

# Fisica Teorica 2 / Advanced Quantum Field Theory (the context)

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# From “Fisica Teorica 1 / Quantum Field Theory”

Historical reason for introducing QFTs:

quantum mechanics, special relativity, and electromagnetism

So far you have seen

QFT in the canonical formalism

applied to the computation of

cross sections using the lowest perturbative order (tree level)

These computations are meaningful only for “weakly interacting” theories, useful typically for some aspects of high-energy physics (e.g., QED or weak interactions).

## Modern formulation of QFTs (since $\simeq$ late '70s)

QFT is the theoretical framework needed to accurately describe the “low-energy” physics of systems in which fluctuations are important.

In this generalized context one often speaks of **quantum and statistical field theories**, depending whether fluctuations are of quantum or of thermal origin.

If fluctuations are not important classical field theory is enough (fluid mechanics, elasticity, electromagnetism, general relativity), moreover in some cases QFT is well approximated by quantum mechanics.

The meaning of “low-energy” depends on what we are actually studying. It can go from  $\ll 10^{19}$  GeV down to  $\ll 10^{-3}$  eV.

# Particles in the modern formulation of QFTs

Particles are the excitations of the underlying QFT  
(identical particles)

Particles can be either of **fundamental** nature

leptons and quarks in the Standard Model (so far)

or describe **effective excitations** of the system we are considering

pions in chiral effective theory, phonons in a crystal,  
electrons in a semiconductor, Cooper pairs in a superconductor

Both the cases can be described by a QFT as far as we limit ourselves to the appropriate energy range. **It is possible (likely) that today fundamental physics is tomorrow effective low energy physics.**

# How physical description can change with the scale

Lagrangian of the **fundamental** model:

$$L = \sum_i \left( \frac{1}{2} m \dot{x}_i^2 - \frac{\kappa}{2} (x_{i+1} - x_i - a)^2 \right)$$

Introducing  $\phi_i = x_i - x_i^{(\text{eq})}$  and considering distances  $\gg a$  (i.e., momenta  $\ll 1/a$ ) we get the **effective** theory described by the Lagrangian density ( $v_s^2 = \kappa/m$ )

$$\mathcal{L} = \frac{1}{2} \left( \partial_t \phi(t, x) \right)^2 - \frac{v_s^2}{2} \left( \partial_x \phi(t, x) \right)^2 = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi$$

**Original symmetry:** shift ( $x_i \rightarrow x_i + \Delta$ ;  $\phi(t, x) \rightarrow \phi(t, x) + \Delta$ )

**Emergent symmetries:** translation and Lorentz (with  $c \rightarrow v_s$ )

**Classical FT:** elasticity. **Quantum FT:** acoustic phonons.

# High-energy physics / condensed matter physics

The traditional (?) subdivision in high-energy physics or low-energy/condensed matter physics is useless and harmful in theoretical physics.

The modern understanding of several key concepts of theoretical physics has benefited greatly from the contamination between high/low energy perspectives.

## Spontaneous symmetry breaking

low-energy  $\rightarrow$  high-energy  $\rightarrow$  everything

## Higgs mechanism

high-energy  $\rightarrow$  low-energy  $\rightarrow$  high-energy  $\rightarrow$  everything

## Renormalization & renormalization group

high-energy  $\rightarrow$  low-energy  $\rightarrow$  everything

# High-energy: Standard Model

## Standard Model of Elementary Particles

		three generations of matter (fermions)						
		I	II	III				
mass		$\approx 2.4 \text{ MeV}/c^2$	$\approx 1.275 \text{ GeV}/c^2$	$\approx 172.44 \text{ GeV}/c^2$	0	$\approx 125.09 \text{ GeV}/c^2$		
charge		$2/3$	$2/3$	$2/3$	0	0		
spin		$1/2$	$1/2$	$1/2$	1	0		
		<b>u</b> up	<b>c</b> charm	<b>t</b> top	<b>g</b> gluon	<b>H</b> Higgs		
<b>QUARKS</b>		$\approx 4.8 \text{ MeV}/c^2$	$\approx 95 \text{ MeV}/c^2$	$\approx 4.18 \text{ GeV}/c^2$	0			
		$-1/3$	$-1/3$	$-1/3$	0			
		$1/2$	$1/2$	$1/2$	1			
		<b>d</b> down	<b>s</b> strange	<b>b</b> bottom	<b><math>\gamma</math></b> photon			
		$\approx 0.511 \text{ MeV}/c^2$	$\approx 105.67 \text{ MeV}/c^2$	$\approx 1.7768 \text{ GeV}/c^2$	$\approx 91.19 \text{ GeV}/c^2$			
		-1	-1	-1	0			
		$1/2$	$1/2$	$1/2$	1			
		<b>e</b> electron	<b><math>\mu</math></b> muon	<b><math>\tau</math></b> tau	<b>Z</b> Z boson			
<b>LEPTONS</b>		$< 2.2 \text{ eV}/c^2$	$< 1.7 \text{ MeV}/c^2$	$< 15.5 \text{ MeV}/c^2$	$\approx 80.39 \text{ GeV}/c^2$			
		0	0	0	$\pm 1$			
		$1/2$	$1/2$	$1/2$	1			
		<b><math>\nu_e</math></b> electron neutrino	<b><math>\nu_\mu</math></b> muon neutrino	<b><math>\nu_\tau</math></b> tau neutrino	<b>W</b> W boson			
							<b>SCALAR BOSONS</b>	
							<b>GAUGE BOSONS</b>	

**Caution:** the masses of the light quarks can not be interpreted as masses in the standard sense of the word.

## High-energy: Standard Model

So far the most economic and successful way to describe particle physics phenomenology.

Described by the QFT based on

- ▶ locality
- ▶ Poincarè invariance
- ▶ gauge symmetry  $U(1) \times SU(2) \times SU(3)$   
(sort-of nonAbelian generalization of QED)
- ▶ renormalizability (more to be said later on)

Electroweak sector dates back to 1967, the strong sector was completed in 1973 (asymptotic freedom and third quark family).

## Accurate results in the EW sector

Magnetic dipole moment and  $g$ -factor:  $\mathbf{M} = g \frac{e}{2m} \mathbf{S}$

For **elementary spin 1/2 particles**: tree level  $g = 2$ ,  
one loop  $a \equiv \frac{g-2}{2} = \frac{\alpha}{2\pi} \simeq 0.001\,161\,...$  (Schwinger '48)

$a_e^{(\text{exp})} = 0.001\,159\,652\,180\,59(13)$  and 5 loop QED used to fix  $\alpha$

$$a_\mu^{(\text{exp})} = 0.001\,165\,920\,72(15)$$

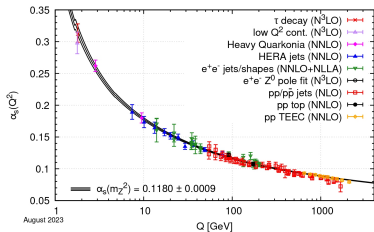
$$a_\mu^{(\text{th})} = 0.001\,165\,920\,33(62)$$

$$(a_\mu^{(\text{exp})} - a_\mu^{(\text{th})}) \times 10^{11} = 39 \pm 62_{\text{th}} \pm 15_{\text{exp}} \text{ Aliberti et al. 2505.21476}$$

“New physics” bounds

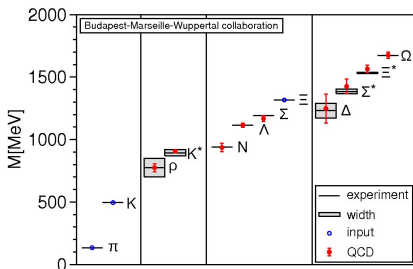
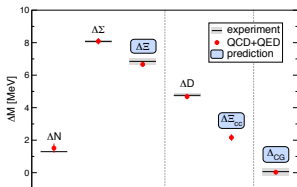
$$\frac{e}{2\Lambda} \bar{\psi} [\gamma_\mu, \gamma_\nu] \psi F^{\mu\nu} \quad \longrightarrow \quad \Lambda \gtrsim 10^8 \text{ GeV}$$

# Accurate results in the strong sector



Perturbative running coupling constant  $\alpha_s(Q^2)$  in  $\overline{\text{MS}}$  scheme at 5 loops compared with experiments

Hadron spectrum from numerical simulations of lattice QCD compared with experiments



# Condensed matter physics

Condensed matter phenomena are typically characterized by the following **properties**:

- ▶ large number of degrees of freedom  $\approx O(10^{23})$ , problems has to be addressed using statistical methods (**stat. mech.**)
- ▶ the degrees of freedom describing the low-energy excitations are often very different from the microscopic fundamental ones (**emergence**)
- ▶ different systems may show common collective behaviors although they differ microscopically (**universality**)

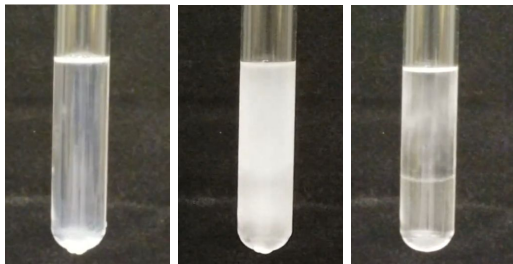
**Examples**: magnetism, theory of metals, superconductivity,  $^3\text{He}$  and  $^4\text{He}$  quantum liquids, quantum Hall effects, optical lattices,...

Depending on the physical setup **thermal** or **quantum** fluctuations can be the most important ones.

## Continuous thermal phase transitions

Characterized by the divergence of the correlation length  $\xi$  of the critical modes:  $\xi \sim (T - T_c)^{-\nu}$ . Also  $C_V \sim |T - T_c|^{-\alpha}$ ,  $\chi \sim |T - T_c|^{-\gamma}$ ,  $G(x) \sim |x|^{-1-\eta}$ , ...

Critical opalescence happens when  $\xi \gtrsim$  visible wavelength



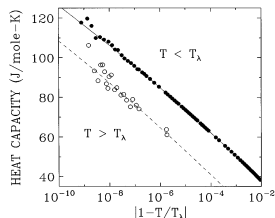
Mixture of methanol and cyclohexane heated at 65-70°C  
(by "Chemistry and Biochemistry Demo lab at OSU",  
see <https://www.youtube.com/watch?v=DIGdbmJvFUw> )

# Continuous thermal phase transitions: precision measures

$C_V$  at the normal-superfluid transition in  $^4\text{He}$ , up to a few nK from  $T_C$ .

order parameter: cond. wavefunction complex field  $\varphi(x)$  with U(1) symmetry

QFT:  $\mathcal{L} = |\partial_\mu \varphi|^2 + r |\varphi|^2 + u |\varphi|^4$



Results for the 3d XY universality class

		$\nu$	ref.
<b>Experiment</b>	$^4\text{He}$	0.66758(6)	Physica B <b>284-288</b> (2000), 49
	$^4\text{He}$	0.6709(1)	cond-mat/0310163
<b>QFT</b>	6,7-loop 3d-exp	0.6703(15)	cond-mat/9803240
	7-loop $\epsilon$ -exp	0.67076(38)	2005.12714
	conformal bootstrap	0.67175(10)	1912.03324
<b>Stat mech</b>	MC+HT	0.6717(1)	cond-mat/0605083
	MC	0.671718(23)	2507.19265

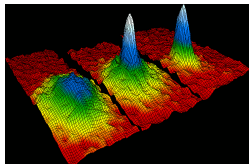
# Phase transitions in the quantum regime

BEC in bosonic gases, when

$$\lambda_{\text{deB}} = \left( \frac{h^2}{mT} \right)^{1/2} \approx d_{\text{atoms}} = (N/V)^{-1/3}$$

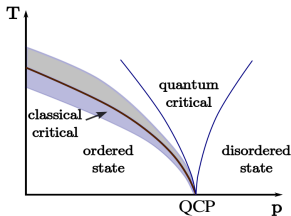
Observed in weakly interacting gases of alkali atoms (rubidium, sodium, lithium, ...).

velocity distribution of rubidium atoms  $\longrightarrow$



At a **quantum critical point** scaling laws describe the **interplay between quantum and thermal fluctuations**.

Interesting phenomenology for quantum magnetism, high- $T$  superconductors, quantum many particle systems, ...



# Why are QFTs so effective?

The main ingredients behind the success of QFTs are

1. **symmetries** (exact and/or spontaneously broken)
2. **renormalization** and **renormalization group**

**Symmetries** constraint the dynamics and the spectrum of a theory. In the presence of **spontaneous symmetry breaking** massless or light excitations can be present, whose interactions are greatly constrained by symmetry arguments.

**Renormalization group** in the Wilson formulation is a systematic way of **disentangling the low-energy features of a theory from its high-energy details**.

**QFTs work when different scales exist**, providing effective descriptions at low-energy, even when the microscopic details are unknown.

Modern point of view on renormalizability: **non-renormalizable interactions are irrelevant at low-energy**.

## Reductionism or not?

**P. A. M. Dirac** (?) "... and the rest is chemistry!"

**V. Weisskopf** "Intensive research goes for the fundamental laws, extensive research goes for the explanation of phenomena in terms of known fundamental laws"

**P. W. Anderson** "The behavior of large and complex aggregates of elementary particles, it turns out, is not to be understood in terms of a simple extrapolation of the properties of a few particles. Instead, at each level of complexity entirely new properties appear, and the understanding of the new behaviors requires research which I think is as fundamental in its nature as any other"

It is rather ironic to note that QFTs have been initially developed for fundamental physics, but the modern approach to QFTs is that they are effective descriptions of something more fundamental.

## Reductionism or not?

**Q:** Doesn't all this mean that quantum field theory, for all its successes, is an approximation that may have little to do with the underlying theory? And isn't renormalization a bad thing, since it implies that we can only probe the high energy theory through a small number of parameters?

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**A:** Nobody ever promised you a rose garden.

J. Polchinski "Effective Field Theory and the Fermi Surface"  
hep-th/9210046

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