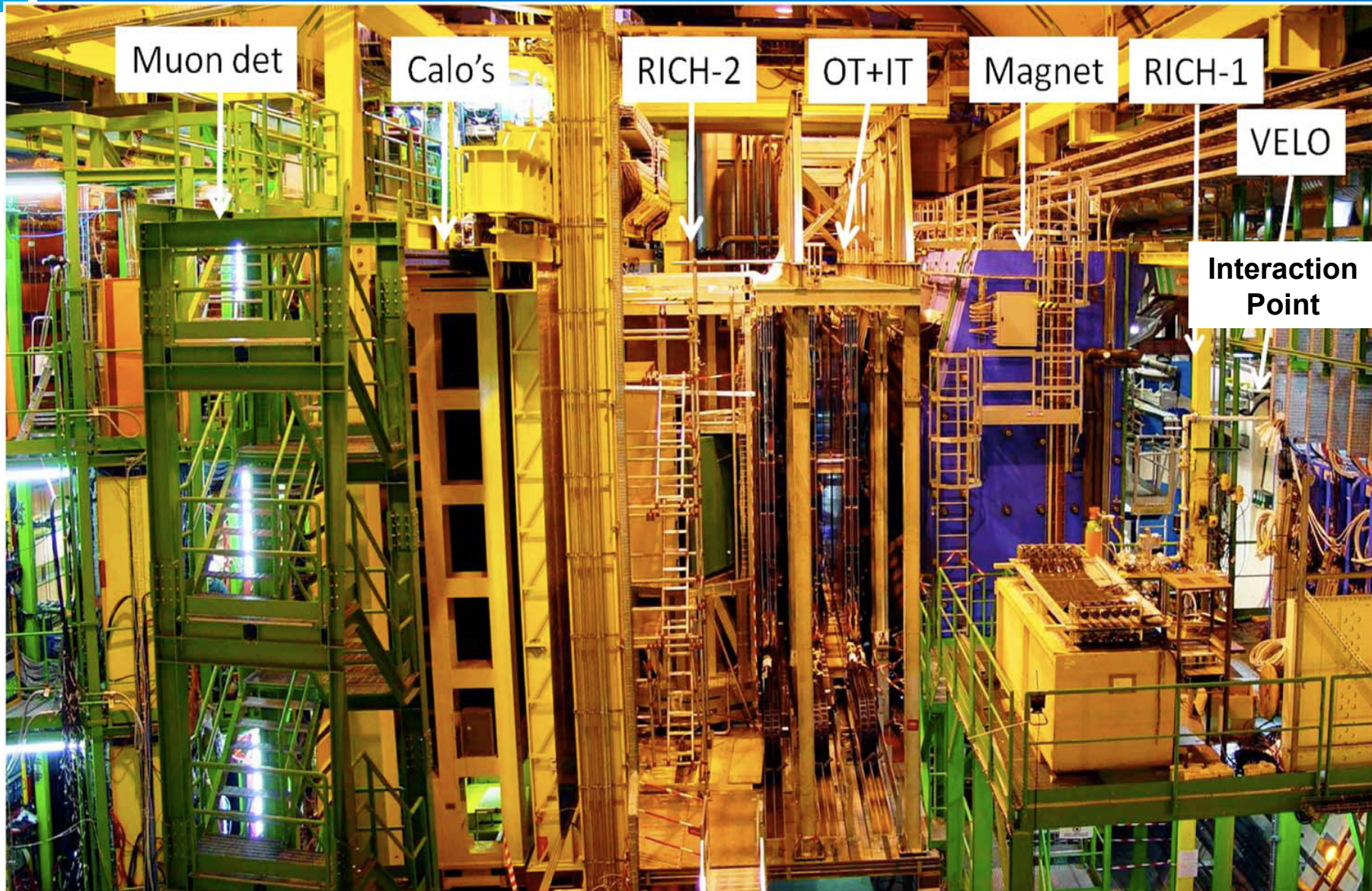
pp collision Point

Tracking System

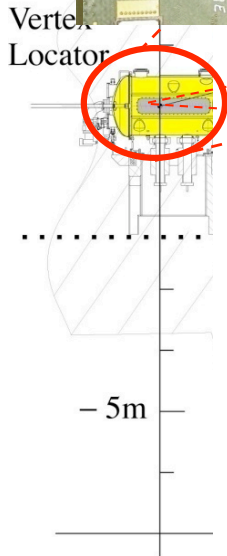
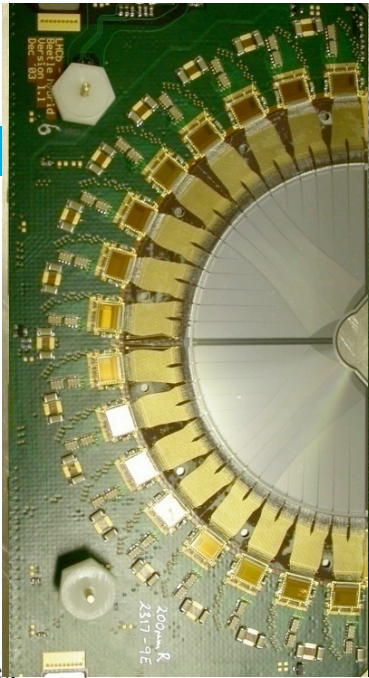
Calorimeters



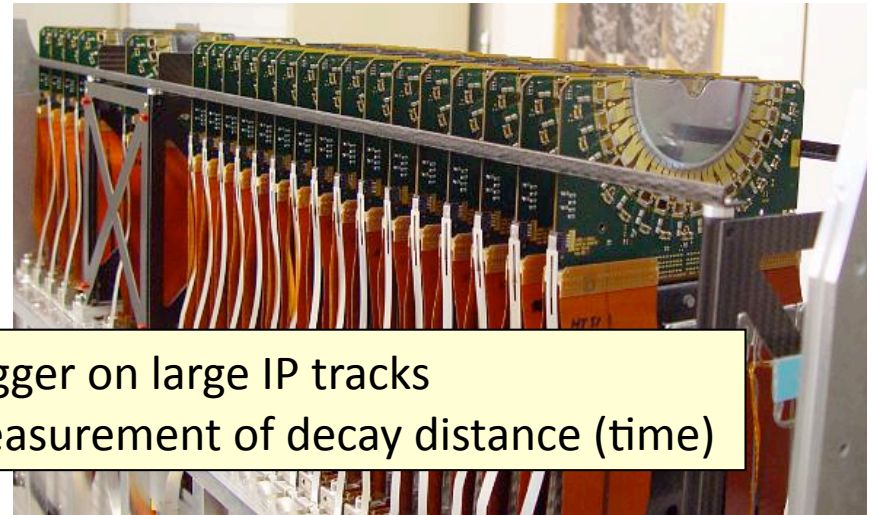
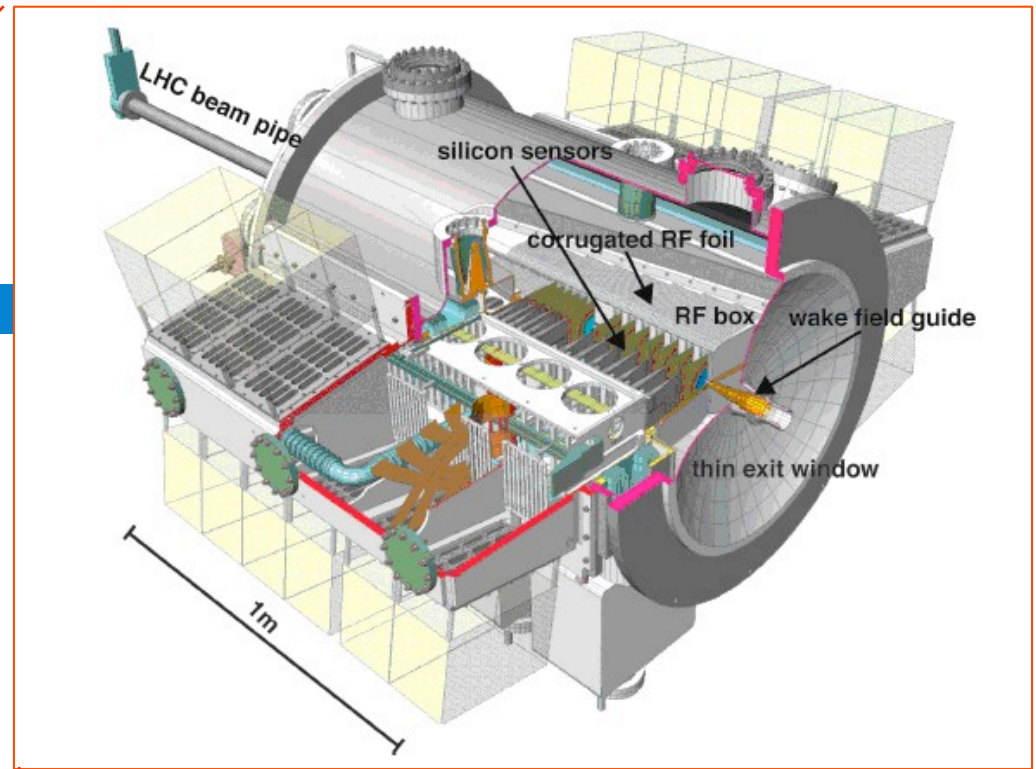
The LHCb Detector



B-Vertex Measurement

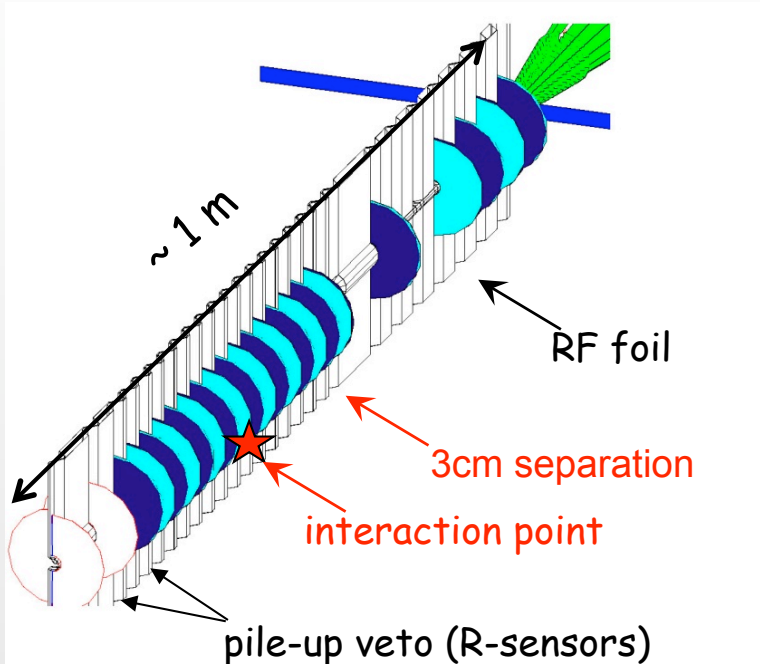


Vertex Locator (Velo)
21 stations of silicon strip detectors ($r-\phi$)
 $\sim 8 \mu\text{m}$ hit resolution
 $\sim 25 \mu\text{m}$ IP resolution

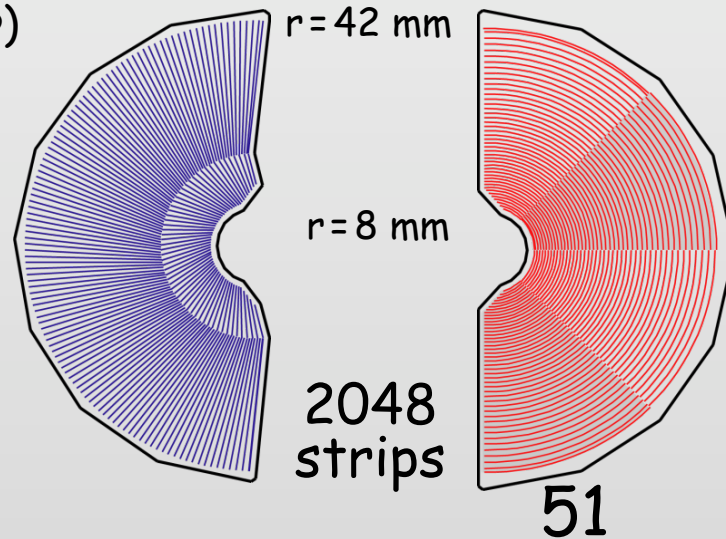


- Trigger on large IP tracks
- Measurement of decay distance (time)

VELO - VErteX LOcator

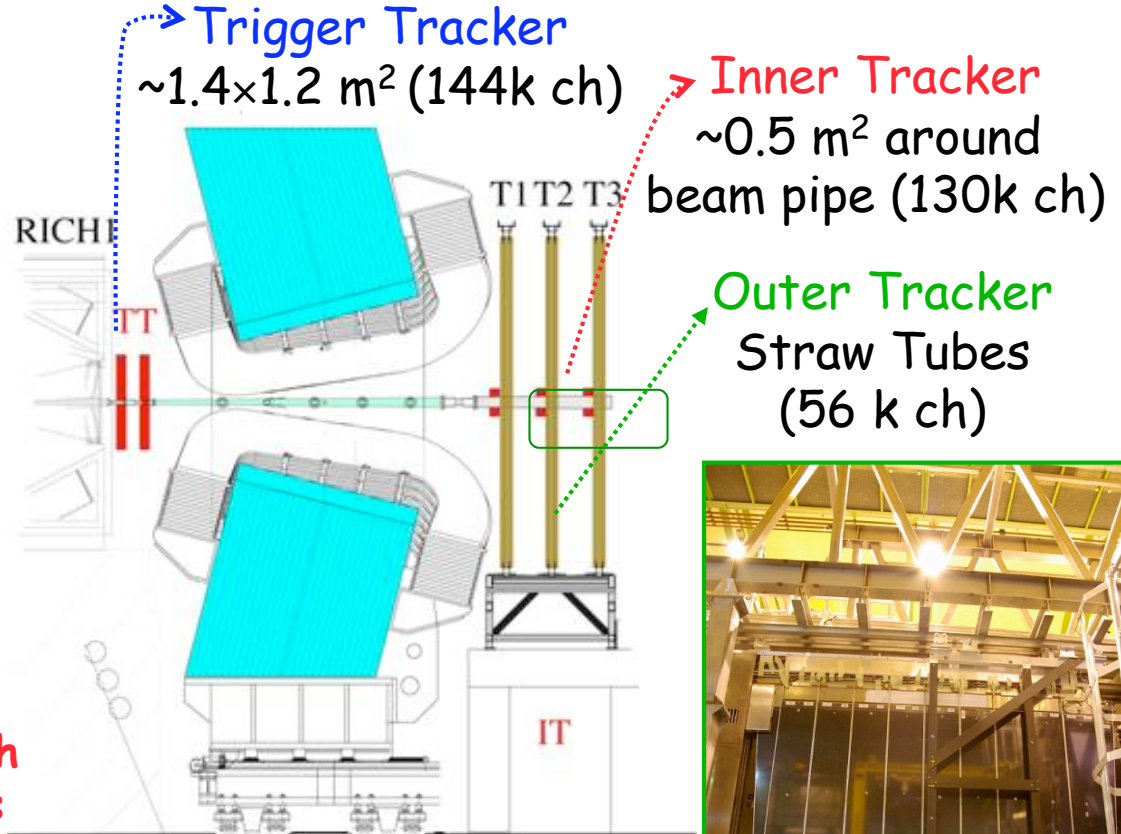
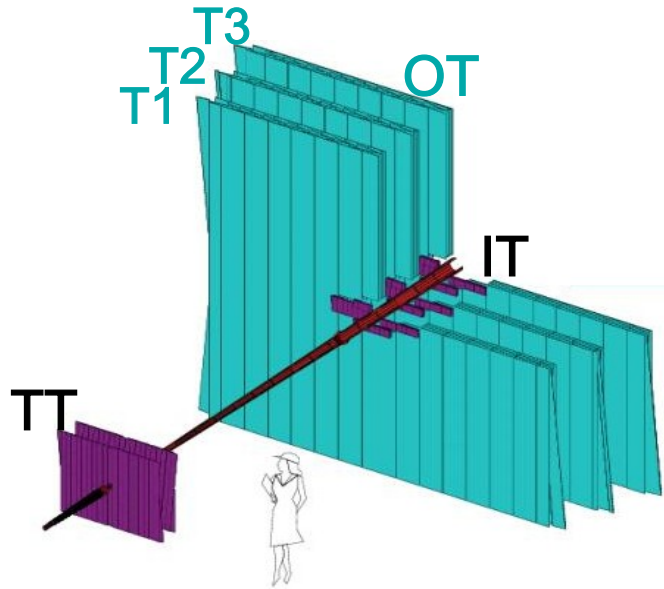


- Detector halves retractable (by 30mm) from interaction region before LHC is filled (to allow for beam excursions before injection and ramping)
- 21 tracking stations
- Unique R- Φ geometry, 40-100 μ m pitch, 300 μ m thick
- Optimized for
 - tracking of particles originating from beam-beam interactions
 - fast online 2D (R-z) tracking
 - fast offline 3D tracking in two steps (R-z then Φ)

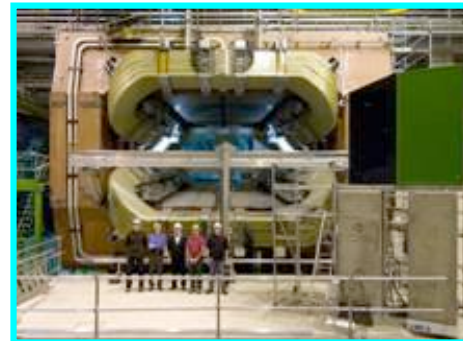
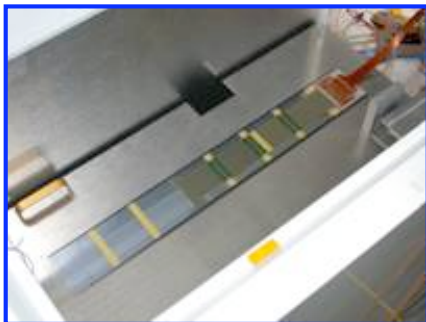


Tracking System

TT + 3 stations (T1, T2, T3), each with 4 detection planes (0°, +5°, -5°, 0°)



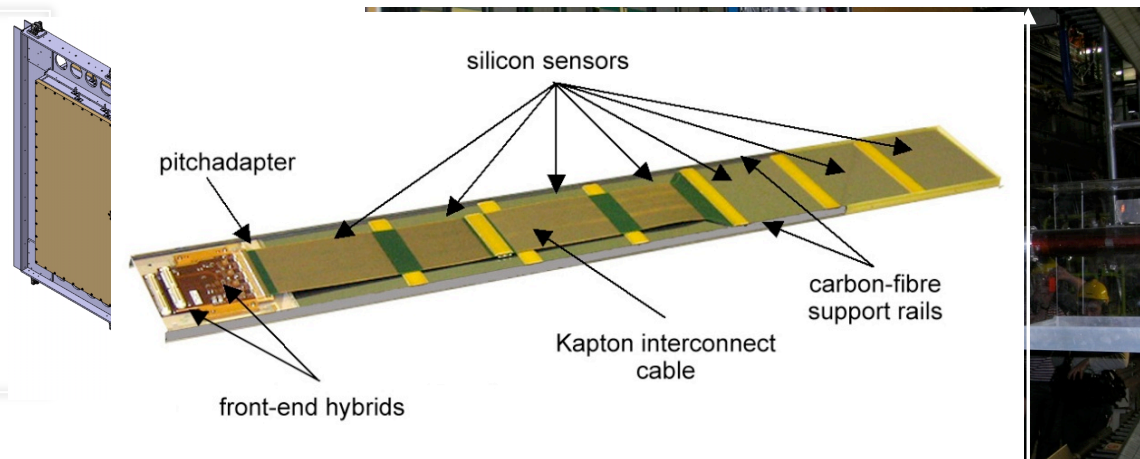
Similar sensors for TT & IT:
 Si μ -strip with pitch $\sim 200 \mu\text{m}$
TT: 128 Modules **IT: ladders with**
(7 Si sensors) **1 or 2 sensors**



Silicon Trackers

Trigger Tracker:

- 500 μm thick Si μ -strip sensors
- 7-sensor long ladders, 183 μm pitch
- Area of 8.2 m^2 covered with 896 sensors, 280 r/o sectors,
- 99.7% of channels functional



- Provides tracking info for triggering
- Tracking of low momentum particles

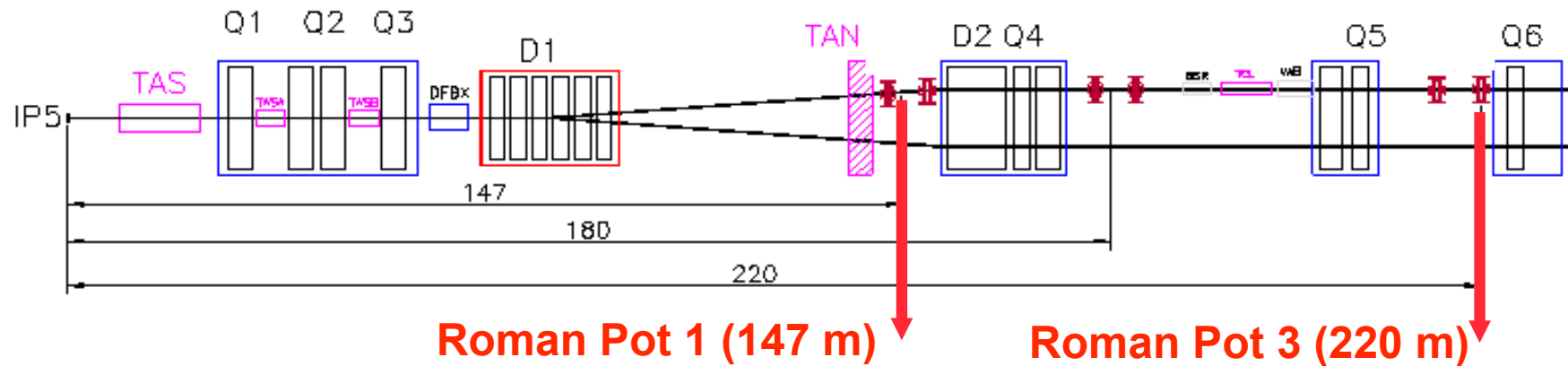
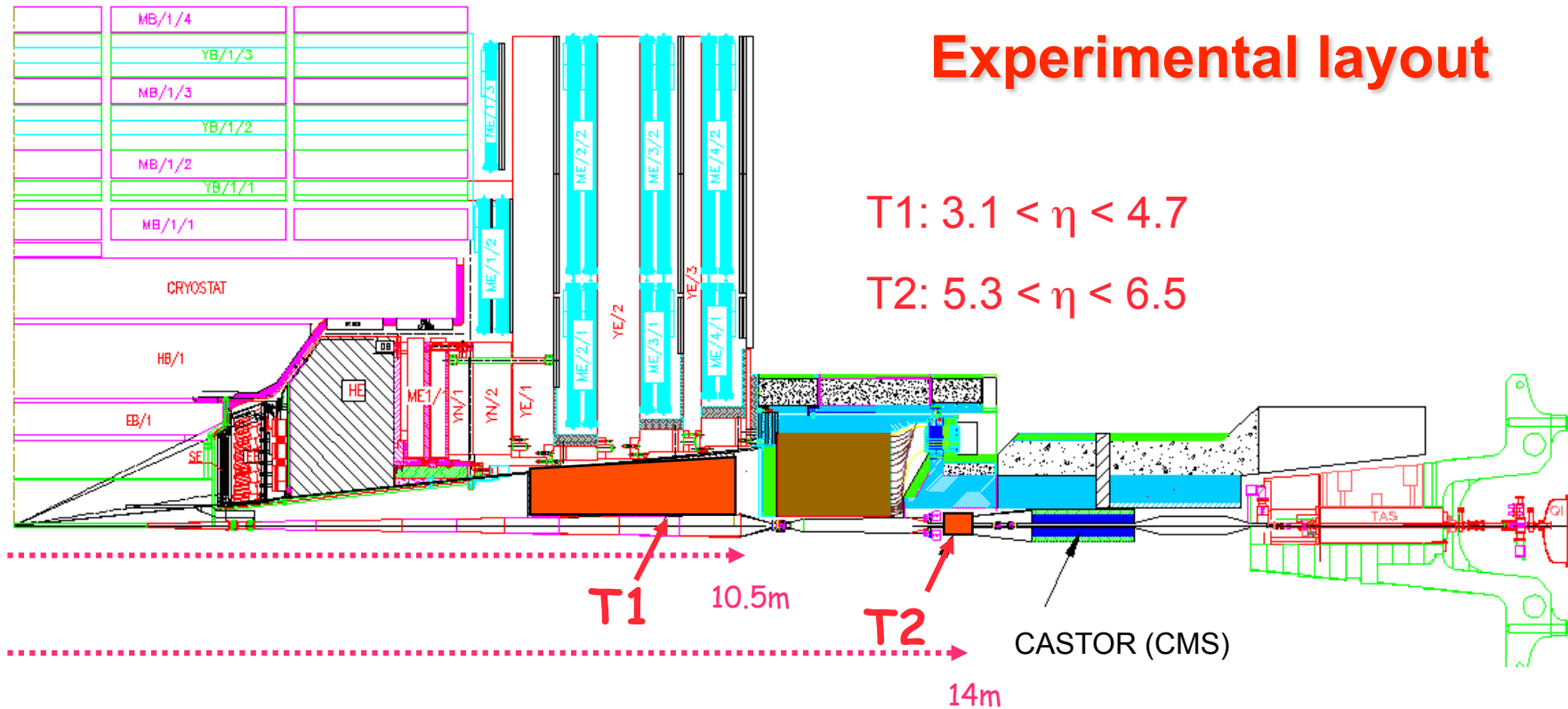
Inner Tracker:

- 320 (410) μm for 1 (2)-sensor ladders
- Readout pitch 198 μm pitch
- Area of 4.2 m^2 covered with 504 sensors, 336 ladders
- 99.4% of channels functional

- Provides tracking in high flux region ($5 \times 10^5 \text{cm}^{-2} \text{s}^{-1}$), 2% of area 20% of tracks

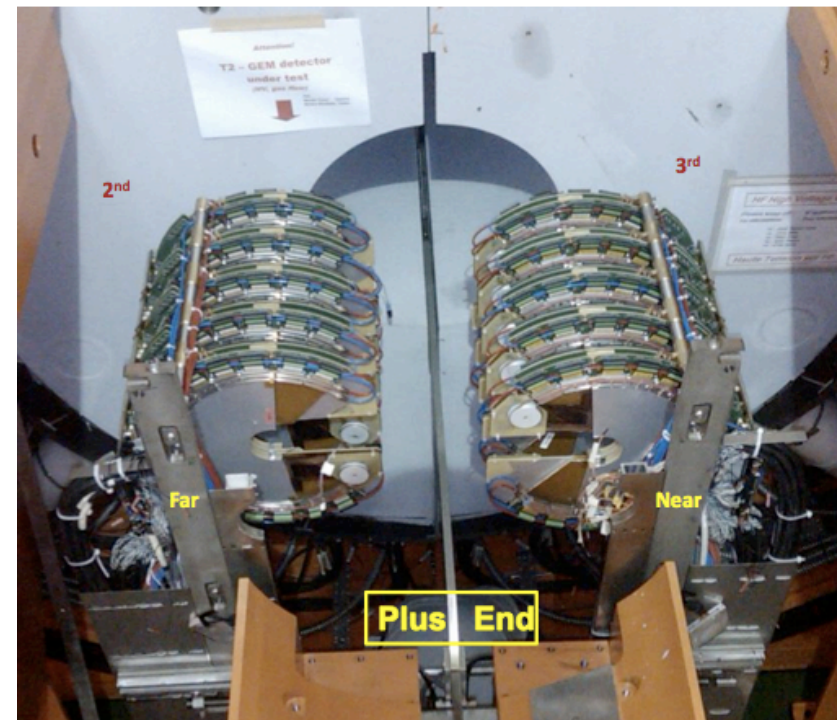
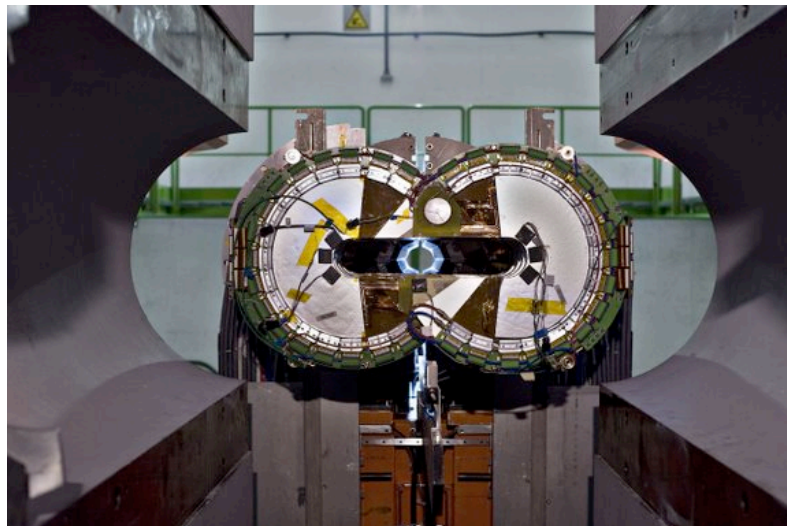
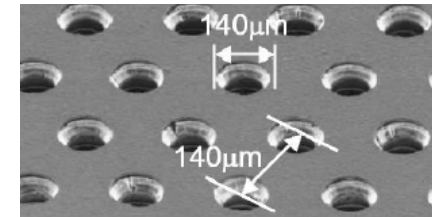
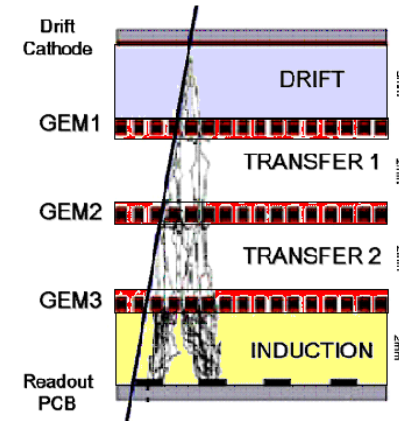


Experimental layout

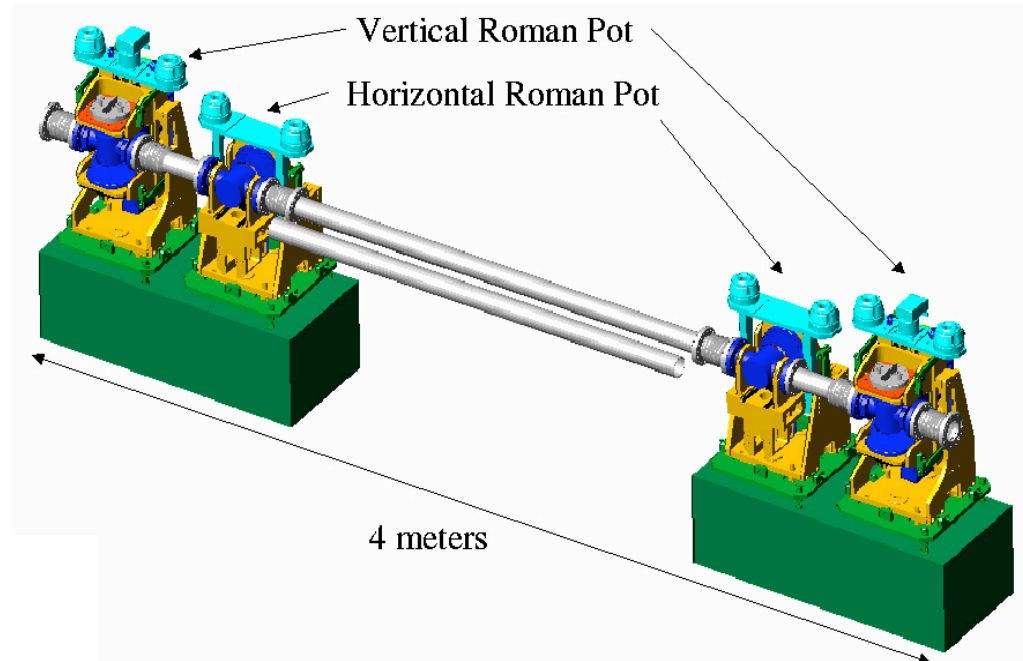
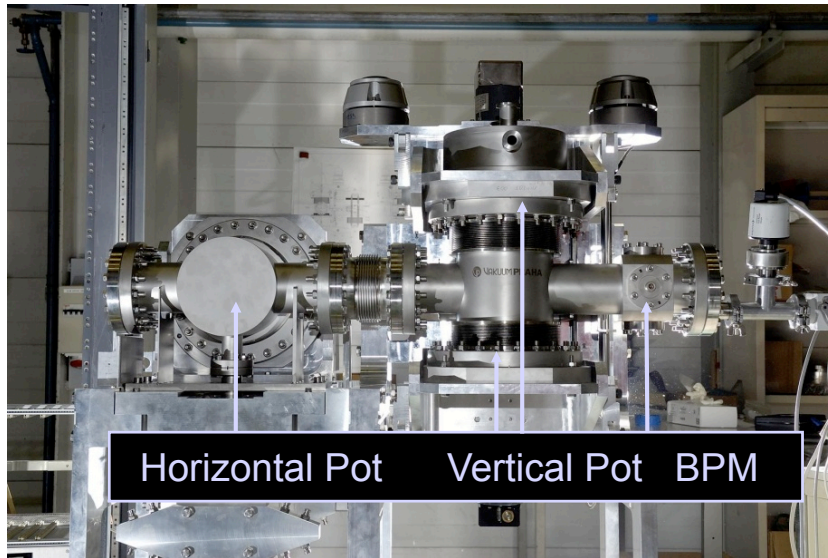


T2 Telescope

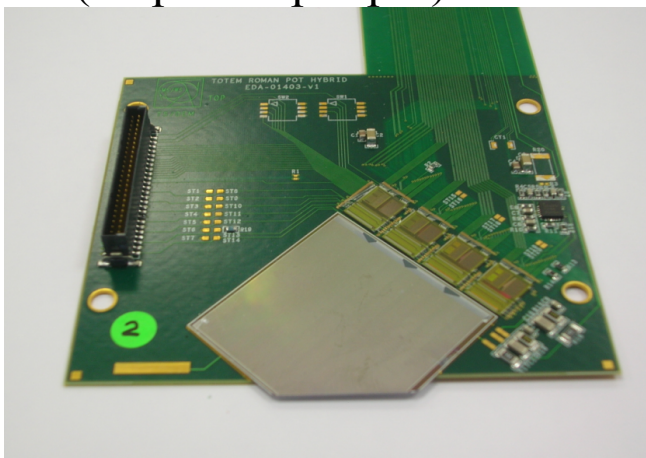
- ◆ Gas Electron Multiplier (GEM)
- ◆ $5.3 < |\eta| < 6.5$
- ◆ 10 half-planes @ 13.5 m from IP5
- ◆ Half-plane:
 - 512 strips (width $80 \mu\text{m}$, pitch of $400 \mu\text{m}$), radial coordinate
 - $65 \times 24 = 1560$ pads ($2 \times 2 \text{ mm}^2 \rightarrow 7 \times 7 \text{ mm}^2$), radial and azimuth coord.
 - Resolution: $\sigma(R) \sim 100 \mu\text{m}$, $\sigma(\varphi) \sim 1^\circ$
- ◆ Primary vertex reconstruction (beam-gas interaction removal)
- ◆ Trigger using (super) pads



Leading Proton Detection: TOTEM Roman Pots



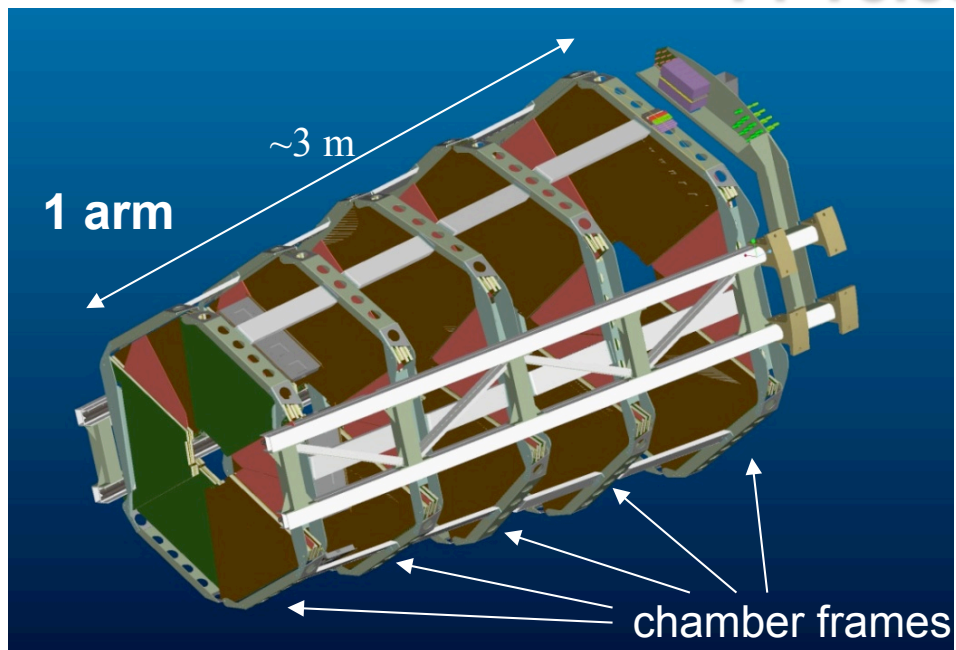
“edgeless” Si strip detectors
(10 planes per pot)



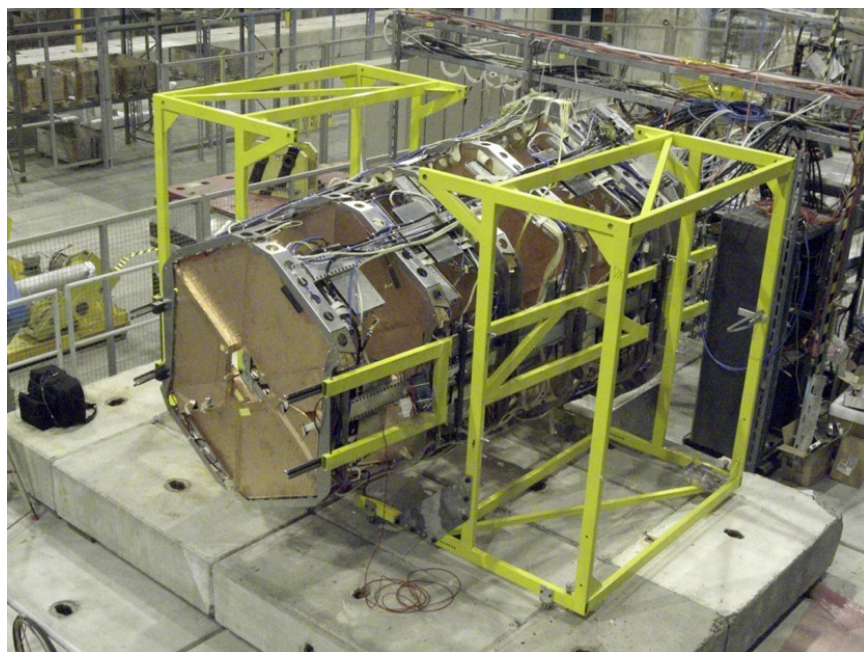
Leading proton detection at distances down $10 \sigma_{\text{beam}} + d$
Need “edgeless” detectors (efficient up to physical edge)
to minimise width d of dead space.

TOTEM: specially designed silicon strip detectors (CTS),
efficient within $50 \mu\text{m}$ from the edge

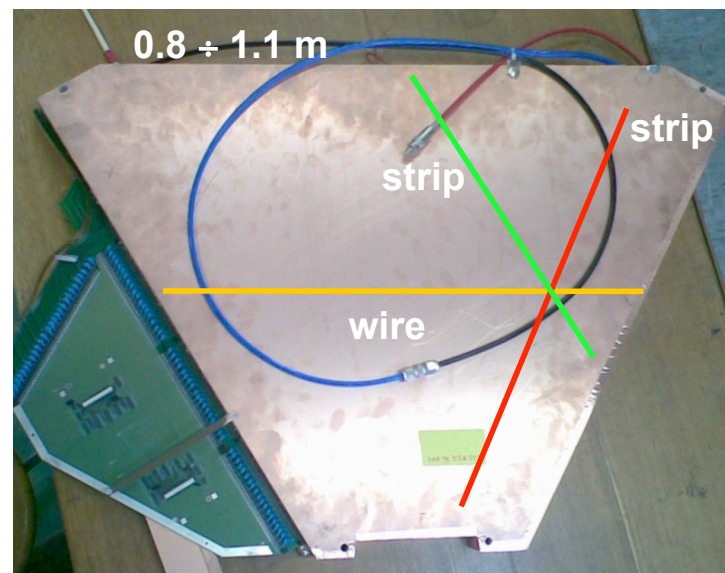
T1 Telescope



- ◆ Cathode Strip Chambers (CSC)
- ◆ $3.1 < |\eta| < 4.7$
- ◆ 5 planes with measurement of three coordinates per plane, $\sigma \sim 1$ mm
- ◆ Primary vertex reconstruction (beam-gas interaction removal)
- ◆ Trigger with anode wires
- ◆ Connected to VFAT chips
- ◆ Successful ageing studies
(~ 5 years at $L_{\text{inst}}=10^{30}$ cm $^{-2}$ s $^{-1}$)



Installation as soon as possible

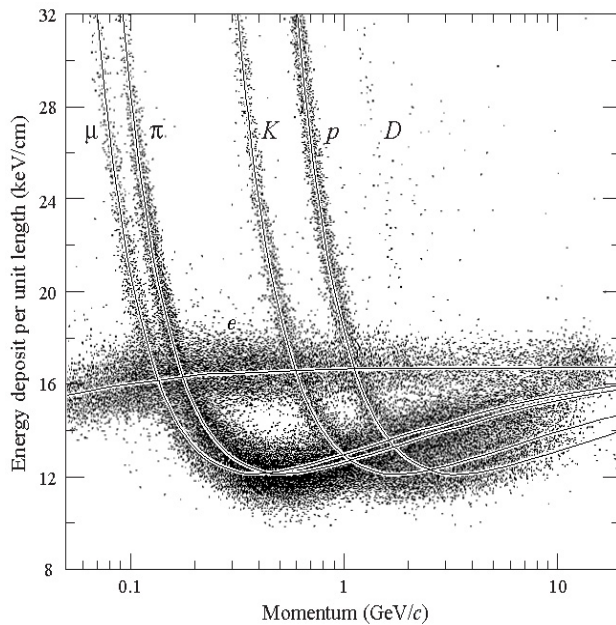
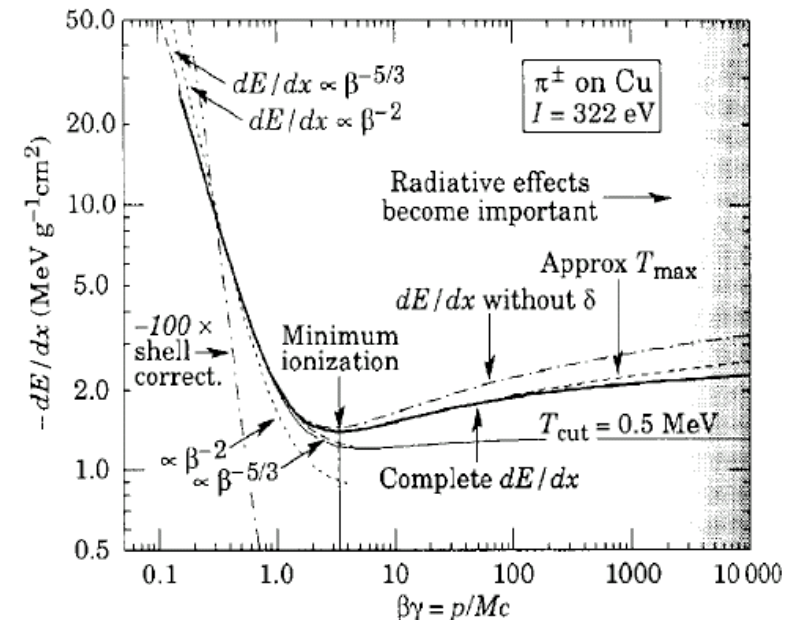


Backup slides

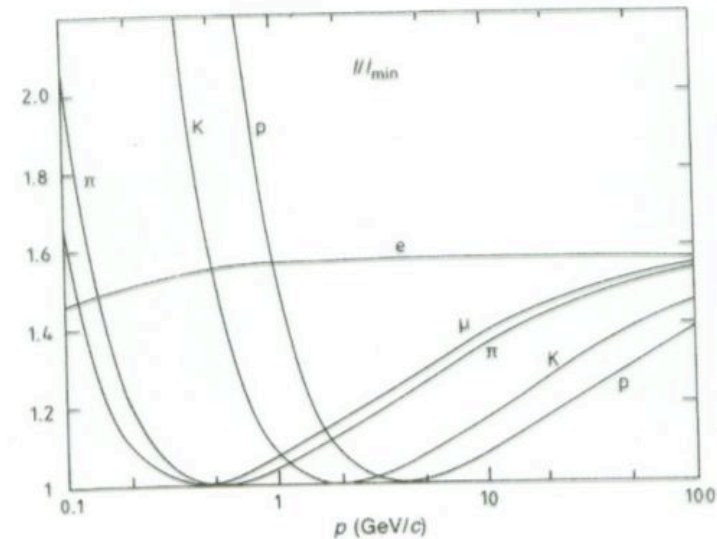
Particle Identification by Energy Loss

Energy loss depends on the $\beta\gamma$ of the particle and is \approx independent from the mass of the particle. As a function of particle momentum $p = Mc\beta\gamma$ the energy loss depends on the mass of the particle.

By measuring the energy loss and the momentum of the particle, the mass of the particle can be measured: \rightarrow Particle Identification !



TPC in magnetic field



$$\frac{1}{\rho} \frac{dE}{dx} = -4\pi r_e^2 m_e c^2 Z_1^2 \frac{p^2 + M^2 c^2}{p^2} N_A \frac{Z}{A} \left[\ln \frac{2m_e c^2 F}{I} \frac{p^2}{M^2 c^2} - \frac{p^2}{p^2 + M^2 c^2} \right]$$

Ionization: the Bethe-Bloch formula

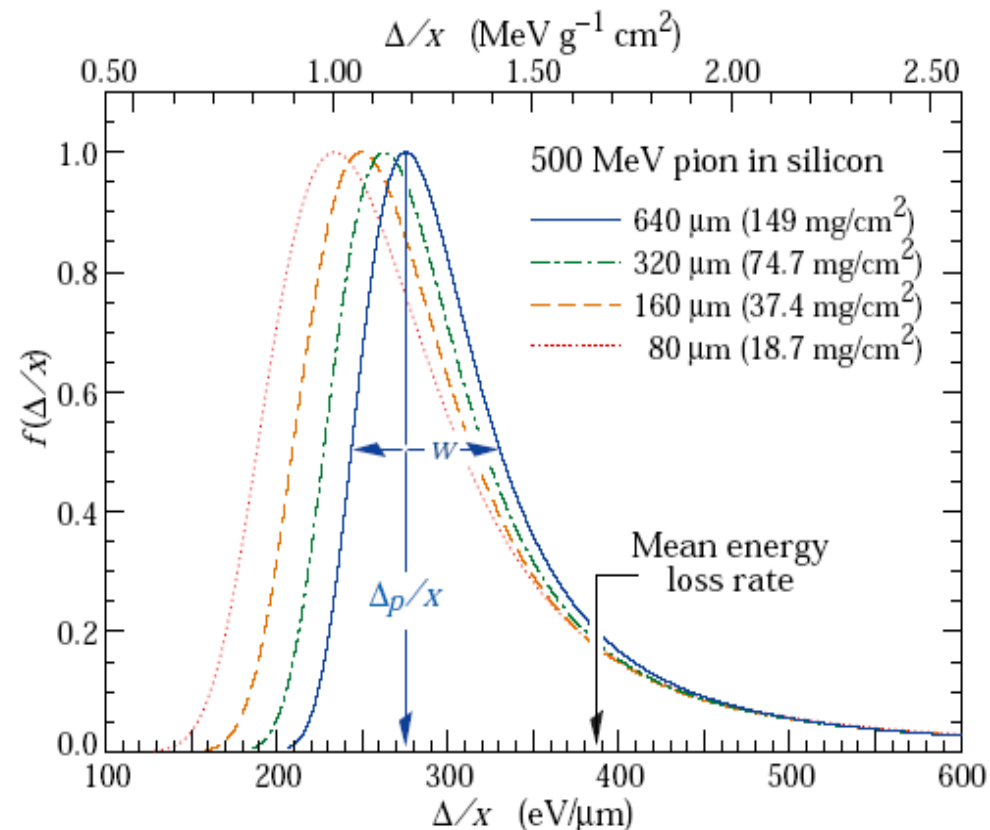
- Bethe-Bloch formula only gives the average energy loss, and do not take into account fluctuations from event to event.
- Large high energy tail - δ rays

δ -rays : electrons that have sufficient energy to ionize further atoms through subsequent interactions on their own.

Landau distribution:

$f(\Delta/X)$: Probability for energy loss Δ in a thickness X of matter.

Very asymmetric distribution:
average and most probable energy loss must be distinguished !



Multiple Scattering

A particle traversing material undergoes successive deflections due to multiple elastic scattering from nuclei.

The probability that the particle is deflected by an angle θ after travelling a distance x in the material is well approximated (actually tails are larger than Gaussian tails) by a Gaussian distribution with sigma of:

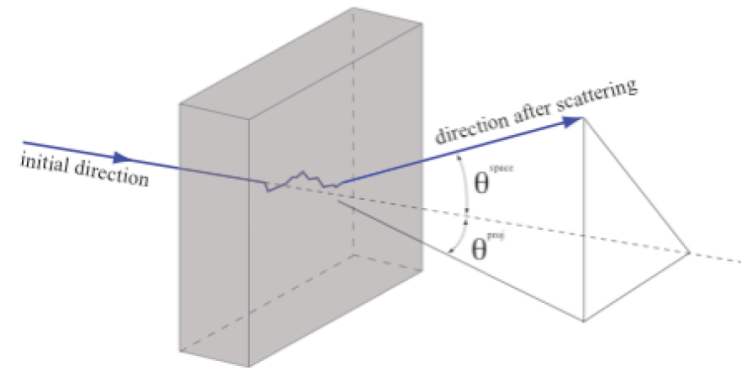
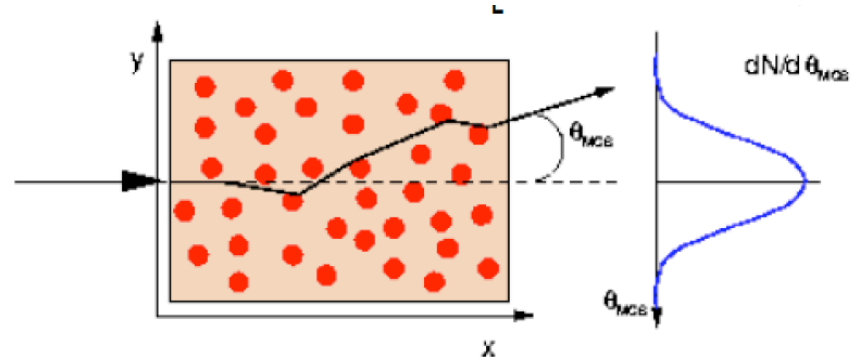
$$\theta_{MCS} = \theta_{rms} = \frac{13.6 \text{ MeV}}{\beta c p} z \sqrt{\frac{x}{X_0}} \left[1 + 0.038 \ln \left(\frac{x}{X_0} \right) \right]$$

- X_0 Radiation length of the material
- z Charge of the particle
- p Momentum of the particle

Radiation Length X_0 has 2 definitions:

- ◇ Mean distance over which high energy electron losses all but $1/e$ of its energy by Bremsstrahlung.
- ◇ 7/9ths of the mean free path for pair production by a high energy photon.

$$X_0 \approx \frac{716.4 \text{ g cm}^{-2} A}{Z(Z+1) \ln(287/\sqrt{Z})}$$



	X_0 (g cm ⁻²)	X_0 (cm)
Air	37	30,000
Silicon	22	9.4
Lead	6.4	0.56